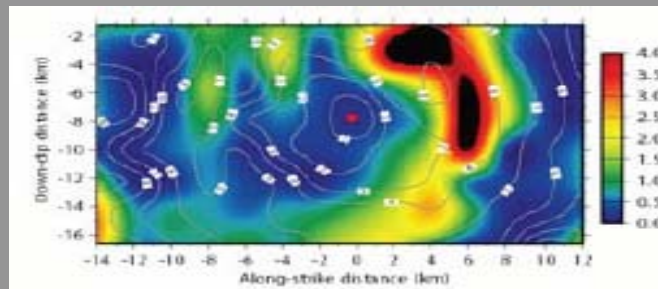




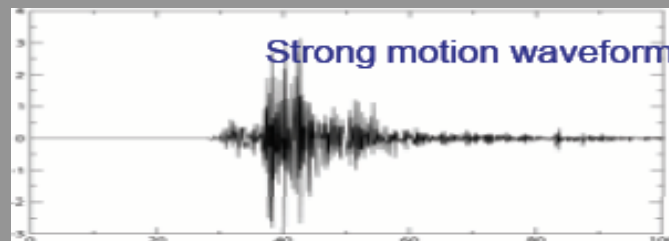
**Fault governing laws
(constitutive equations)**

Seismologists need traction

- ✓ To apply fracture mechanics on mathematical planes representing the fault surfaces;
- ✓ To numerically simulate the spontaneous rupture nucleation, propagation, healing and arrest in dynamic earthquake models;

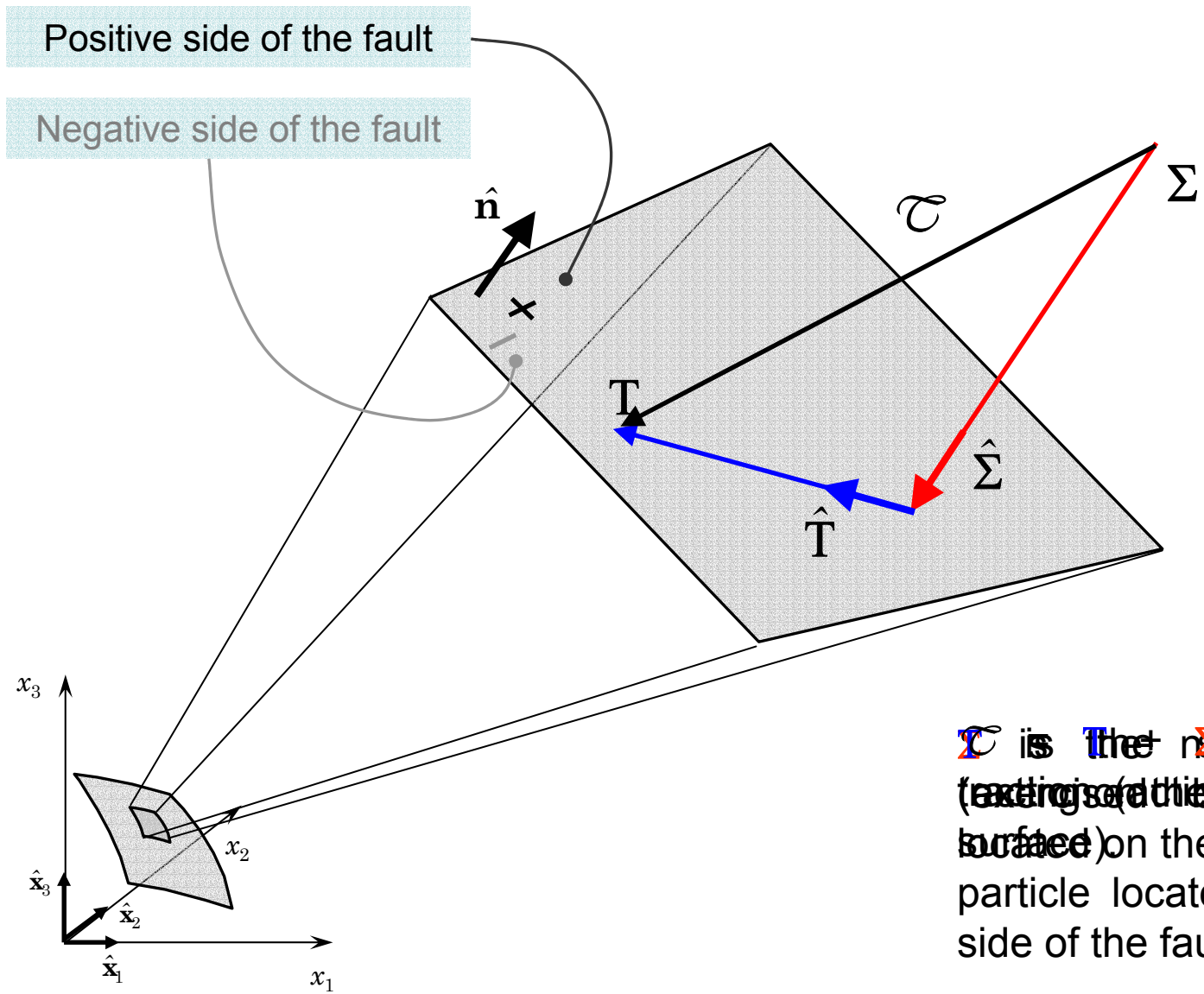


- ✓ To model seismic wave propagation in the surrounding medium;



- ✓ To predict ground shaking.

Notations and symbols



\mathcal{T} is the traction (acting on the surface) on the +ve side on a particle located on the -ve side of the fault surface)

\hat{S} is the shear (acting on the surface) on the +ve side on a particle located on the -ve side of the fault surface)

$$\mathcal{T}^{(\hat{n})} = \mathbf{T}^{(\hat{n})} + \Sigma^{(\hat{n})}$$

total traction (acting on the fault surface).

$$\mathcal{T}_j^{(\hat{n})} = n_i \sigma_{ij}^{eff}$$

Cauchy's formula, where $\mathcal{T}^{(\hat{n})} = (\mathcal{T}_1^{(\hat{n})}, \mathcal{T}_2^{(\hat{n})}, \mathcal{T}_3^{(\hat{n})})$,

$\mathbf{n} = (n_1, n_2, n_3)$ and

$$\sigma_{ij}^{eff} = \sigma_{ij} + p_{fluid} \delta_{ij} = \begin{bmatrix} -\sigma_{n_1}^{eff} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & -\sigma_{n_2}^{eff} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & -\sigma_{n_3}^{eff} \end{bmatrix}$$

where: $\sigma_{n_i}^{eff} = \sigma_{n_i} - p_{fluid} = -\sigma_{ii} - p_{fluid}$ and stresses are assumed to be negative for compression

$$T_j^{(\hat{n})} = n_i \sigma_{ij}^{eff} - n_j (n_i \sigma_{ik}^{eff} n_k)$$

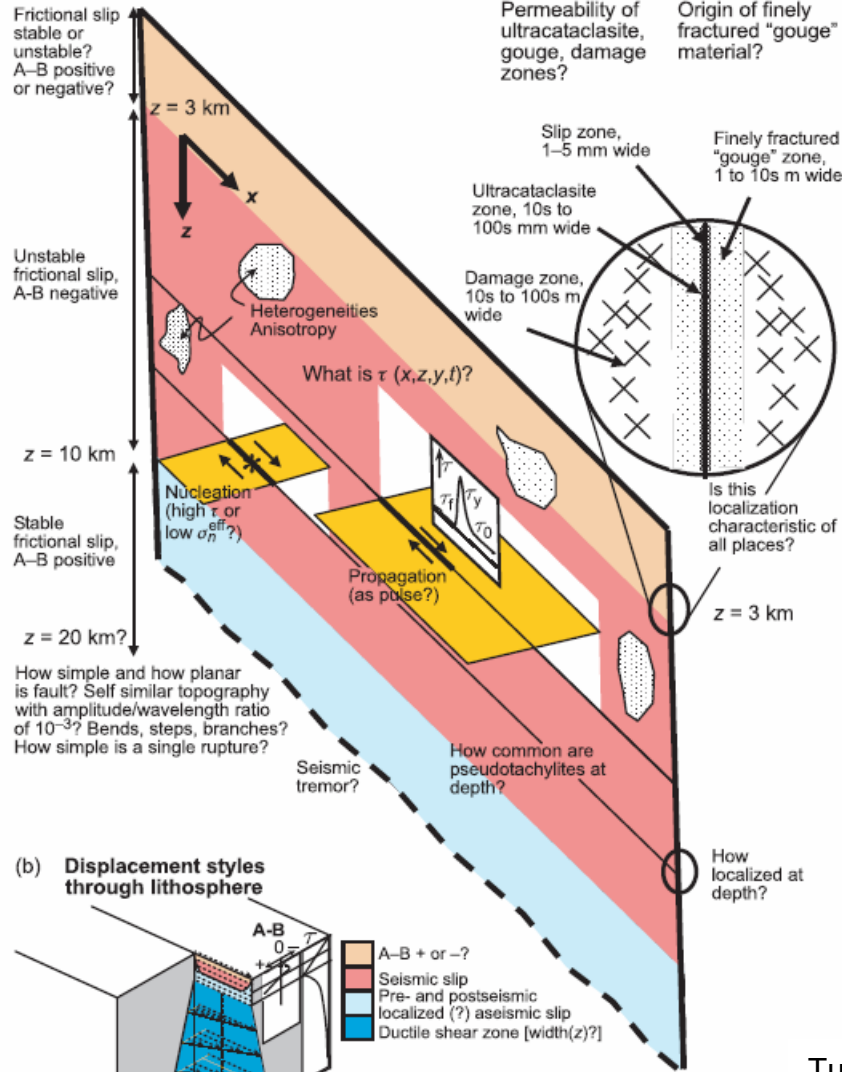
shear traction

$$\Sigma_j^{(\hat{n})} = n_j (n_i \sigma_{ik}^{eff} n_k)$$

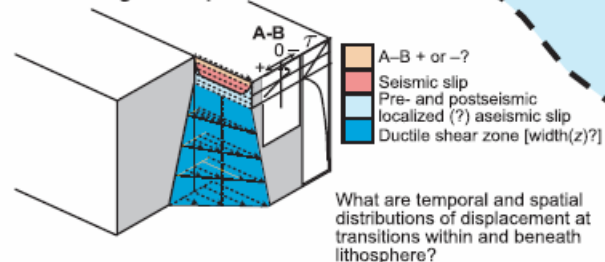
normal traction

Fault models

(a) Seismogenic part of fault



(b) Displacement styles through lithosphere



Tullis et al. (2007, MIT Press)

