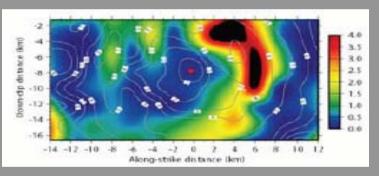
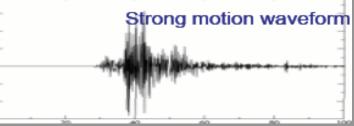
# Fault governing laws (constitutive equations)

# Seismologists need traction

- To apply fracture mechanics on mathematical planes representing the fault surfaces;
- To numerically simulate the spontaneous rupture nucleation, propagation, healing and arrest in dynamic earthquake models;

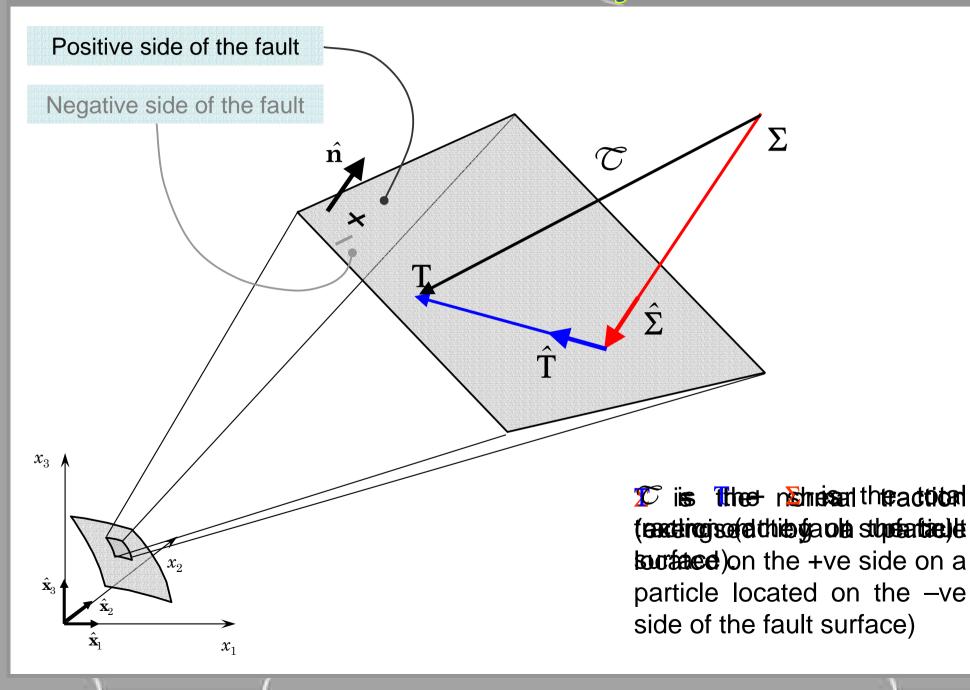


To model seismic wave propagation in the surrounding medium;



To predict ground shaking.

# Notations and symbols



 $\mathcal{C}^{(\hat{n})} = \mathbf{T}^{(\hat{n})} + \boldsymbol{\Sigma}^{(\hat{n})}$   $\underbrace{\text{total}}_{j} \text{ traction (acting on the fault surface).}$   $\mathcal{C}^{(\hat{n})}_{j} = n_{j}\sigma_{jj}^{\text{eff}}$   $Cauchy' \text{ s formula, where } \mathcal{C}^{(\hat{n})} = (\mathcal{C}^{(\hat{n})}_{1}, \mathcal{C}^{(\hat{n})}_{2}, \mathcal{C}^{(\hat{n})}_{3}),$   $\mathbf{n} = (n_{1}, n_{2}, n_{3}) \text{ and}$ 

$$\sigma_{ij}^{eff} = \sigma_{ij} + p_{fluid} \,\delta_{ij} = \begin{bmatrix} -\sigma_{n_1}^{eff} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & -\sigma_{n_2}^{eff} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & -\sigma_{n_3}^{eff} \end{bmatrix}$$

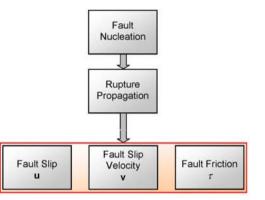
where:  $\sigma_{n_i}^{\text{eff}} = \sigma_{n_i} - p_{\text{fluid}} = -\sigma_{ii} - p_{\text{fluid}}$  and stresses are assumed to be negative for compression

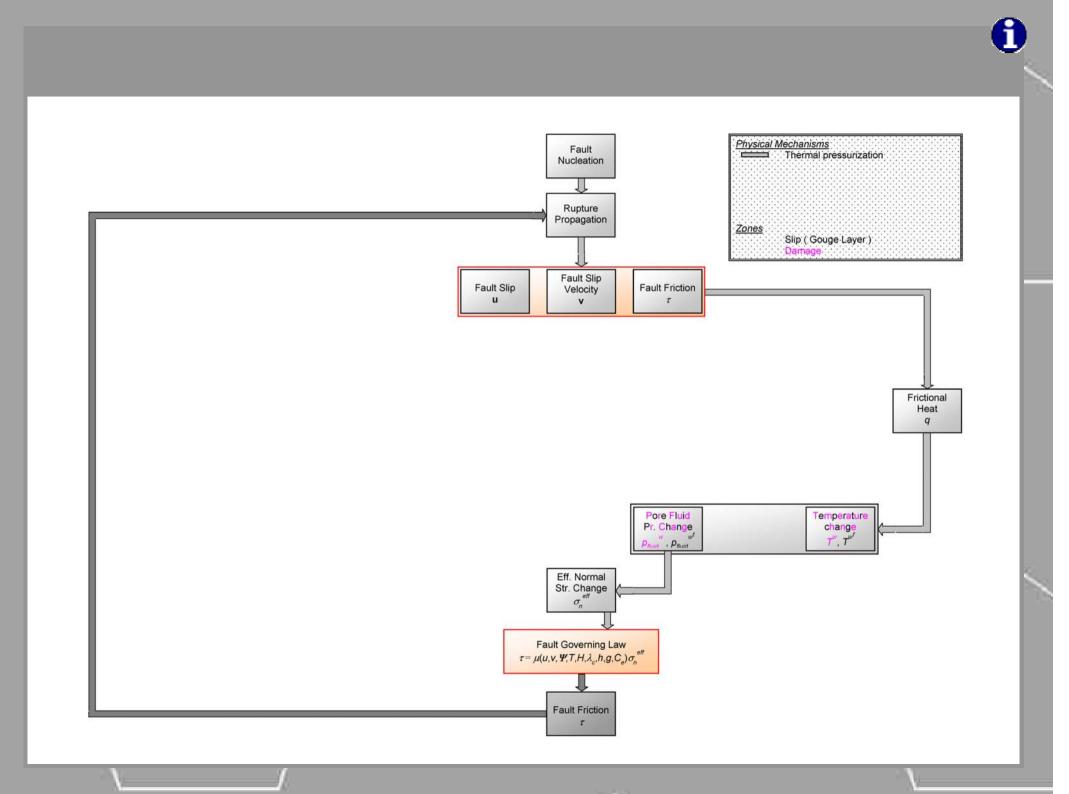
 $T_{j}^{(\hat{\mathbf{n}})} = n_{i}\sigma_{ij}^{\text{eff}} - n_{j}(n_{i}\sigma_{ik}^{\text{eff}}n_{k}) \qquad \text{shear traction}$ 