Internal Structure of Principal Faults of the North Branch San Gabriel Fault

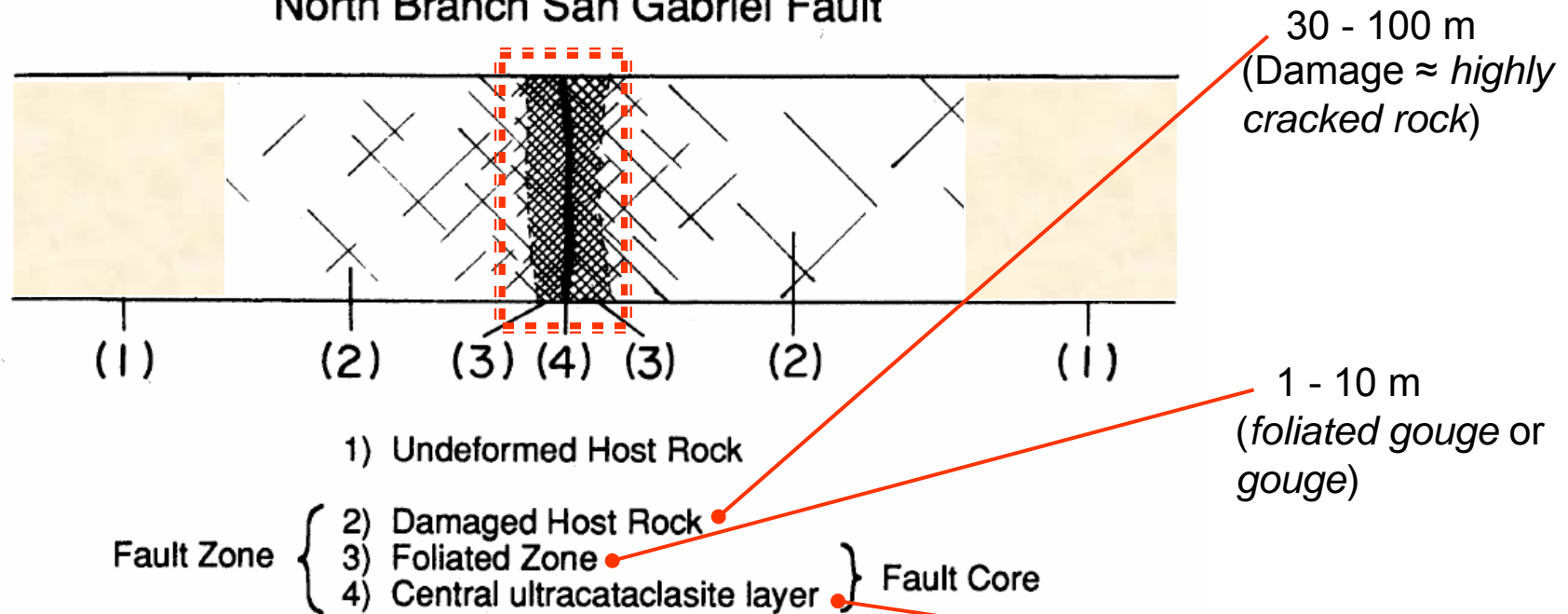
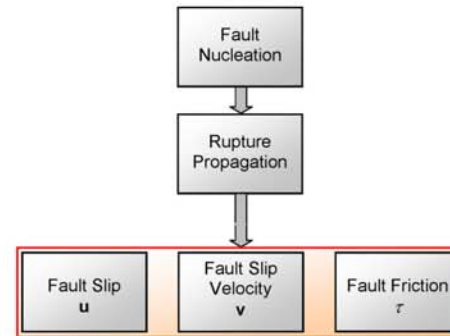
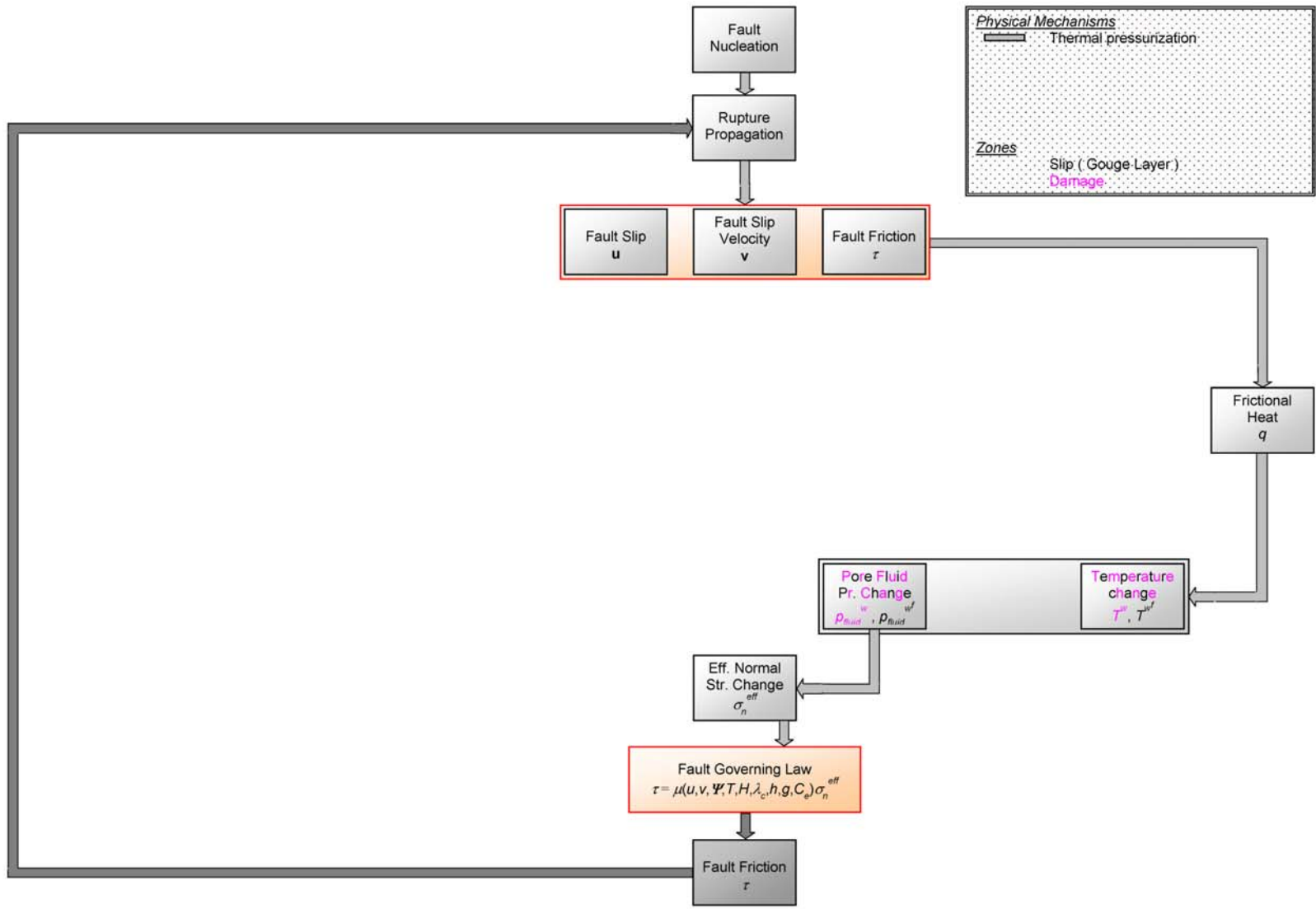


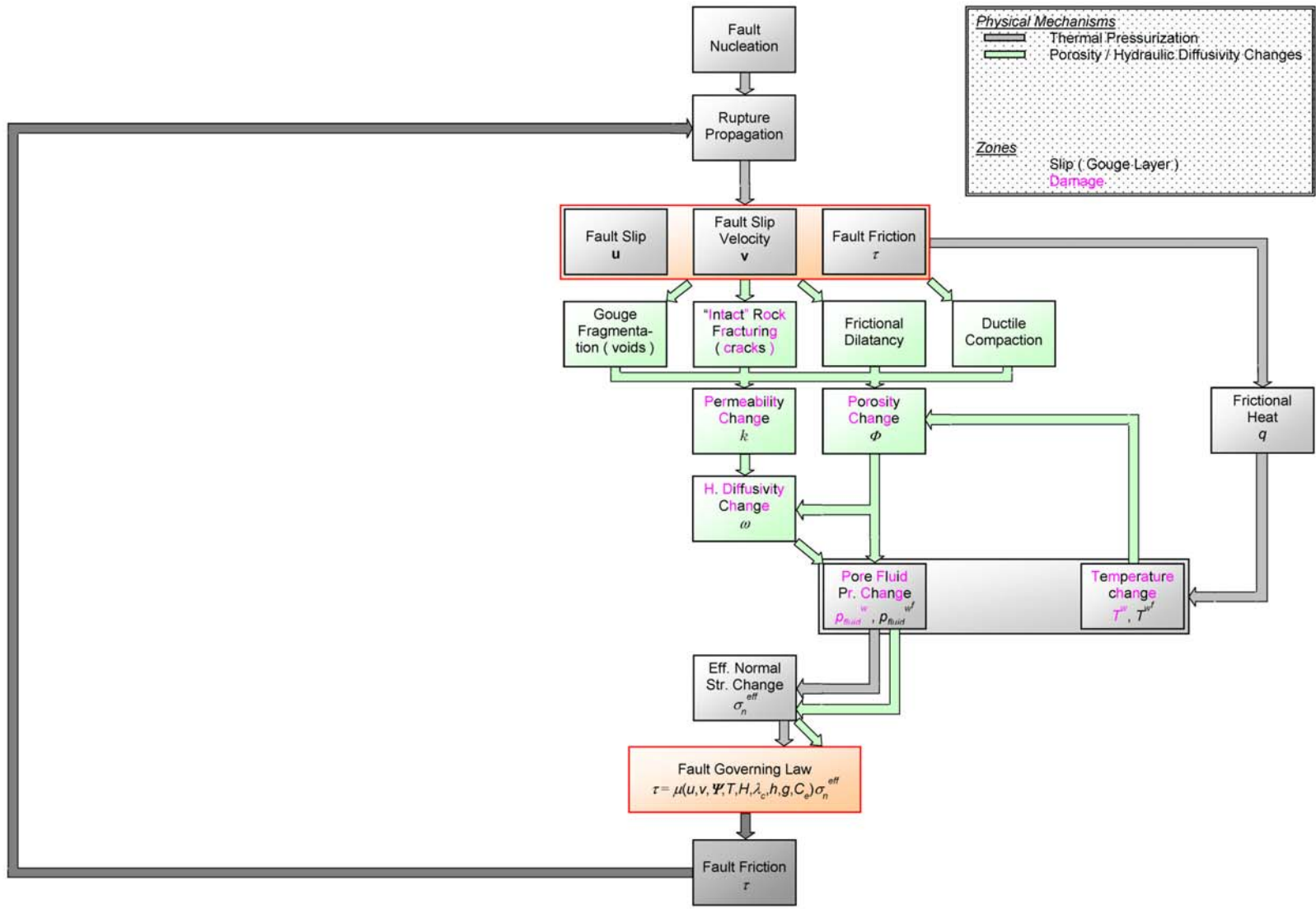
Fig. 2. Schematic section across the North Branch San Gabriel fault zone illustrating position of the structural zones of the fault. The diagram is not to scale.

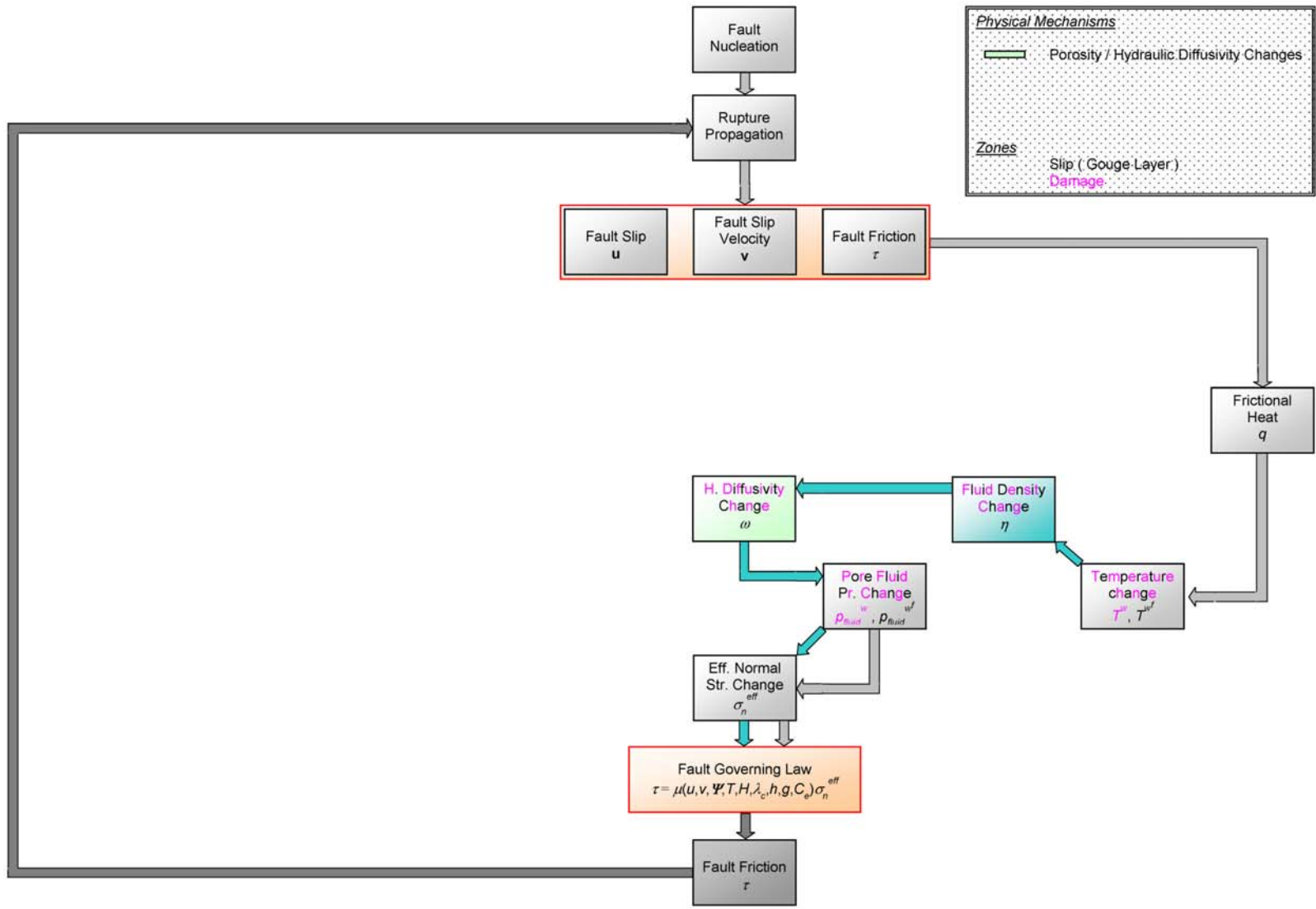
Chester, Evans and Biegel, *J. Geoph. Res.*, 1993
 Sibson, BSSA, 2003
 Chester and Chester, SSA, SCEC meetings 2004