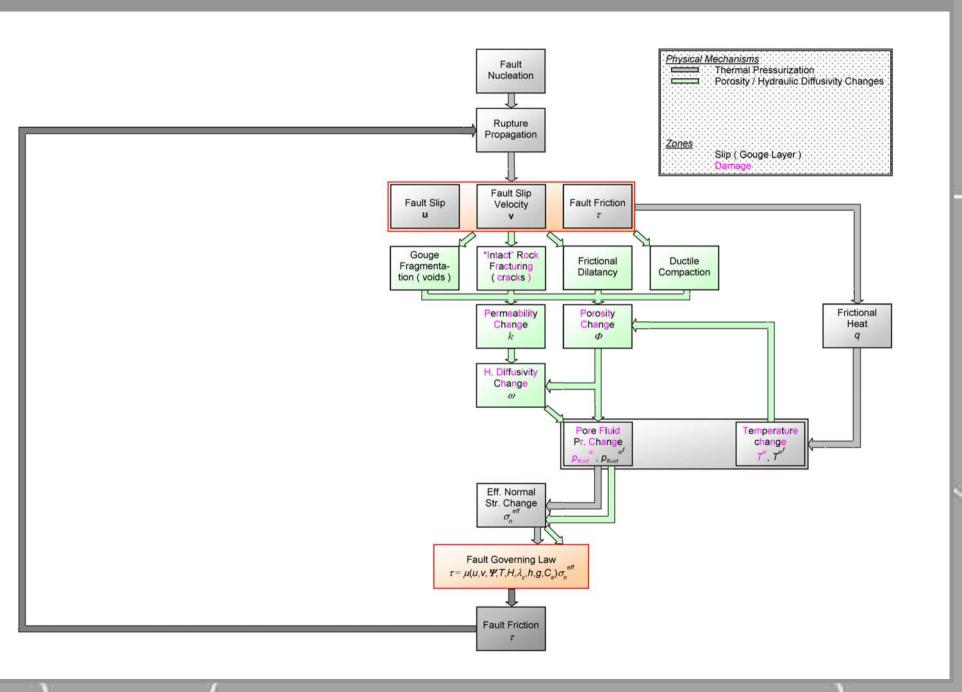
 $\Sigma_i^{(\hat{\mathbf{n}})} = n_i (n_i \sigma_{ik}^{eff} n_k)$ 

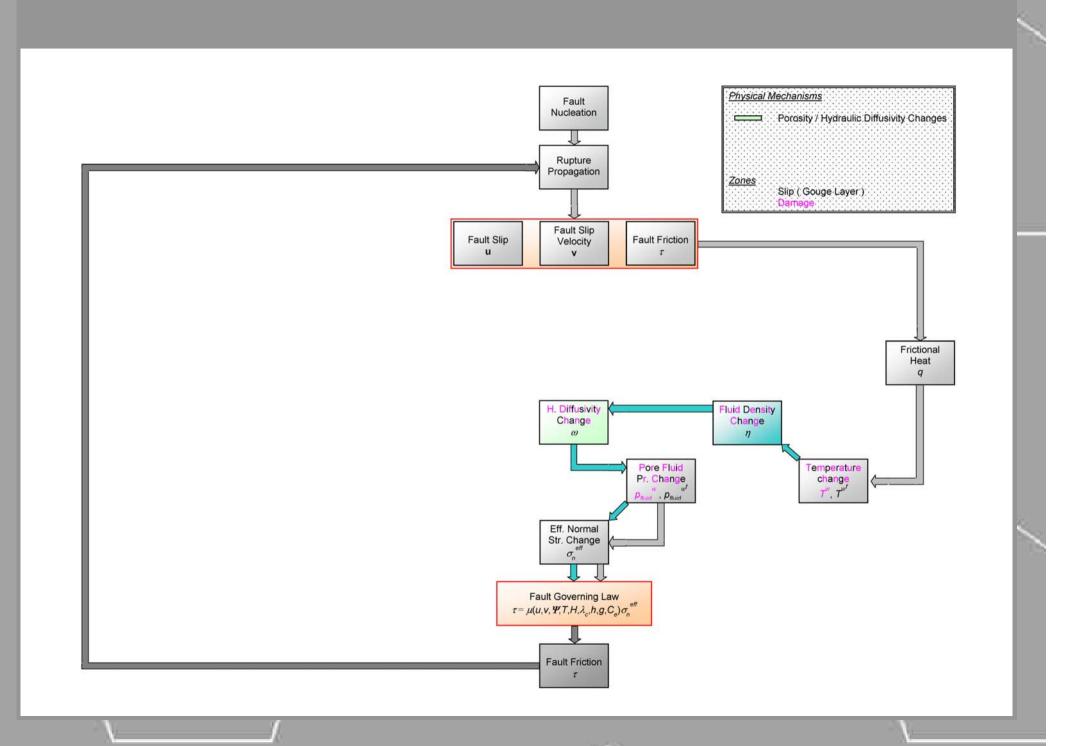
normal traction

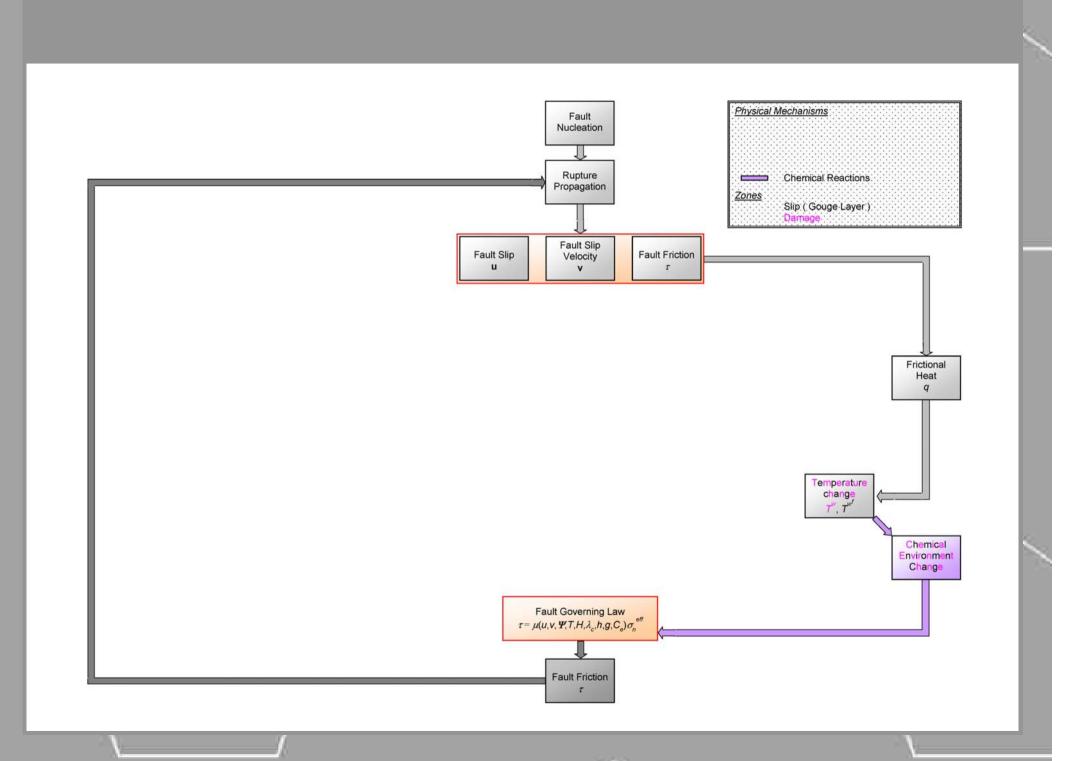
# **Physical Phenomena in Faulting**

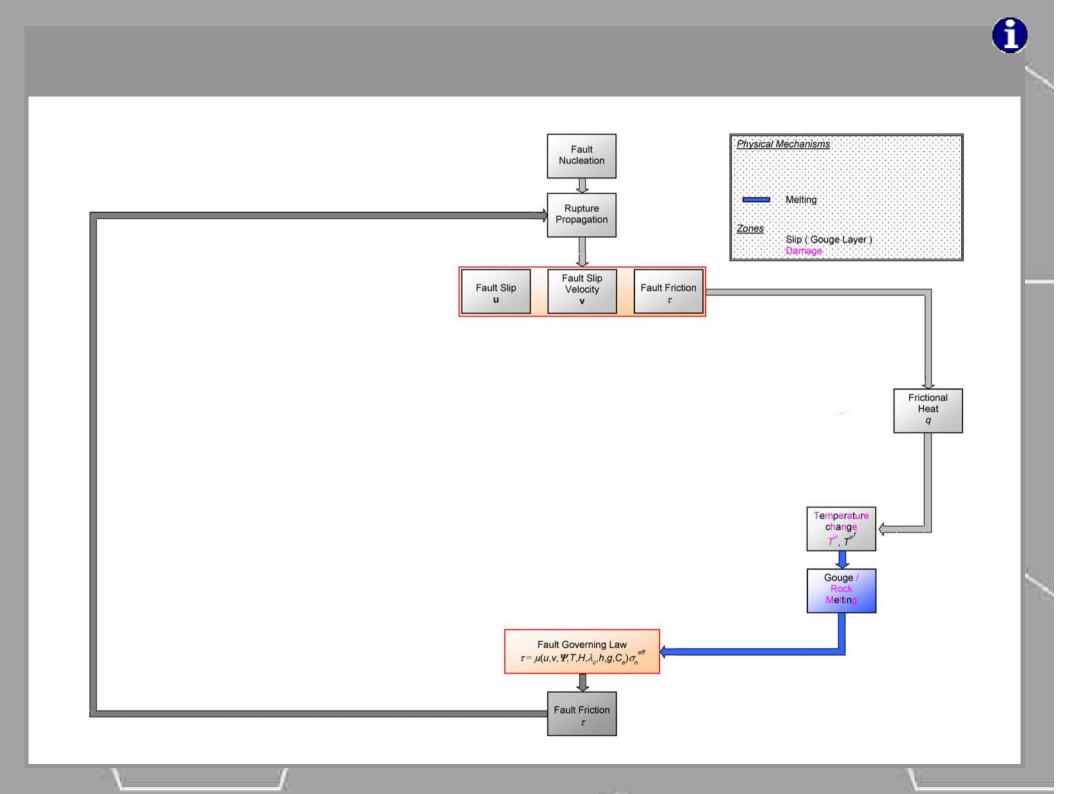


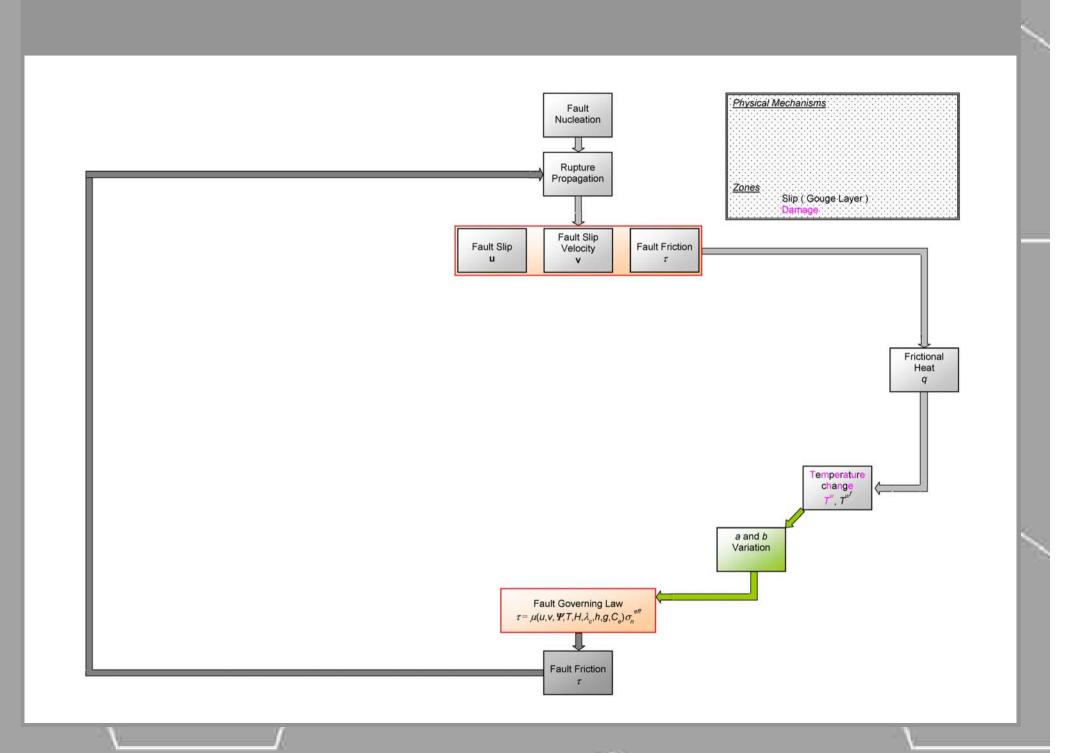


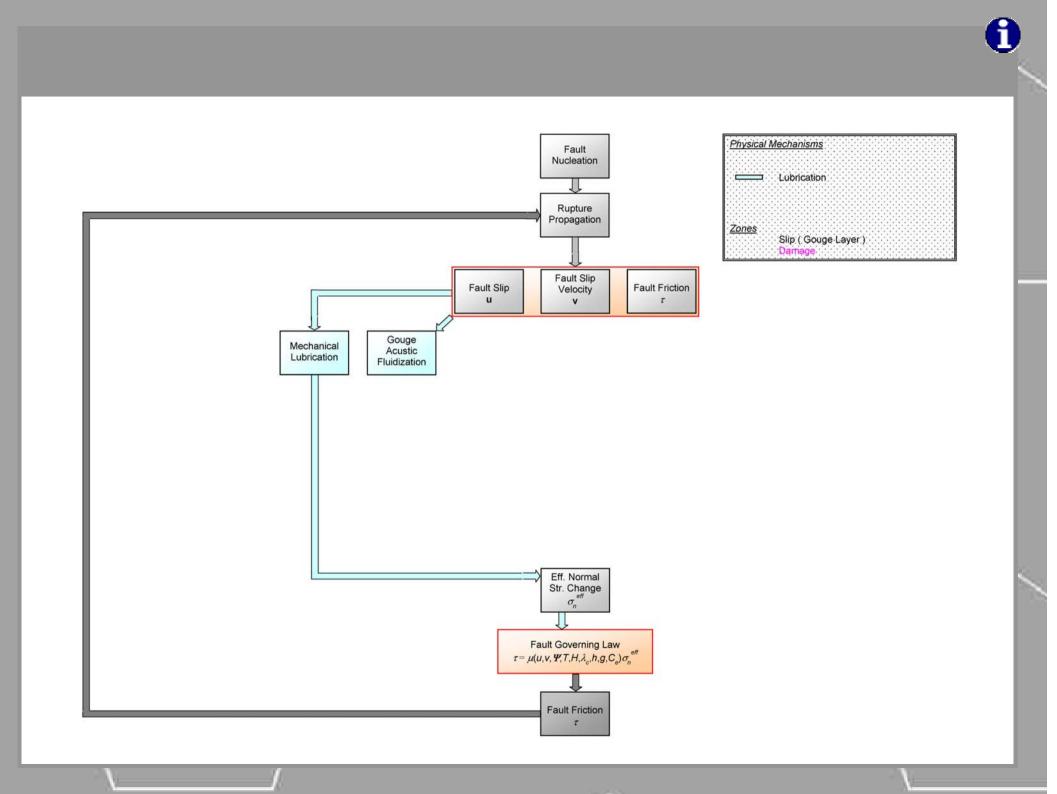


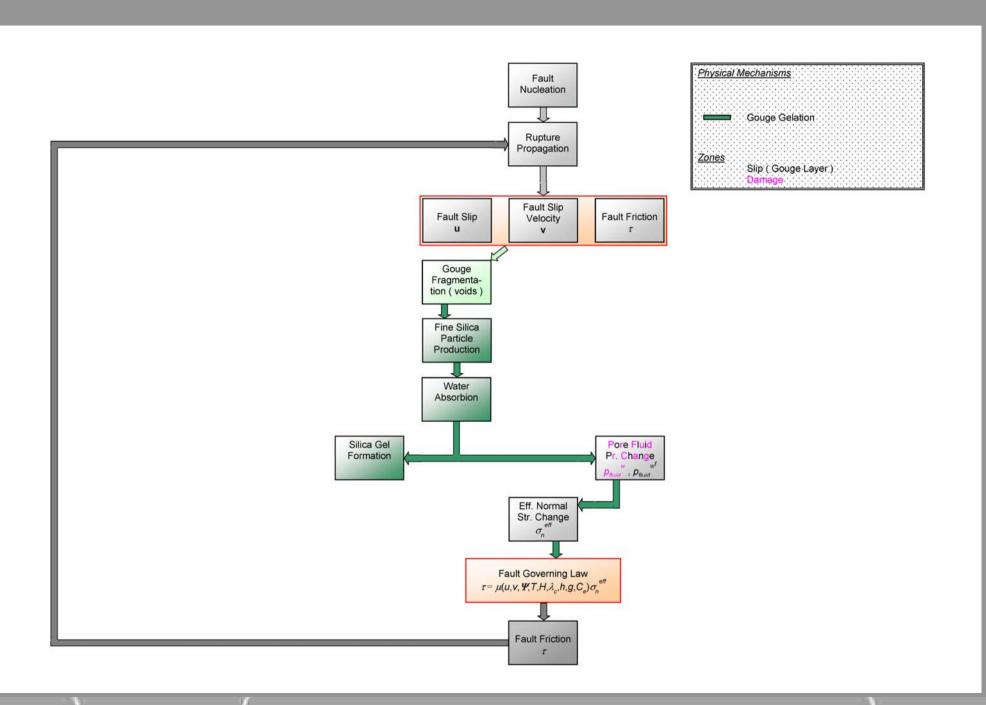


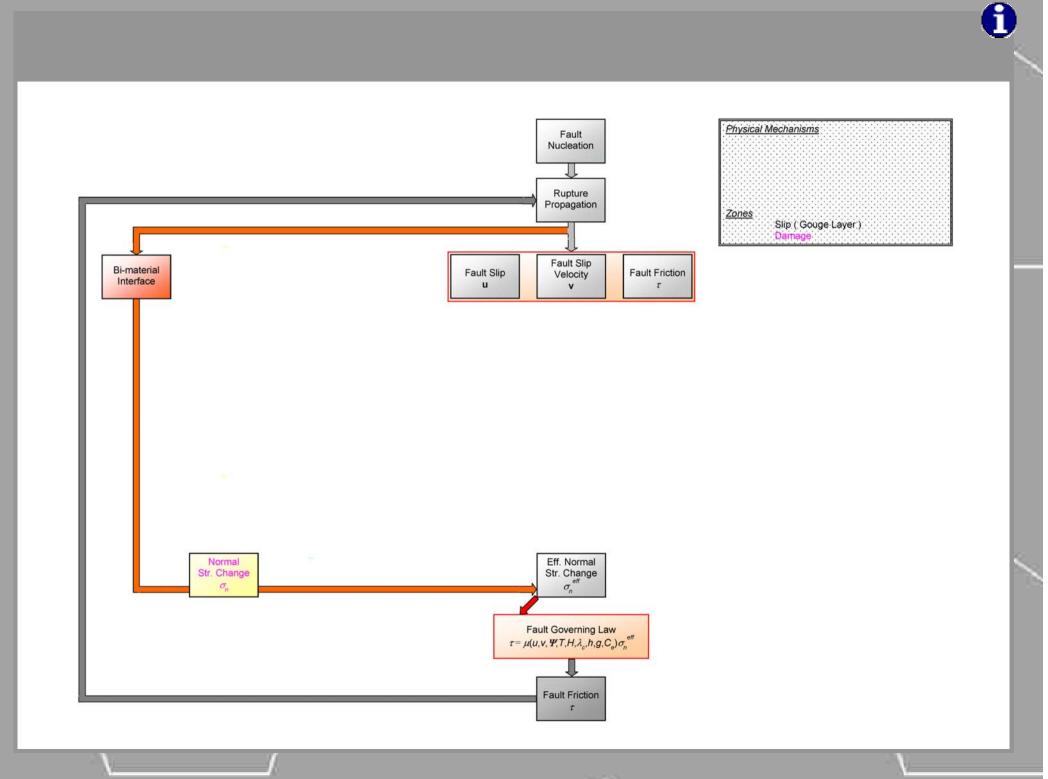


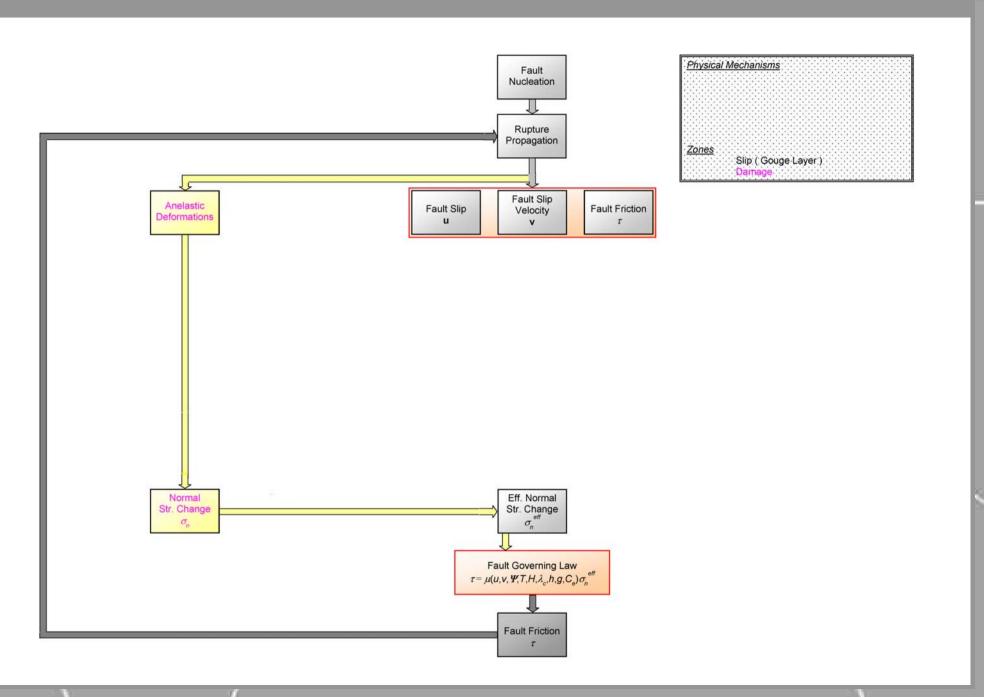




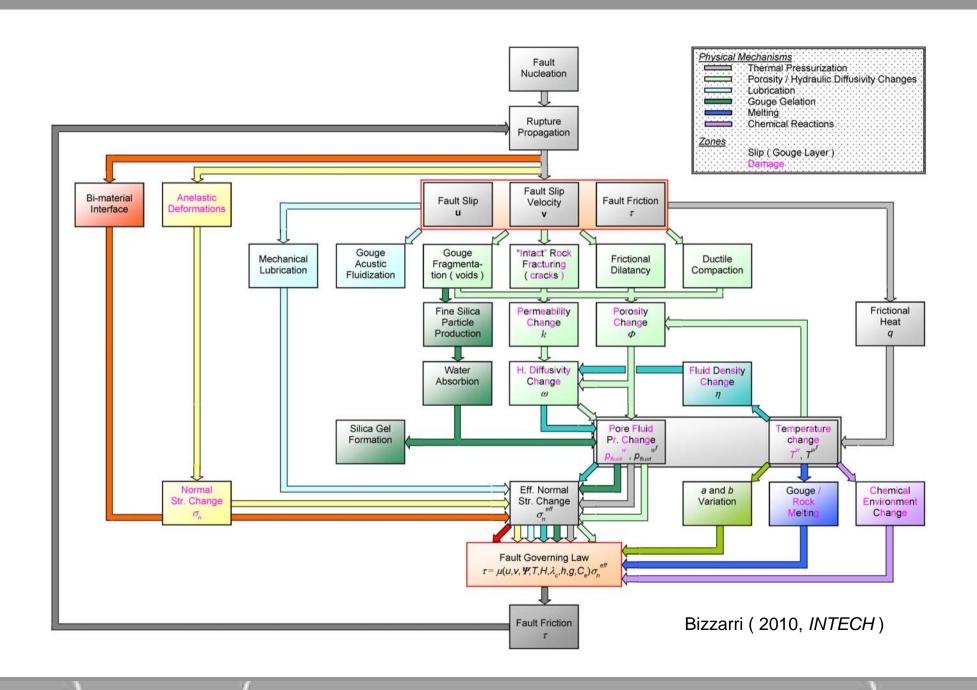








1 m



## Fracture Criteria & Constitutive Laws

### **1.** FRACTURE CRITERION

Condition that specify, at a given fault point and at a given time, if there is a rupture or not.

- It can be expressed in terms of energy, in terms of maximum frictional resistence, and so on.
- It is based on (*i*) the *Benioff (1951)* hypothesis: The fracture occours when the stress in a volume reaches the rock strength

or, analogoulsy,

(*ii*) the *Reid* (1910) statement: The fracture takes place when the stress attains a value greater than the rock can endure.

### 2. CONSTITUTIVE LAW

Analytical relation existing between the components of the stress tensor and physical observable(s), like the slip, the slip velocity, the state variable, etc..

- From a mathematical point of view it is a Fault Boundary Condition (FBC) that controls earthquake dynamics and its complexity in space and in time.
- Its simplest form consider only two frictional levels,  $\tau_u$  and  $\tau_f$ ; it accounts for stress drop (or stress realease), but the process is instantaneous: there is a singularity at crack tip.