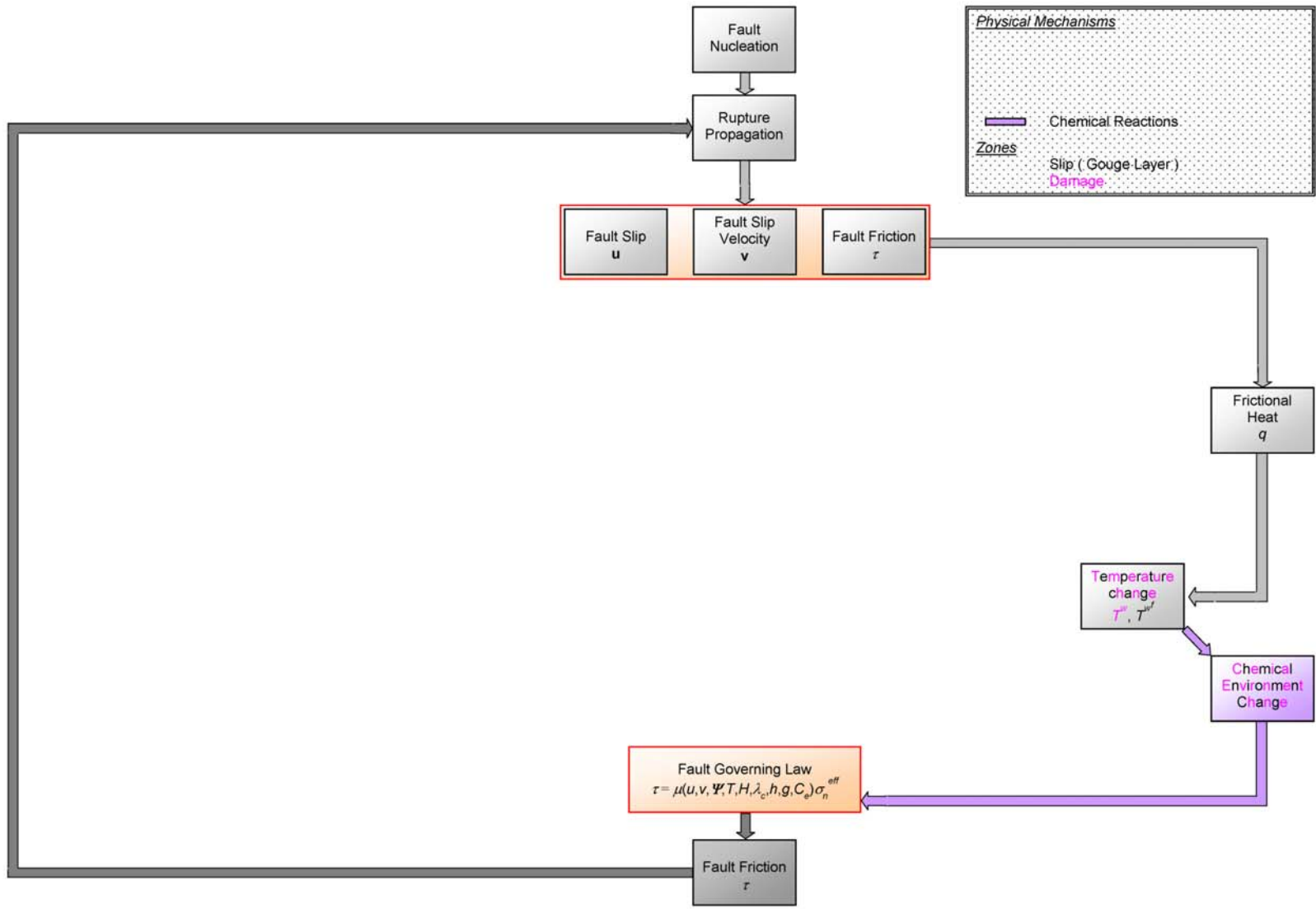
Physical Phenomena in Faulting

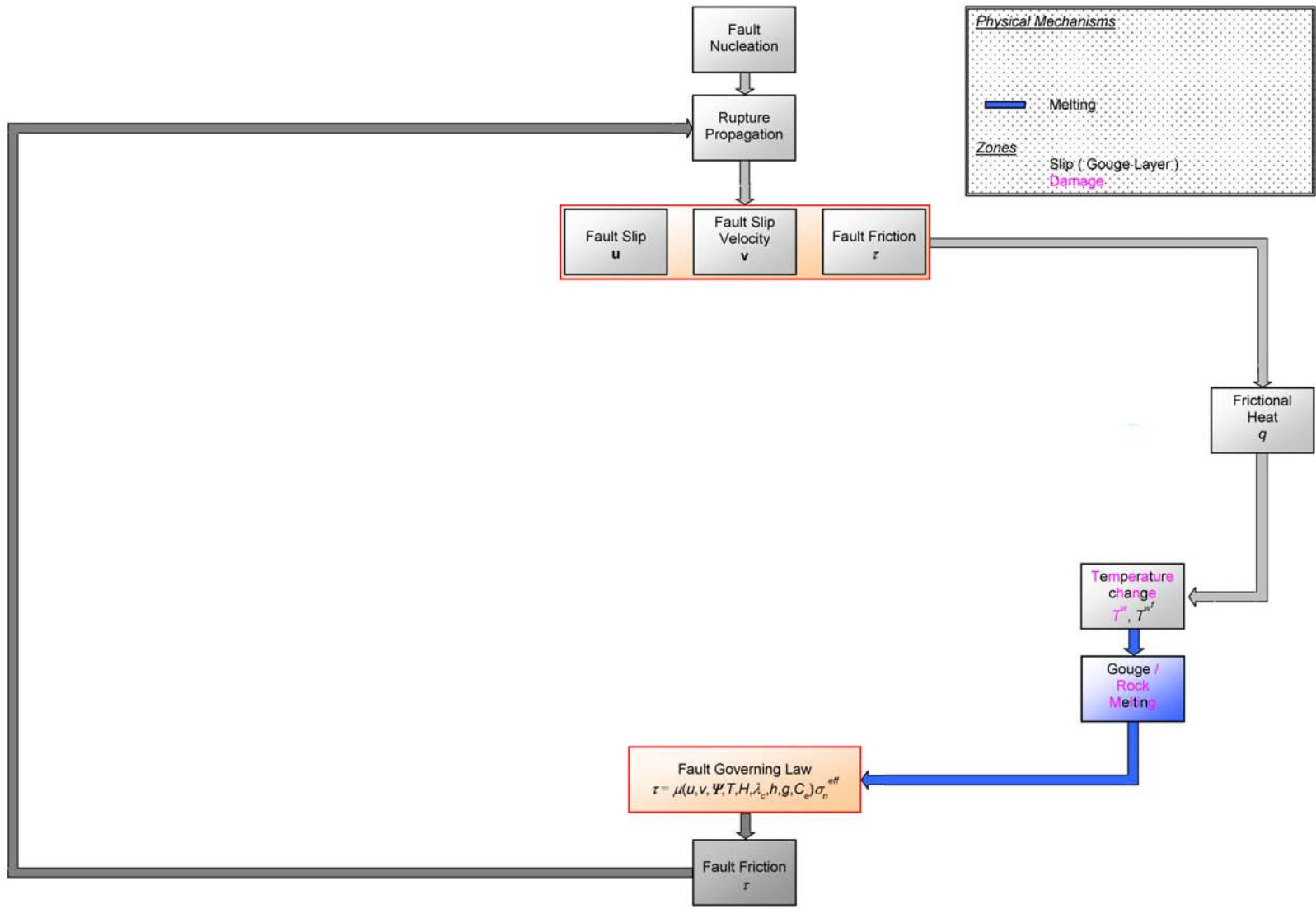


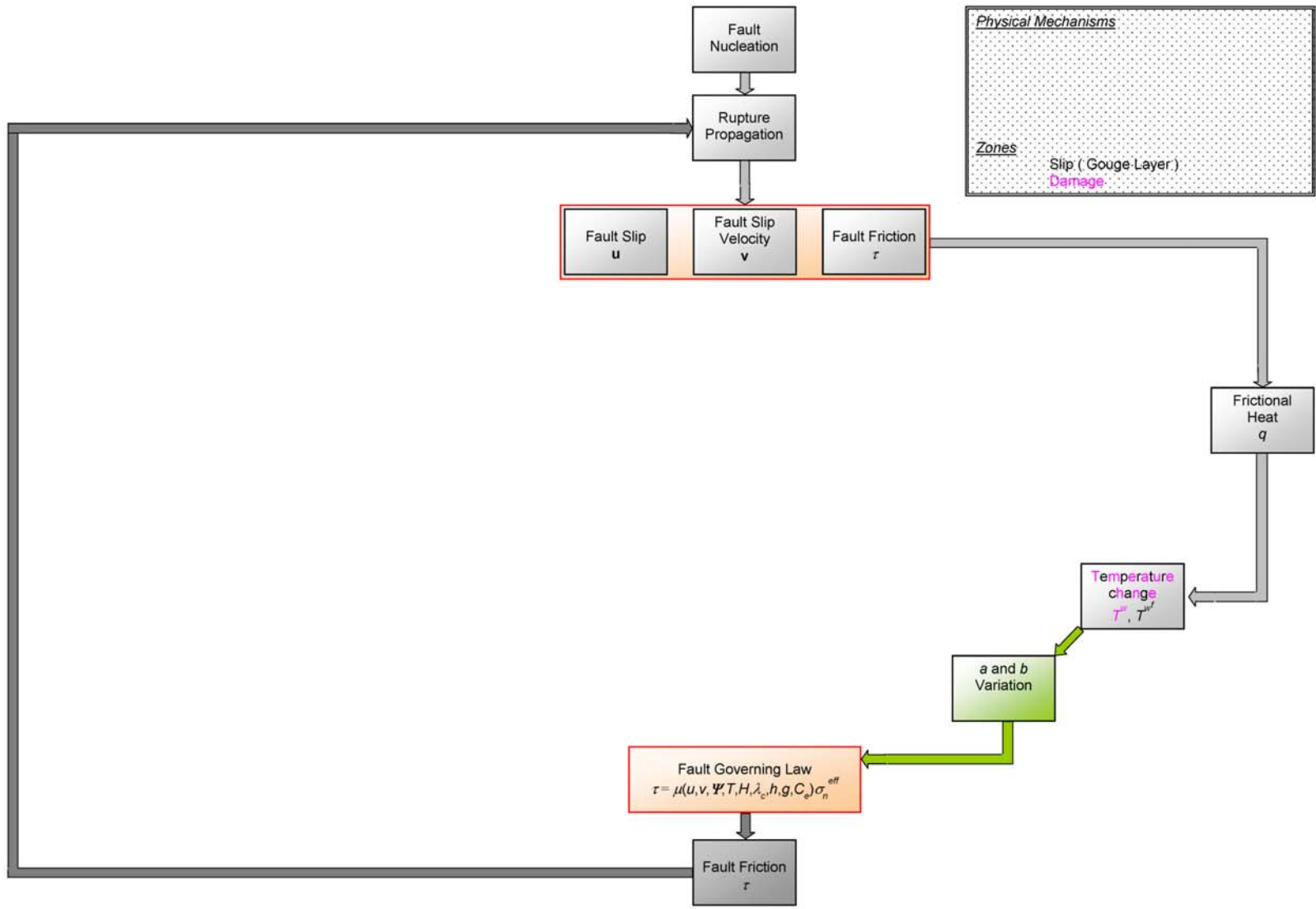


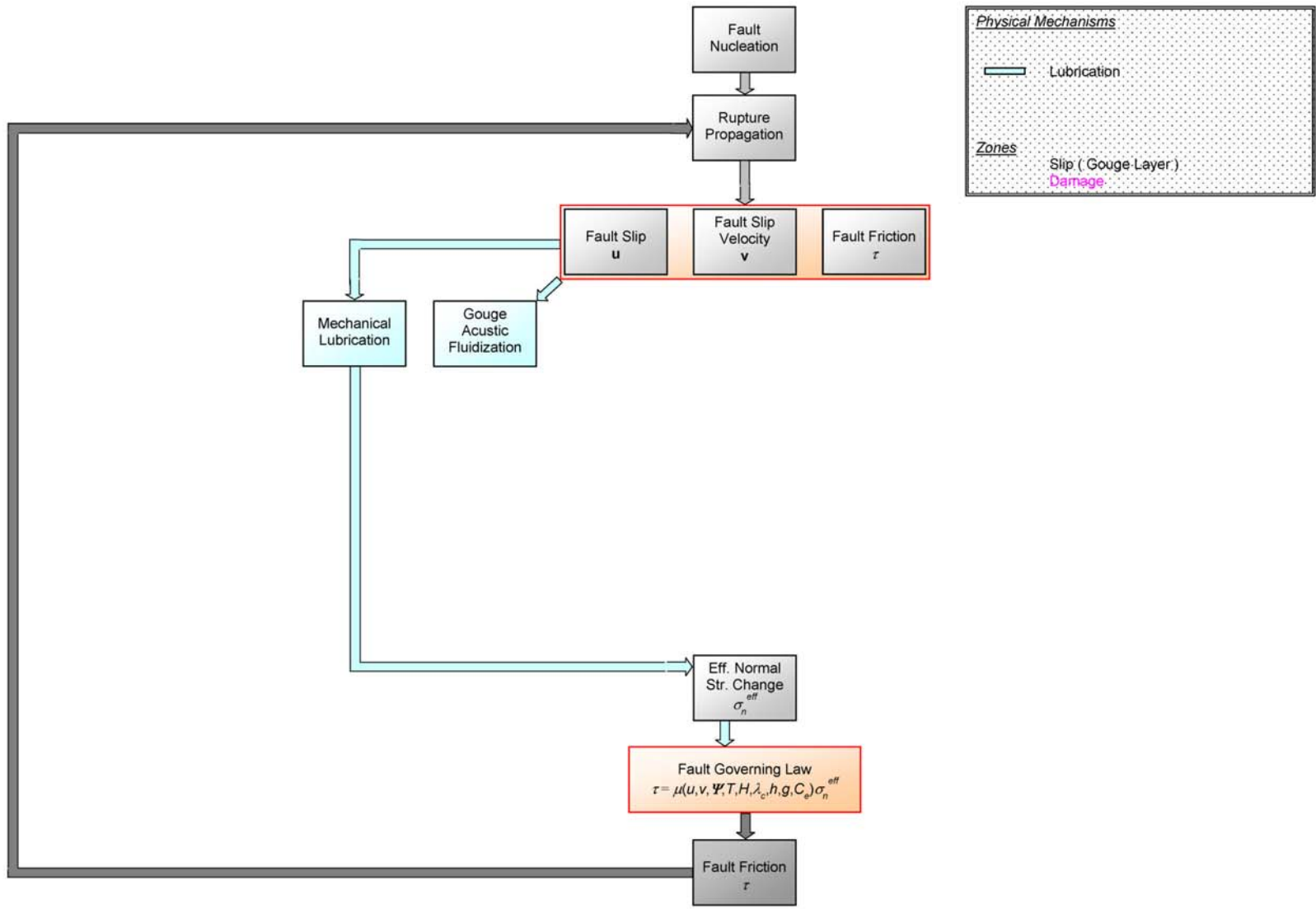


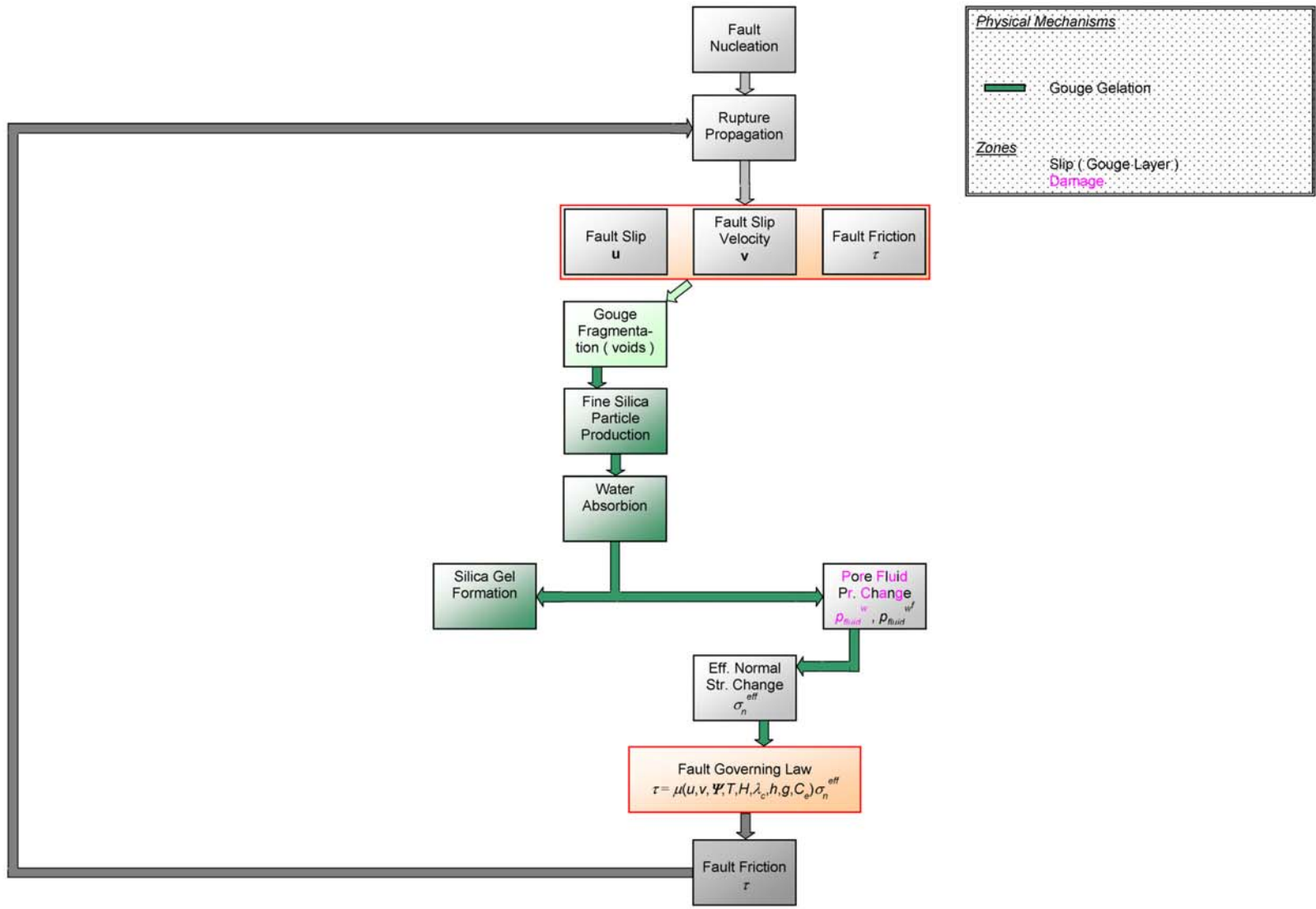


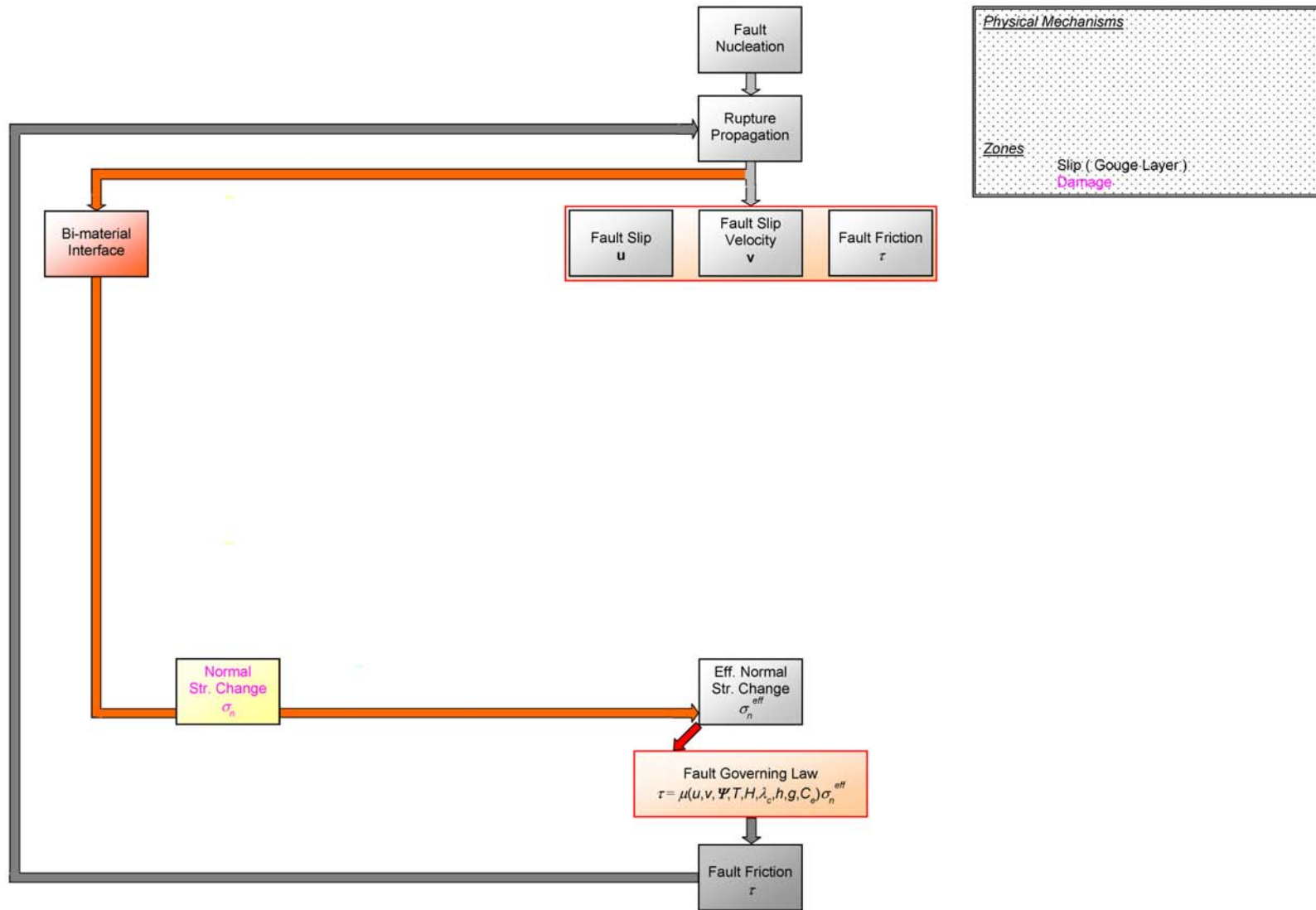


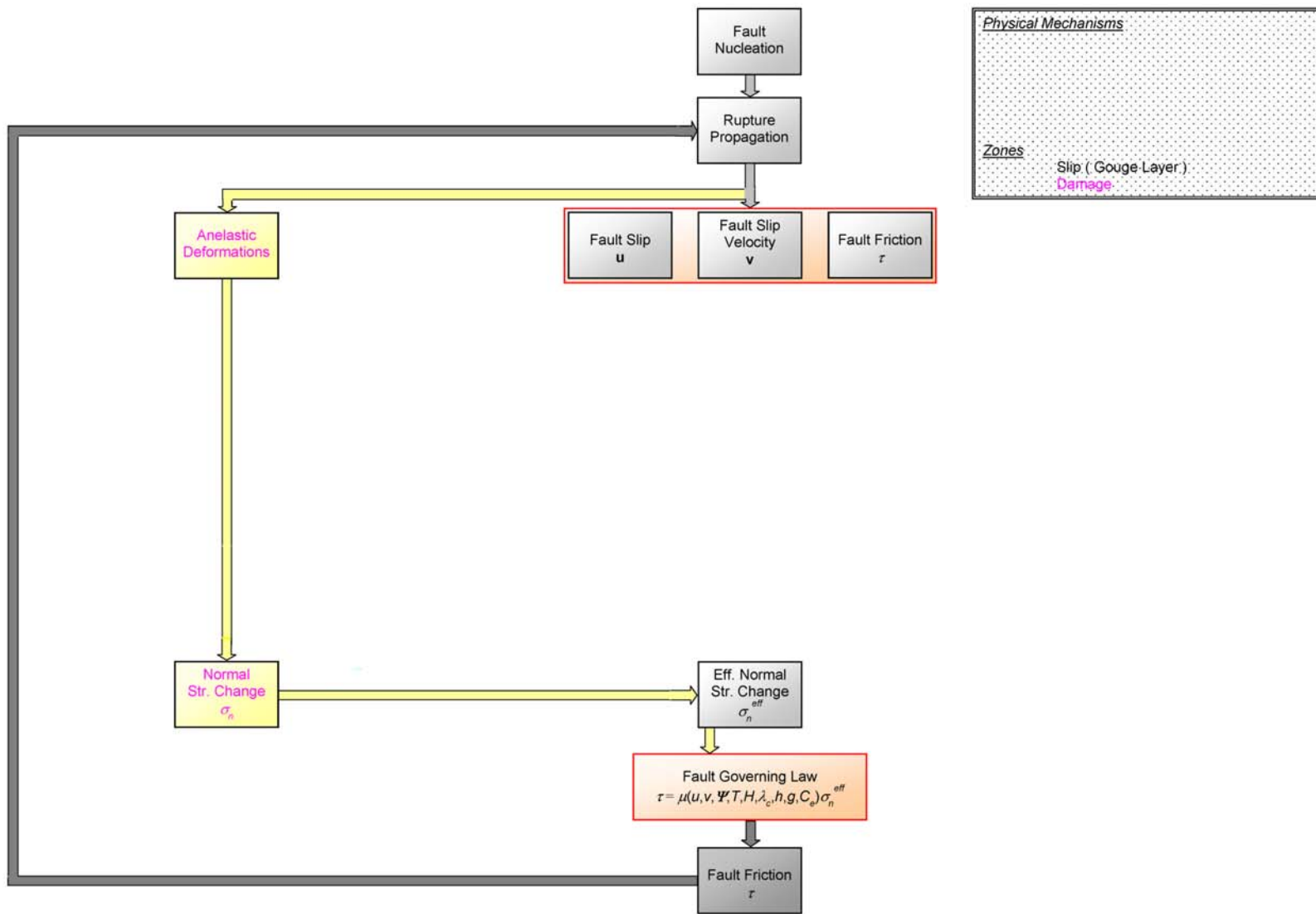


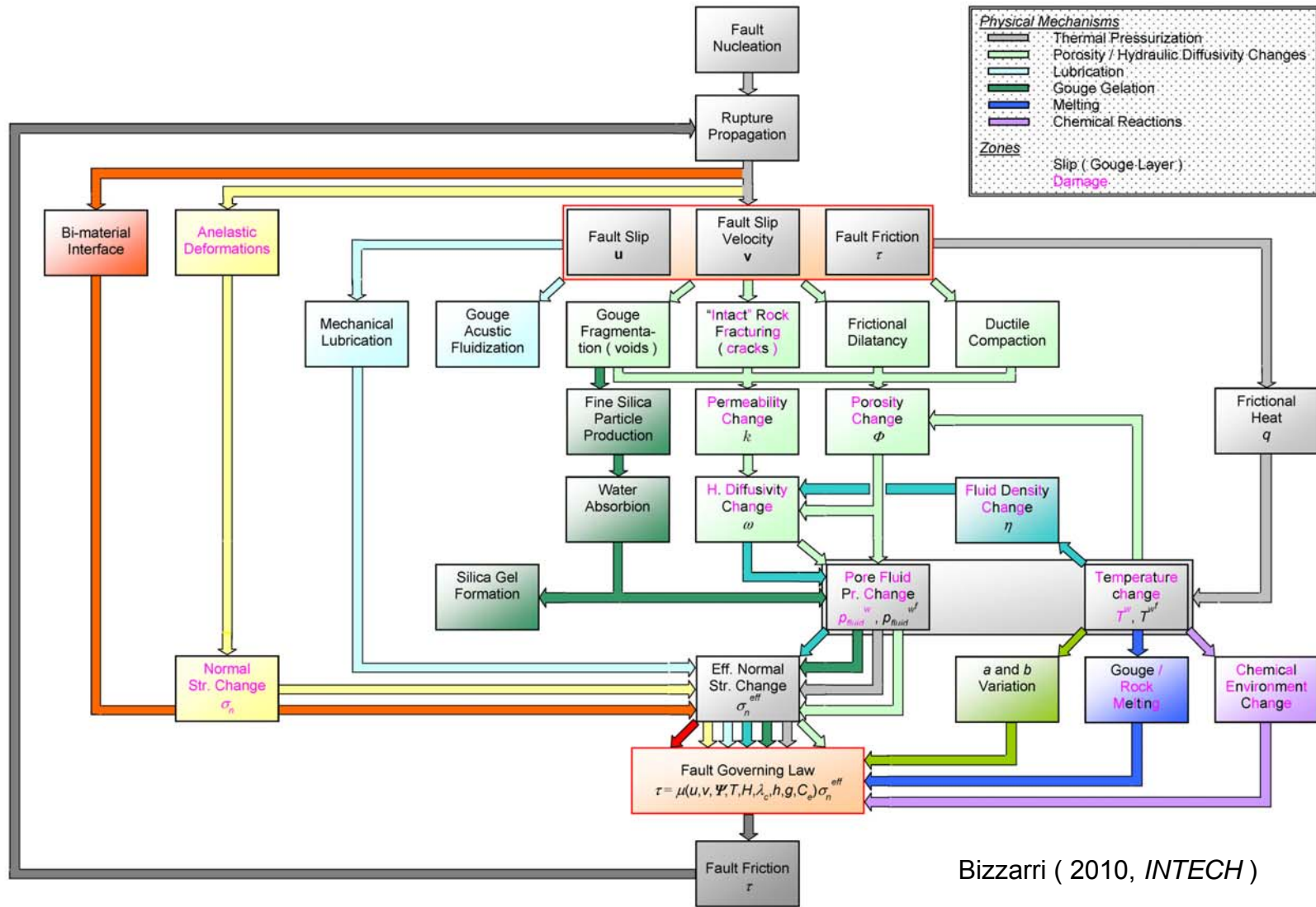












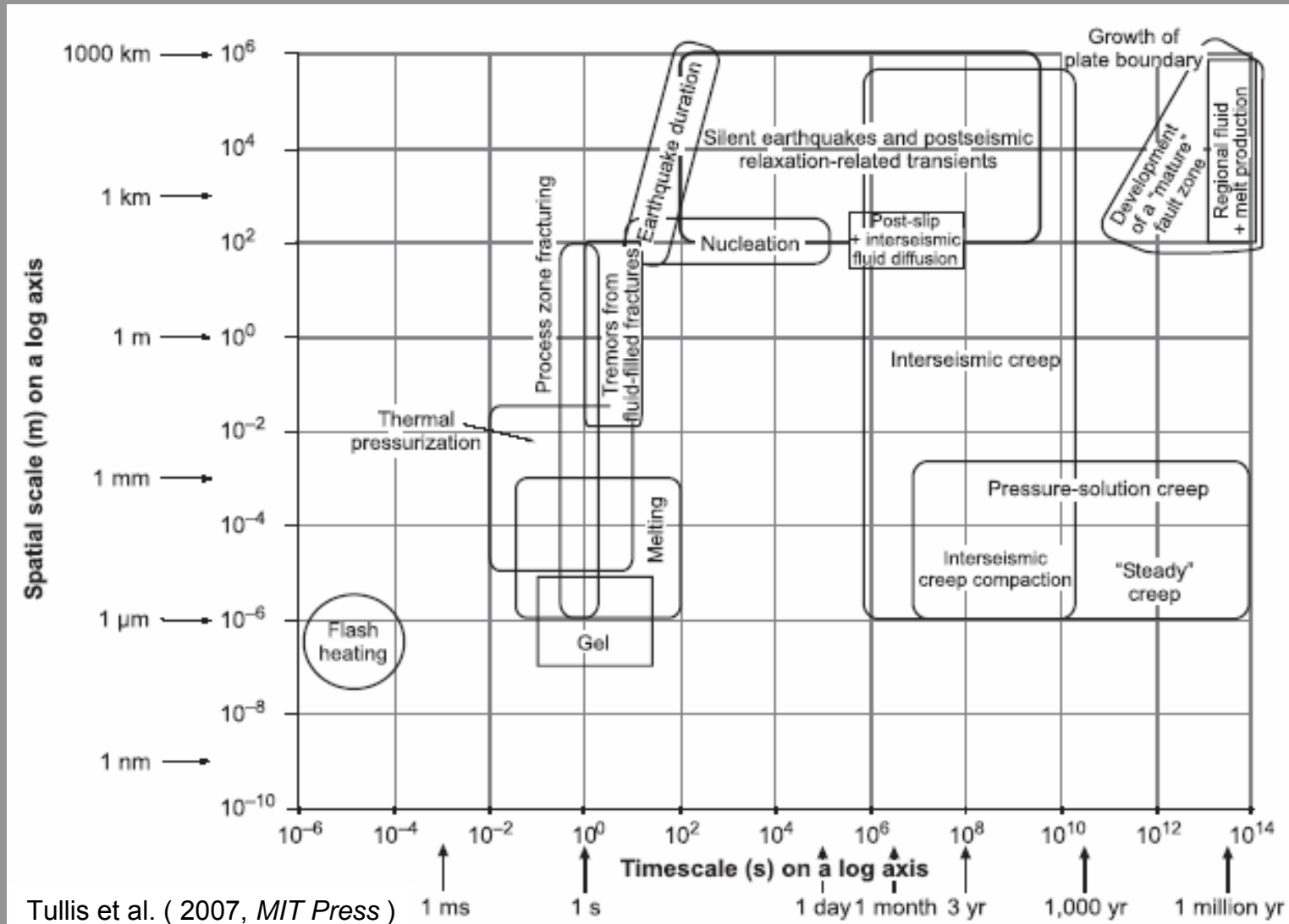
Bizzarri (2010, INTECH)

Occam's razor

- ✓ We follow the logical principle of simplicity (i.e., the Occam's razor):

The simplest way to describe the fault complexity is to **start from the beginning** (i.e., canonical formulations of the governing equations) and then **add** to the model **one by one** all additional phenomena until the empirical (instrumentally recorded) evidence can be explained.

Spatial and temporal scales