- Cohesive zone models: Barenblatt (1959a, 1959b), Ida (1972), Andrews (1976a, 1976b). In these models the singularity is removed and the sress release occours over a breakdown zone distance  $X_b$  and in a breakdown zone time  $T_b$ .
- Friction laws (Rate and State dependent f. l.): Dieterich (1976), Ruina (1980, 1983). They accounts for fault spontaneous nucleation, re – strengthening, healing, etc..

#### CONSTITUTIVE LAW (continues)

- "The central issue is *whether* faults obey simple friction laws, and if so, what is the friction coefficient associated with fault slip " (*Scholz and Hanks, 2004*).

### CONSTITUTIVE LAW (continues)

- In full of generality we can express the constitutive ( or governing ) as:

$$\tau = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{eff}(\sigma_n, p_f)$$

#### where:



- u is the Slip (i. e. displ. disc.) modulus,
- v is the Slip Velocity modulus (its time der.),

 $\Psi = (\Psi_1, ..., \Psi_N)$  is the State Variable vector,

- T is the Temperature ( accounting for Ductility, Plastic Flow, Melting and Vaporization ),
- *H* is the Humidity,
- $\lambda_c$  is the Characteristic Length of surface (accounting for Roughness and Topography of asperity contacts),
- h is the Hardness,
- g is the Gouge ( accounting for Surface Consumption and Gouge formation ),
- $C_{e}$  is the Chemical Environment

# Strength & Constitutive Laws

1. THE STRENGTH PARAMETER

- Hystorically introduced by Das and Aki ( 1977a, 1977b) to have a quantitative extimate of the ability to fracture for a fault
- Its expression can be generalized as:

$$S = (\mu_u \sigma_n^{\text{eff}} - \tau_0) / (\tau_0 - \mu_f \sigma_n^{\text{eff}})$$

where  $\mu$  are the friction coefficient.

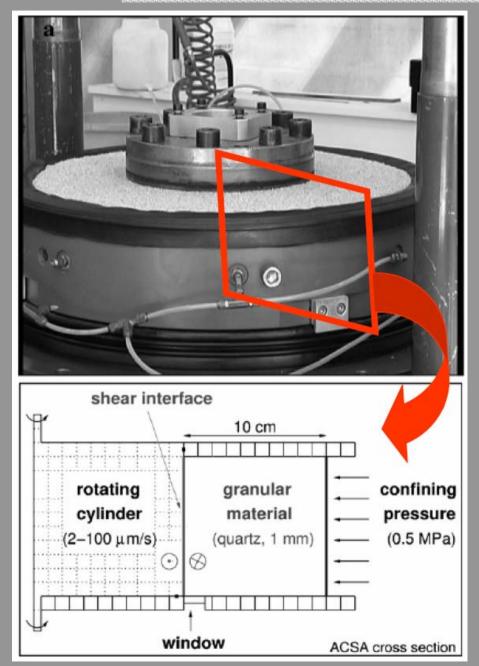
- We can also define

### 2. THE FAULT STRENGTH

 Is the parameter that quantify the Strenght in the more general case, in which a fault is described by a rhealistic friction laws

 $S^{fault} = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{eff}(\sigma_n, p_{fluid})$ 

#### Annular simple shear apparatus



 $u_{tot}$  < 50 m  $v = 1 \ \mu$ m/s - 0.1 mm/s  $\sigma_n^{eff}$  < 1 MPa

Chambon et al. ( 2006a, 2006b, *JGR*, **111**, B09308, B09309 )

#### High velocity rotary friction apparatus

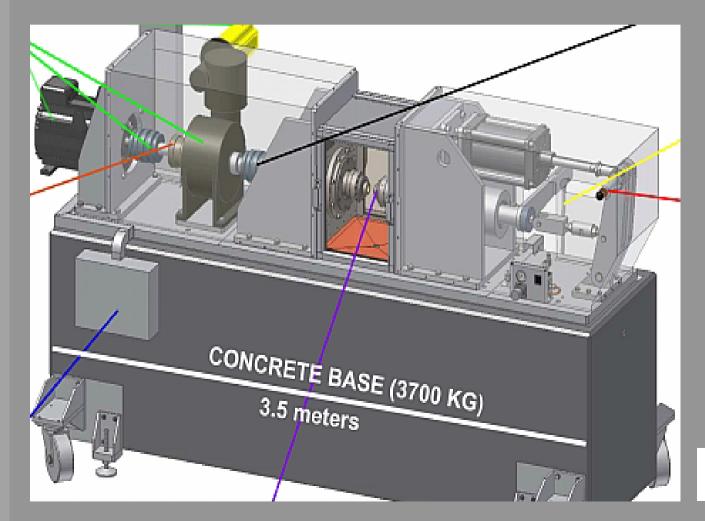


 $u_{tot}$  = infinite  $v = 0.1 \ \mu m/s - 10 \ m/s$  $\sigma_n^{eff} < 20 \ MPa$ 

Shimamoto and Tsutumi (2004, Str. Geol., **39**)

### High velocity rotary friction apparatus @ INGV

 $u_{tot}$  = infinite  $v = 1 \ \mu m/s - 9 \ m/s$  $\sigma_n^{eff} < 70 \ MPa$ 



Niemeijer et al. (2009, *AGU Fall Meeting*)

# **Time - weakening Friction Law**

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{(t - t_r)}{t_0} \right] \sigma_n^{eff} & , t - t_r < t_0 \\ \mu_f \sigma_n^{eff} & , t - t_r \ge t_0 \end{cases} \xrightarrow{\text{ilaw}} \text{II}$$

 $t_r = t_r(\xi)$  is the rupture onset time in every fault point  $\xi$  (when u > 0). <u>Andrews (1985</u>), Bizzarri et al. (2001) and other following Bizzarri's papers

 $t_0$  is the characteristic time – weakening duration.