Fracture Criteria & Constitutive Laws

1. FRACTURE CRITERION

- Condition that specifies, at a given fault point and at a given time, if there is a rupture or not.
- It can be expressed in terms of **energy**, in terms of **maximum frictional resistance**, and so on.
- It is based on (i) the *Benioff* (1951) hypothesis: The fracture occurs when the stress in a volume reaches the rock strength
or, analogously,
(ii) the *Reid* (1910) statement: The fracture takes place when the stress attains a value greater than the rock can endure.

2. CONSTITUTIVE LAW

- Analytical relation existing between the components of the stress tensor and physical observable(s), like the slip, the slip velocity, the state variable, etc..
- From a mathematical point of view it is a **Fault Boundary Condition (FBC)** that controls earthquake dynamics and its complexity in space and in time.
- Its simplest form consider only **two frictional levels**, τ_u and τ_f ; it accounts for stress drop (or stress release), but the process is instantaneous: there is a singularity at crack tip. 
- **Cohesive zone models**: *Barenblatt (1959a, 1959b)*, *Ida (1972)*, *Andrews (1976a, 1976b)*. In these models the singularity is removed and the stress release occurs over a breakdown zone distance X_b and in a breakdown zone time T_b . 
- Friction laws (Rate and State dependent f. l.): *Dieterich (1976)*, *Ruina (1980, 1983)*. They accounts for fault spontaneous nucleation, re – strengthening, healing, etc..

CONSTITUTIVE LAW (continues)

- “ The central issue is *whether* faults obey simple friction laws, and if so, what is the friction coefficient associated with fault slip “ (*Scholz and Hanks, 2004*).