## **Position - weakening Friction Law**

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{x}{R_0} \right] \sigma_n^{eff} & , -R_0 < x < 0 \\ \mu_f \sigma_n^{eff} & , -L < x < -R_0 \end{cases}$$

 $\mathbf{PW}$ 

x is the position on the fault (extending up to -L).

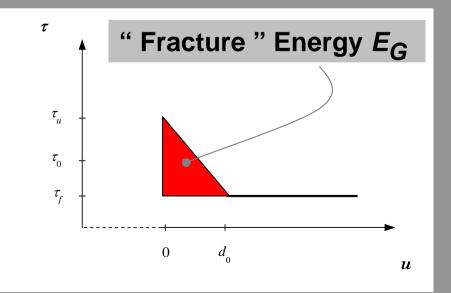
Palmer and Rice (1973)

 $R_0$  is the characteristic position – weakening distance.



### 1. LINEAR SLIP - WEAKEING LAW

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{u}{d_0} \right] \sigma_n^{eff} & , u < d_0 \\ \mu_f \sigma_n^{eff} & , u \ge d_0 \end{cases} \quad \begin{array}{l} \text{ilaw} = 21 \\ \text{sw} \\ \text{sw} \end{cases}$$



Barenblatt (1959a, 1959b), <u>Ida</u> (<u>1972</u>), Andrews (1976a, 1976b), and many authors thereinafter

 $d_0$  is the characteristic slip – weakening distance

#### ilaw = 22

IW

### 2. NON – LINEAR SLIP – WEAKEING LAW

$$\tau = \begin{cases} \left[ \mu_u - \frac{\mu_u - \mu_f}{d_0} \left( u - \frac{(1 - p_{IW})d_0}{2\pi} \sin\left(\frac{2\pi u}{d_0}\right) \right) \right] \sigma_n^{eff} &, u < d_0 \\ \mu_f \sigma_n^{eff} &, u \ge d_0 \end{cases} \end{cases}$$

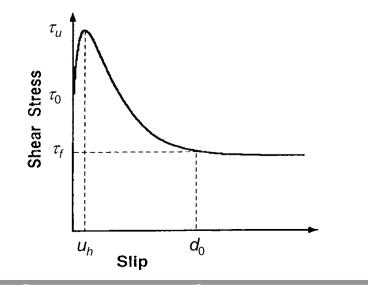
Ionescu and Campillo (1999)

#### 3. NON LINEAR SLIP – WEAKEING LAW WITH SLIP – HARDENING

$$\tau = \left\{ \begin{bmatrix} \left( \frac{\tau_0}{\sigma_n^{eff}} - \mu_f \right) \left( 1 + \alpha_{OY} \ln \left( 1 + \frac{u}{\beta_{OY}} \right) \right) \end{bmatrix} e^{-\frac{u}{d_0}} + \mu_f \right\} \sigma_n^{eff}$$

$$\mathbf{u}_h : \frac{\mathrm{d}\tau}{\mathrm{d}u}\Big|_{u_h} = 0; \qquad \begin{cases} u_h = rd_0 \quad (\mathrm{e.\,g.} \ r = 0.1) \\ \tau(u_h) = \tau_u \end{cases}$$

$$\mathbf{W}_h : \frac{\mathrm{d}\tau}{\mathrm{d}u}\Big|_{u_h} = 0;$$



<u>Ohnaka and Yamashita (1989)</u> and the following papers by Ohnaka and coworkers

 $u_h$  is associated with the preparatory phase of the imminent macroscopic failure in the cohesive zone. It accounts for micro–cracking

#### 4. NON LINEAR SLIP - WEAKEING LAW WITH EXPONENTIAL DECAY

$$\tau = \left[ \left( \mu_u - \mu_f \right) e^{-\frac{u}{d_0}} + \mu_f \right] \sigma_n^{eff}$$

ilaw = 24

EW

### 5. POWER LAW SLIP - WEAKEING

$$\tau = \left\{ \mu_u - \left(\mu_u - \mu_f\right) \left[ \left(\frac{p_{PW}}{p_{PW} + 1}\right) \frac{u}{d_0} \right]^{p_{PW}} \right\} \sigma_n^{eff}$$

ilaw = 25

PW

# **Rate - Dependent Friction Law**

$$\tau = \frac{\upsilon_*}{\upsilon + \upsilon_*} \,\mu_u \sigma_n^{\text{eff}}$$

Burridge and Knopoff ( 1967 ), <u>Carlson and Langer ( 1989 )</u>, Madariaga and Cochard ( 1994 ), Cochard and Madariaga ( 1994 )

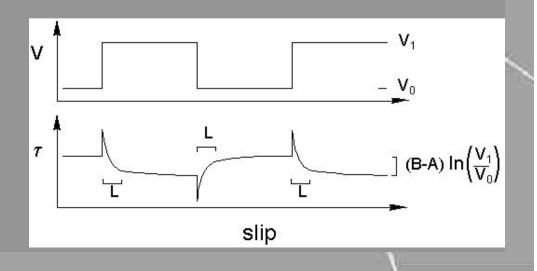
# Rate - and State - Dependent Friction Laws

**1.** DIETERICH IN REDUCED FORMULATION

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_* - \alpha \ln \left( \frac{v_*}{v} \bullet \right) + b \ln \left( \frac{\Psi v_*}{L} \bullet \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} \end{cases} \text{ In } \\ \end{bmatrix} \sigma_n^{eff} \\ DR$$

However, while in velocity stepping experiments the traction response following the velocity variation is directly controlled by the parameter *L*, its effects are much less evident during the dynamic rupture propagation. Bizzarri and Cocco (2005)

#### Response to an abrupt jump in load



### 2. RUINA – DIETERICH

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_{*} - a \ln \left( \frac{v_{*}}{v} \right) + b \ln \left( \frac{\Psi v_{*}}{L} \right) \right] \sigma_{n}^{eff} \\ \frac{d}{dt} \Psi = - \frac{\Psi v}{L} \ln \left( \frac{\Psi v}{L} \right) \end{cases} \text{ RD} \end{cases}$$

<u>Ruina (1980, 1983)</u>, Beeler et al. (1984), Roy and Marone (1996)

### 3. DIETERICH – RUINA WITH VARYING NORMAL STR.

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_{*} - a \ln \left( \frac{v_{*}}{v} \right) + b \ln \left( \frac{\Psi v_{*}}{L} \right) \right] \sigma_{n}^{eff} & \text{ilaw} = 31 \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} - \left( \frac{\alpha_{LD} \Psi}{b \sigma_{n}^{eff}} \right) \frac{d}{dt} \sigma_{n}^{eff} & \text{DR} \end{cases} \\ \end{cases}$$

<u>Linker and Dieterich (1992)</u>, Dieterich and Linker (1992), Bizzarri and Cocco (2006b, 2006c)

#### 4. RUINA – DIETERICH WITH VARYING NORMAL STR.

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_{*} - a \ln \left( \frac{v_{*}}{v} \right) + b \ln \left( \frac{\Psi v_{*}}{L} \right) \right] \sigma_{n}^{eff} & \text{ilaw} = 32 \\ \frac{d}{dt} \Psi = -\frac{\Psi v}{L} \ln \left( \frac{\Psi v}{L} \right) - \left( \frac{\alpha_{LD} \Psi}{b \sigma_{n}^{eff}} \right) \frac{d}{dt} \sigma_{n}^{eff} & \text{RD} \end{cases}$$

<u>Linker and Dieterich (1992)</u>, Bizzarri and Cocco (2006b, 2006c)

### 5. DIETERICH IN REDUCED FORM REGULARIZED

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_{*} - a \ln \left( \frac{v + v_{*}}{v + v_{*}} \right) + b \ln \left( \frac{\Psi(v - v)}{L} + 1 \right) \right] \sigma_{n}^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi(v + v)}{L} \end{cases} \\ \end{bmatrix} \text{DE}$$

 $v_r$  is a regularization fault slip velocity

<u>Perrin et al. ( 1995 )</u>, Cocco et al. ( 2004 )

### 6. RUINA REGULARIZED

$$\tau = \left[ \begin{array}{c} \mu_{*} - \alpha \ln \left( \frac{v_{*} - v_{r}}{v_{+} + v_{r}} \right) + \frac{\Psi}{\sigma_{n}^{eff}} \right] \sigma_{n}^{eff} \\ \frac{d}{dt} \Psi = -\frac{v_{+} v_{r}}{L} \left( \Psi + b \ln \left( \frac{v_{*} - v_{r}}{v_{*} - v_{r}} \right) \right) \end{array} \right]$$
 ilaw = 34  
RE

 $v_r$  is a regularization fault slip velocity

Bizzarri (2002, unpublished work)

### 7. DIETERICH IN REDUCED FORM WITH HEALING

$$\begin{cases} \tau = \left[ \begin{array}{c} \mu_{*} - \alpha \ln \left( \frac{v_{*}}{v} + 1 \right) + b \ln \left( \frac{\Psi v_{*}}{L} + 1 \right) \right] \sigma_{n}^{eff} \\ \frac{d}{dt} \Psi = \frac{\gamma_{fh} - \Psi}{L} - \frac{\Psi v}{L} \end{cases} \text{ DH} \end{cases}$$

 $\gamma_{fh} = 1 \text{ s}$ 

 $t_{fh}$  is the time for healing (slip duration)

Evolution law proposed by <u>Nielsen et</u> <u>al. (2000)</u> and by Nielsen and Carlson (2000). Used in this form by Cocco et al. (2004)

### 9. PRAKASH – CLIFTON

$$\tau = \left[ \begin{array}{c} \mu_{*} - \alpha \ln \left( \frac{v_{*}}{v} \right) + b \ln \left( \frac{\Psi v_{*}}{L} \right) \right] \left( \frac{\mathrm{d}}{\mathrm{d}t} \Psi_{1} + \frac{\mathrm{d}}{\mathrm{d}t} \Psi_{2} \right) \\ \frac{\mathrm{d}}{\mathrm{d}t} \Psi = 1 - \frac{\Psi v}{L} \\ \frac{\mathrm{d}}{\mathrm{d}t} \Psi_{1} = -\frac{v}{L_{1}} \left( \Psi_{1} - \alpha_{PC_{1}} \sigma_{n}^{eff} \right) \\ \frac{\mathrm{d}}{\mathrm{d}t} \Psi_{2} = -\frac{v}{L_{2}} \left( \Psi_{2} - \alpha_{PC_{2}} \sigma_{n}^{eff} \right) \end{array} \right]$$

$$I = 1 - \frac{\Psi v}{L_{1}} \left( \Psi_{1} - \alpha_{PC_{1}} \sigma_{n}^{eff} \right)$$

$$I = 1 - \frac{\Psi v}{L_{1}} \left( \Psi_{1} - \alpha_{PC_{2}} \sigma_{n}^{eff} \right)$$

 $\Psi_1$  and  $\Psi_2$  are additional state variables accountinf for the coupling with effective normal stress. The formulation of friction law is not based on the Amonton – Coulamb law.

Coupling with effective normal stress proposed by <u>Prakash and Clifton</u> (1993) and Prakash (1998). Used in this form by Bizzarri (2005, unpublished work)

$$\begin{split} & \Delta \mu \text{ is an initial artificial stress drop} \\ & \Psi_1 \equiv \Psi_0 \; (u - u_1) / (d_1 - u_1) \\ & U_1 \equiv - \; d_1 \; (\mu_{sp} - \mu_u + \Delta \mu) / (\mu_u - \Delta \mu) \\ & d_0 \text{ and } d_1 \text{ are characteristic lengths} \\ & \mu_{sp} \; = \; 0 \; \Rightarrow \; \text{linear SW with } d_1 \text{ as characteristic length} \end{split}$$

Cochard and Madariaga (1994)

## -Free Volume Friction law

$$\begin{cases} \tau = \sigma_d \operatorname{Arcsinh} \left( \frac{\upsilon}{\upsilon_*} \frac{\mathrm{e}^{f_* + \frac{\chi_s + \chi_h}{\chi}}}{1 - m_0} \right) \\ \frac{\mathrm{d}}{\mathrm{d}t} \chi = -R_c \, \mathrm{e}^{-\frac{\chi_c}{\chi}} + \alpha_{FV} \tau \upsilon \\ m_0 = \begin{cases} 1 & , \tau \leq \tau_0 \mathrm{e}^{\frac{\chi_h}{\chi}} \\ \frac{\tau_0}{\tau} \, \mathrm{e}^{\frac{\chi_h}{\chi}} & , \tau > \tau_0 \mathrm{e}^{\frac{\chi_h}{\chi}} \end{cases} \end{cases}$$

 $\begin{array}{ll} \chi \equiv \varPhi - \varPhi_0 & \mbox{free volume variable} \\ \chi_s & \mbox{reference value of } \chi \mbox{ for shearing} \\ \chi_h & \mbox{FV value required to create a Shear Transformation Zone (STZ)} \\ \chi_c & \mbox{FV value for compaction} \\ R_c & \mbox{rate of compaction} \\ \alpha_{FV} & \mbox{scaled dilatancy coefficient} \end{array}$ 

Falk and Langer ( 1998, 2000 ); Lemaitre ( 2002 ); <u>Daub and Carlson</u> (2008 )

## How to relate relevant quantities to contitutive parameters

### **Dynamic Parameters** $E_R$ : radiated energy ruputre speed, $V \approx 0.8$ to 0.9 $\beta$ $E_{NR} = E_F + E_G + \dots$ : non-radiated energy $E_F$ : friction (heat), $E_G$ : fracture energy of : friction Stress 50 $\Delta \sigma_s = \sigma_0 - \sigma_1$ : static stress drop $\Delta \sigma_d = \sigma_0 - \overline{\sigma}_f$ : dynamic stress drop 5 end begin Slip (xS)

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### Support Slides: Parameters, Notes, etc.

To not be displayed directly. Referenced above.

#### **Thermal pressurization:**

Sibson (1973); Lachenbruch (1980); Mase and Smith (1985, 1987); Andrews (2002); Bizzarri and Cocco (2006b, 2006c).