CONSTITUTIVE LAW (continues)

- In full of generality we can express the constitutive (or governing) as:

$$\tau = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{eff}(\sigma_n, \rho_f)$$



where:

1st – order dependencies

- u is the Slip (i. e. displ. disc.) modulus, ←
- v is the Slip Velocity modulus (its time der.), ←
- $\Psi = (\Psi_1, \dots, \Psi_N)$ is the State Variable vector, ←
- T is the Temperature (accounting for Ductility, Plastic Flow, Melting and Vaporization),
- H is the Humidity,
- λ_c is the Characteristic Length of surface (accounting for Roughness and Topography of asperity contacts),
- h is the Hardness,
- g is the Gouge (accounting for Surface Consumption and Gouge formation),
- C_e is the Chemical Environment

Strength & Constitutive Laws

1. THE STRENGTH PARAMETER

- Historically introduced by *Das and Aki (1977a, 1977b)* to have a quantitative estimate of the ability to fracture for a fault
- Its expression can be generalized as:

$$S = (\mu_u \sigma_n^{eff} - \tau_0) / (\tau_0 - \mu_f \sigma_n^{eff})$$

where μ are the friction coefficient.

- We can also define

2. THE FAULT STRENGTH

- Is the parameter that quantify the Strength in the more general case, in which a fault is described by a rhealistic friction laws

$$S^{fault} = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{eff}(\sigma_n, p_{fluid})$$

Towards real – world conditions

$u_{tot} \sim$ several m

$v \sim$ several m/s

$\sigma_n^{eff} = 100 - 200$ MPa

Classical laboratory

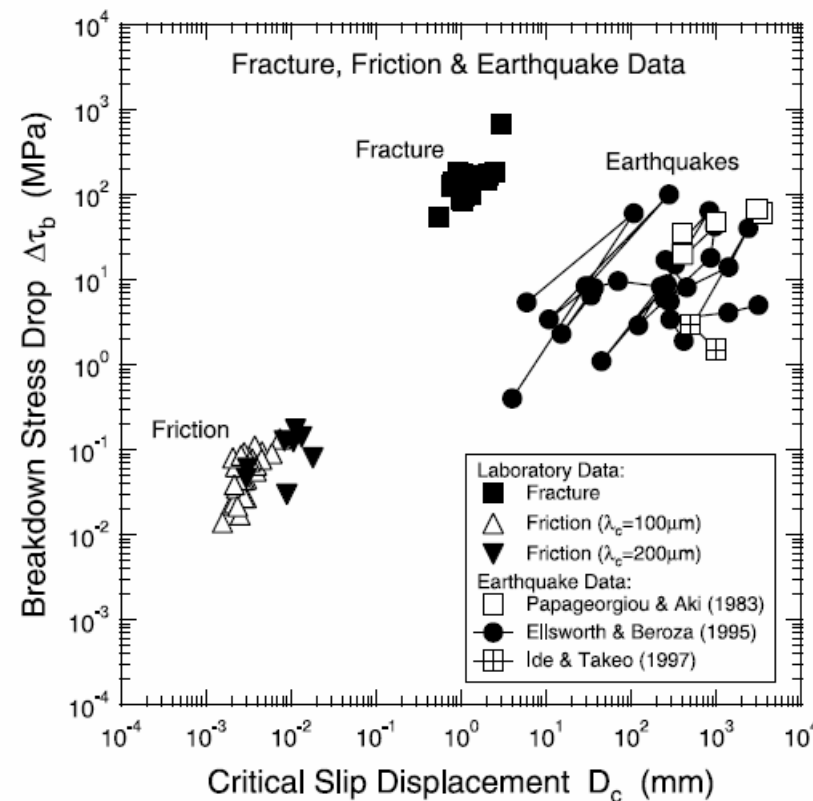
stick – slip experiments

(Dieterich, 1981)

u_{tot} up to 1.4 mm

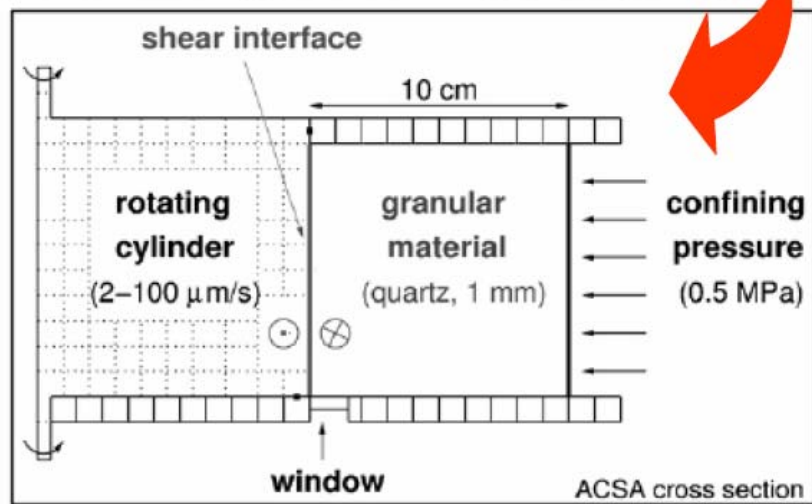
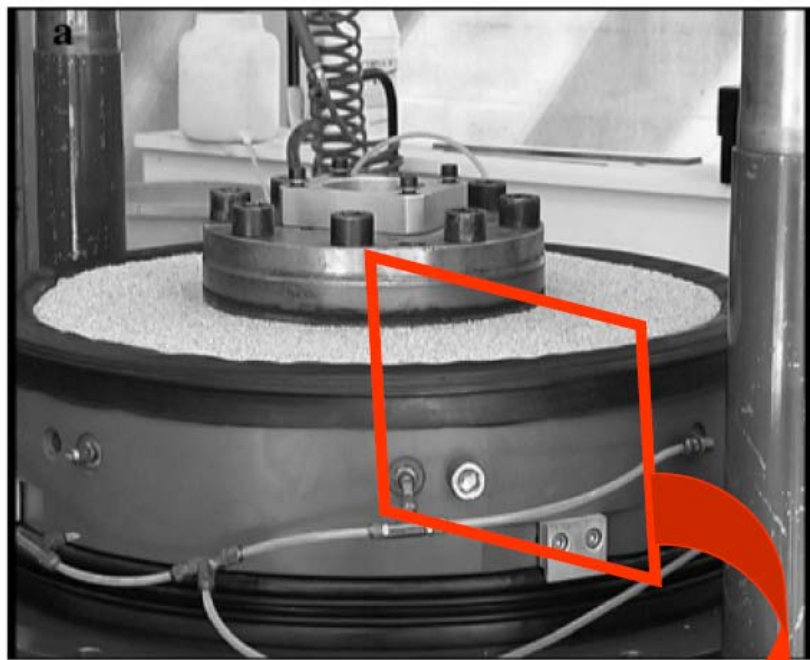
v up to $25 \mu\text{m/s}$

$\sigma_n^{eff} = 10$ MPa



From Ohnaka (2003)

Annular simple shear apparatus



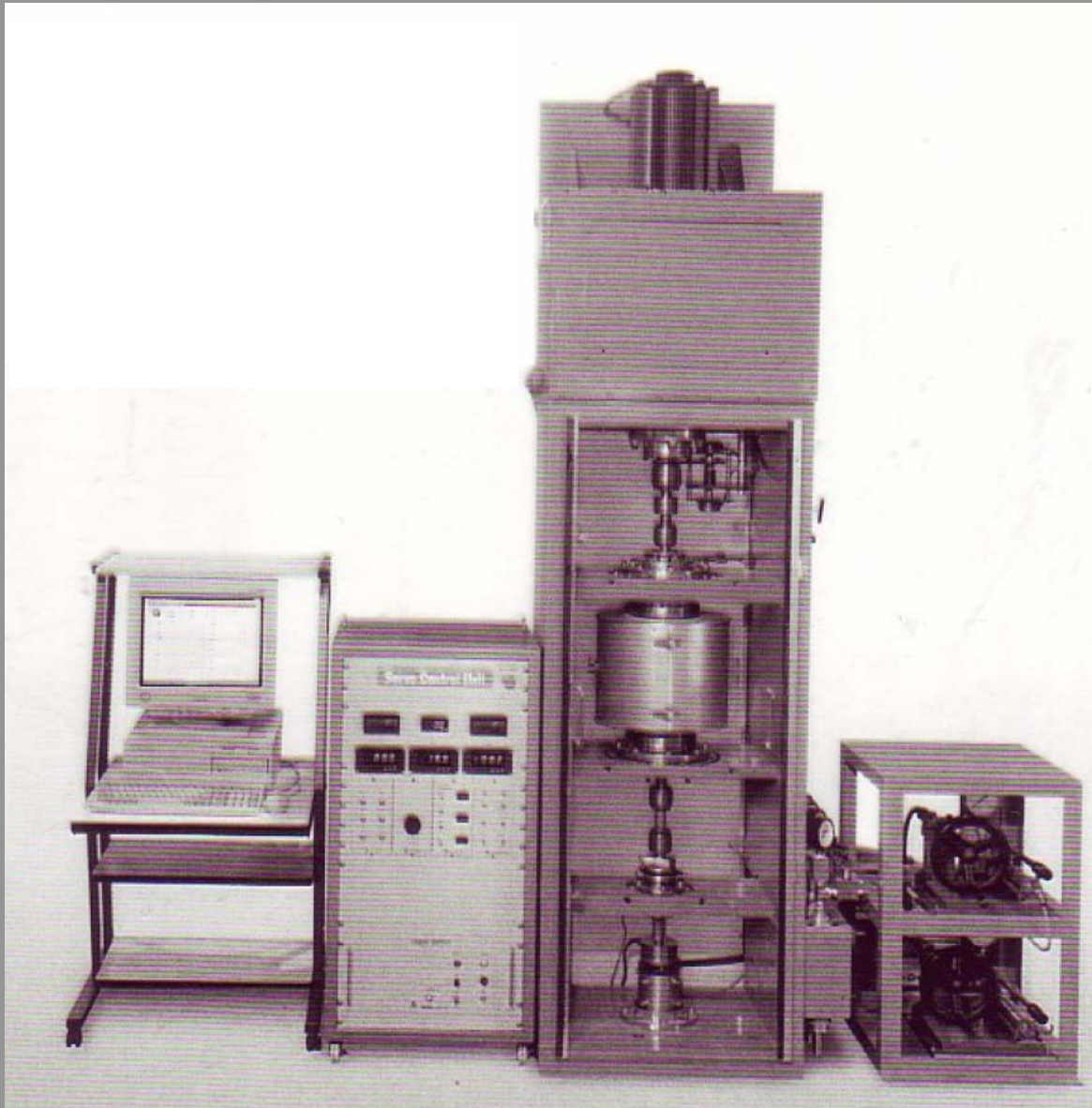
$$u_{tot} < 50 \text{ m}$$

$$v = 1 \text{ } \mu\text{m/s} - 0.1 \text{ mm/s}$$

$$\sigma_n^{eff} < 1 \text{ MPa}$$

Chambon et al. (2006a, 2006b, *JGR*, **111**, B09308, B09309)