Morrow et al. (1984) show that gouge contains water

#### Gouge behaviour:

Marone et al. (1990); Marone and Kilgore (1993); Mair and Marone (1999); Mair et al. (2002)

#### Frictional melting:

Jeffreys (1942); McKenzie and Brune (1972); Richards (1977); Sibson (1977); Cardwell et al. (1978); Allen (1979)



Pseudo tachylyte: Fault vein (*Sibson*, 1975)

#### **Mechanical Iubrication:**

Spray (1993); Brodsky and Kanamori (2001); Kanamori and Brodsky (2001)

#### **Acustic fluidization:**

Melosh (1979, 1996)

#### **Gouge gelation:**

Goldbsy and Tullis (2002); Di Toro et al. (2004)

#### Bi – material interface:

Andrews and Ben – Zion (1997); Harris and Day (1997); Andrews and Harris (2005)



MTL: Fractured mylonite, cataclasite and gouge

#### Humidity effects:

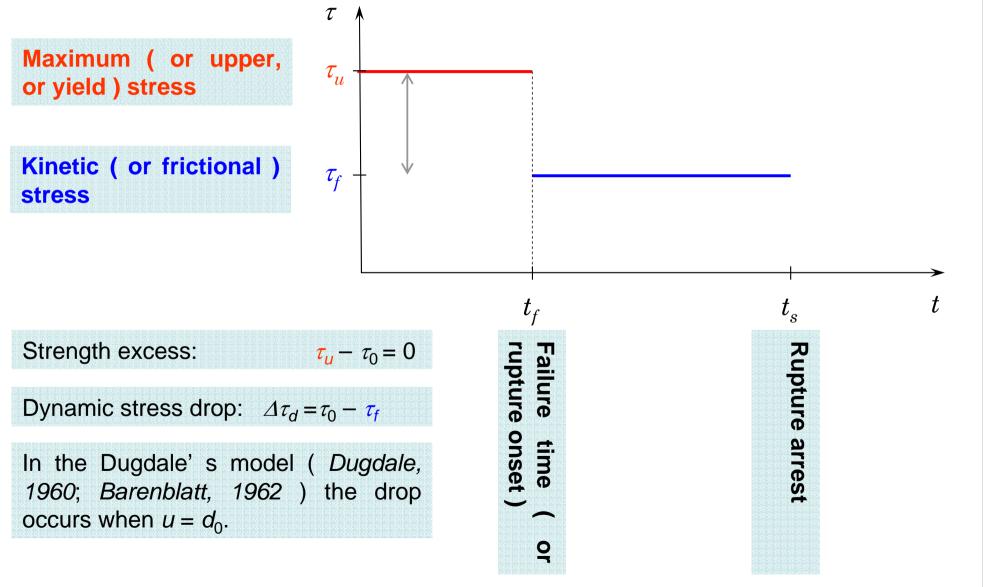
Dieterich and Conrad (1984)

#### **Characteristic length of surface effects:**

Ohnaka and Shen (1999); Ohnaka (2003)

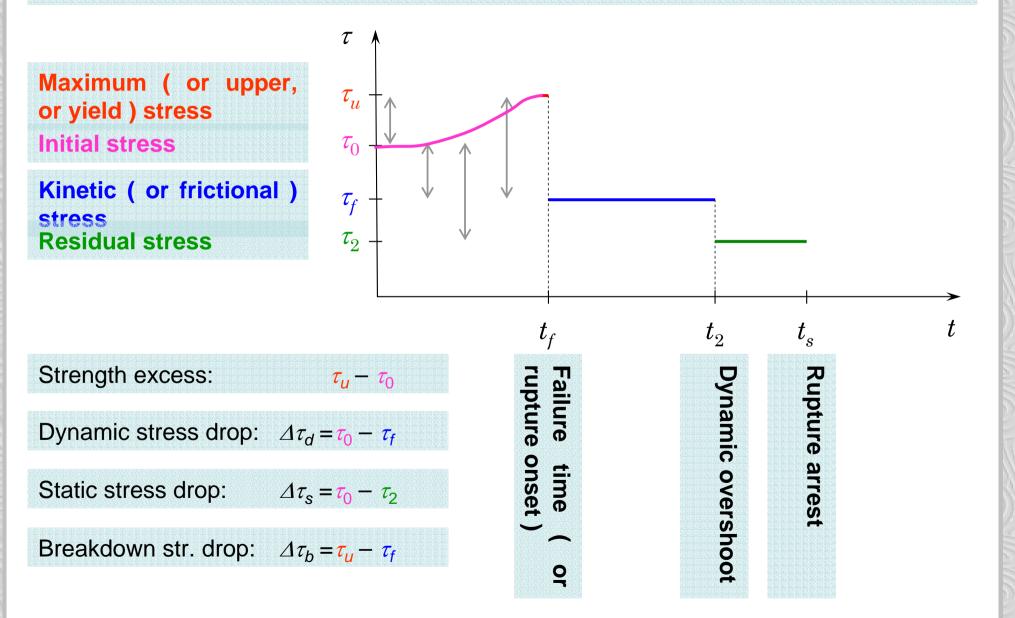
### <mark>elebom noitoiri teelqmi2</mark>





## <mark>elebom noitoiri teelqmi2</mark>

At a particular fault point  $\xi$  (following Savage and Wood, 1971; Scholz, 1990)



Savage and Wood (1971) also define:

Mean stress:  $<\tau>=\frac{1}{2}(\tau_{\mu}+\tau_{2})$ 

Seismic efficiency:  $\eta = E_s/E$ , where:  $E_s$  is the seismic energy

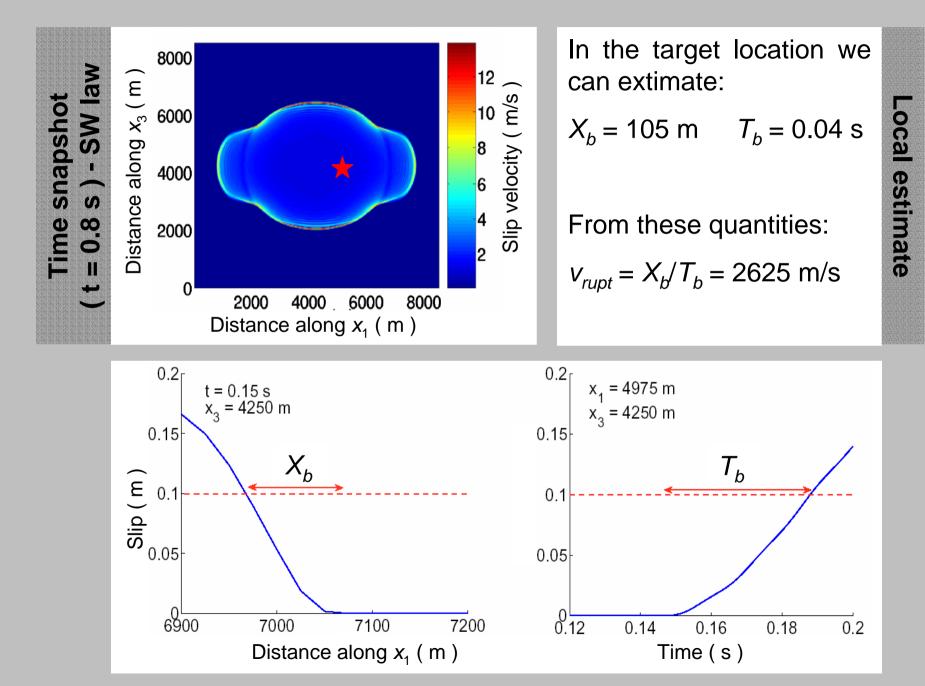
*E* is the total available energy

Apparent stress:  $\tau_a = \eta < \tau >$ 

Direct observation of the absolute stress near an earthquake is not feasible, but it is possible (*Wyss and Brune, 1968*) calculate  $\tau_a$  and stress drop from physical observables.



### <u>Lys copesive zobe</u>



## Slip - hardening effect

\* The slip – hardening (SH) phenomenon has been also found in seismological inversion studies (e.g. Quin, 1990; Miyatake, 1992; Mikumo and Miyatake, 1993; Beroza and Mikumo, 1996; Ide, 1997; Bouchon, 1997).

# Interpretation of the state variable 0