High velocity rotary friction apparatus



$U_{tot} = \text{infinite}$

$v = 0.1 \mu\text{m/s} - 10 \text{ m/s}$

$\sigma_n^{eff} < 20 \text{ MPa}$

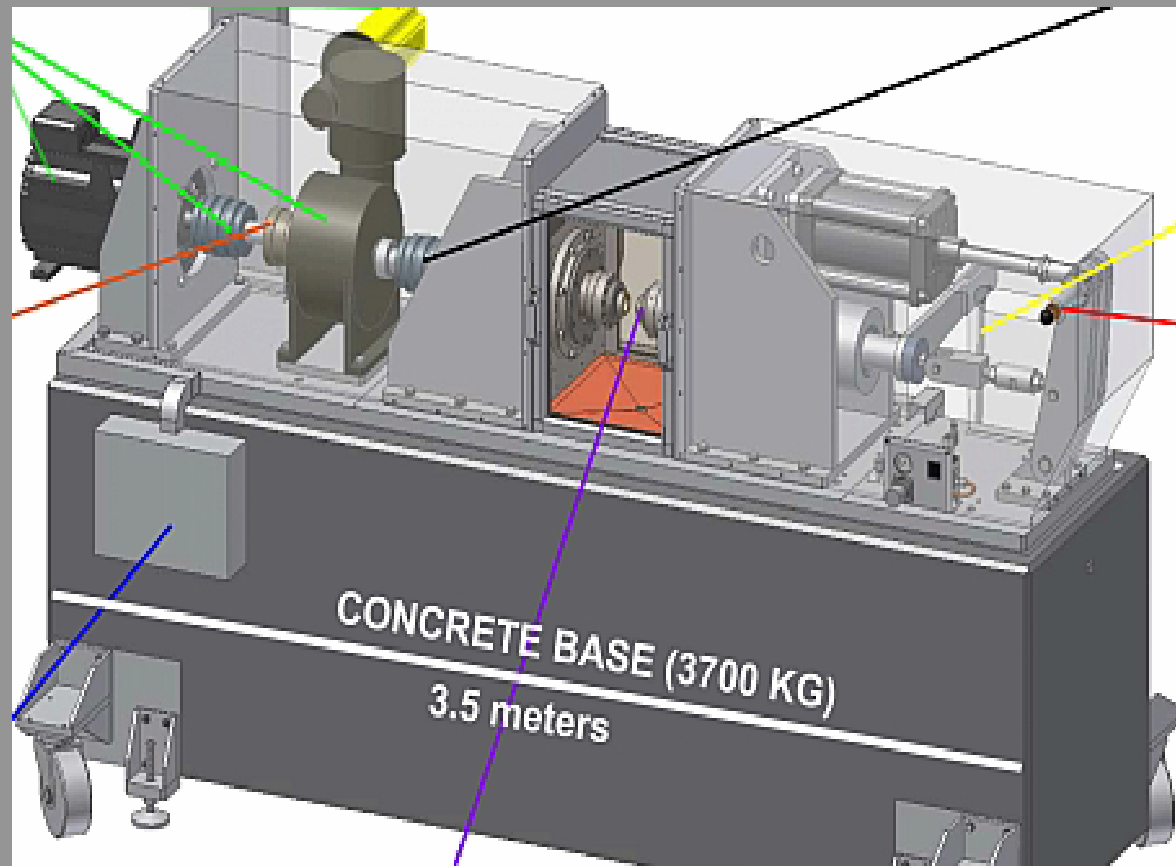
Shimamoto and Tsutumi (2004,
Str. Geol., **39**)

High velocity rotary friction apparatus @ INGV

$$u_{tot} = \text{infinite}$$

$$v = 1 \mu\text{m/s} - 9 \text{ m/s}$$

$$\sigma_n^{eff} < 70 \text{ MPa}$$



Niemeijer et al. (2009, AGU Fall Meeting)

Time - weakening Friction Law

$$\tau = \begin{cases} \left[\mu_u - (\mu_u - \mu_f) \frac{(t - t_r)}{t_0} \right] \sigma_n^{eff} & , t - t_r < t_0 \\ \mu_f \sigma_n^{eff} & , t - t_r \geq t_0 \end{cases}$$

ilaw = 11

TW

$t_r = t_r(\xi)$ is the rupture onset time in every fault point ξ (when $u > 0$).

Andrews (1985), Bizzarri et al. (2001) and other following Bizzarri' s papers

t_0 is the characteristic time – weakening duration.

Position - weakening Friction Law

$$\tau = \begin{cases} \left[\mu_u - (\mu_u - \mu_f) \frac{x}{R_0} \right] \sigma_n^{eff} & , -R_0 < x < 0 \\ \mu_f \sigma_n^{eff} & , -L < x < -R_0 \end{cases}$$

PW

x is the position on the fault
(extending up to $-L$).

Palmer and Rice (1973)

R_0 is the characteristic position –
weakening distance.

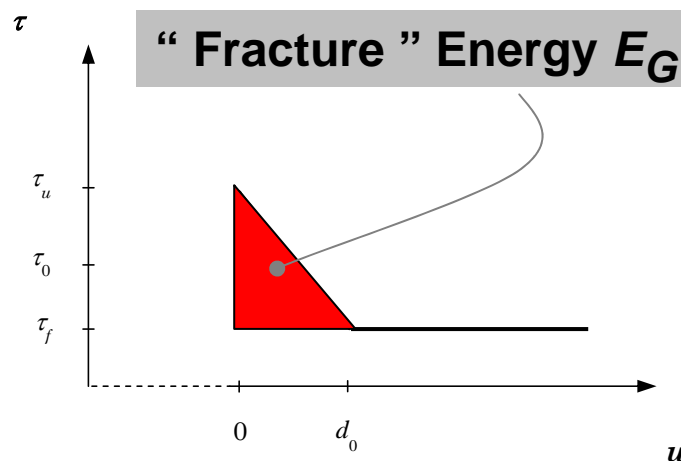
Slip - Dependent Friction Laws

1. LINEAR SLIP – WEAKENING LAW

$$\tau = \begin{cases} \left[\mu_u - (\mu_u - \mu_f) \frac{u}{d_0} \right] \sigma_n^{eff} & , u < d_0 \\ \mu_f \sigma_n^{eff} & , u \geq d_0 \end{cases}$$

ilaw = 21

SW



Barenblatt (1959a, 1959b), Ida (1972), Andrews (1976a, 1976b), and many authors thereafter

d_0 is the characteristic slip – weakening distance

ilaw = 22

IW

2. NON – LINEAR SLIP – WEAKEING LAW

$$\tau = \begin{cases} \left[\mu_u - \frac{\mu_u - \mu_f}{d_0} \left(u - \frac{(1 - p_{IW})d_0}{2\pi} \sin\left(\frac{2\pi u}{d_0}\right) \right) \right] \sigma_n^{eff} & , u < d_0 \\ \mu_f \sigma_n^{eff} & , u \geq d_0 \end{cases}$$

Ionescu and Campillo (1999)

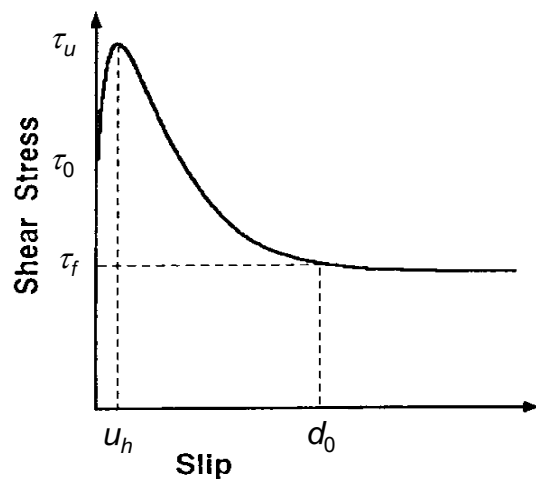
3. NON LINEAR SLIP – WEAKENING LAW WITH SLIP – HARDENING

$$\tau = \left\{ \left[\left(\frac{\tau_0}{\sigma_n^{eff}} - \mu_f \right) \left(1 + \alpha_{OY} \ln \left(1 + \frac{u}{\beta_{OY}} \right) \right) \right] e^{-\frac{u}{d_0}} + \mu_f \right\} \sigma_n^{eff}$$

ilaw = 23

OW

$$u_h : \left. \frac{d\tau}{du} \right|_{u_h} = 0; \quad \begin{cases} u_h = rd_0 \quad (\text{e.g. } r = 0.1) \\ \tau(u_h) = \tau_u \end{cases}$$



Ohnaka and Yamashita (1989) and the following papers by Ohnaka and coworkers

u_h is associated with the preparatory phase of the imminent macroscopic failure in the cohesive zone. It accounts for micro-cracking

4. NON LINEAR SLIP – WEAKENING LAW WITH EXPONENTIAL DECAY

$$\tau = \left[(\mu_u - \mu_f) e^{-\frac{u}{d_0}} + \mu_f \right] \sigma_n^{eff}$$

ilaw = 24

EW

5a. POWER LAW SLIP – WEAKENING

$$\tau = \left\{ \mu_u - (\mu_u - \mu_f) \left[\left(\frac{p_{PW}}{p_{PW} + 1} \right) \frac{u}{d_0} \right]^{p_{PW}} \right\} \sigma_n^{eff}$$

$$i_{law} = 25$$

PW

5b. POWER LAW SLIP – WEAKENING II

$$\tau = \left\{ \mu_f + \frac{\alpha_{CEA}}{\sigma_n^{eff}} (u - d_0)^{p_{CEA}} \right\}$$

$$\alpha_{CEA} = 5.6 \times 10^{-2} \text{ MPa m}^{p_{CEA}}$$

$$p_{CEA} = 0.4$$

Chambon et al. (2006b)

Slip - and Rate - Dependent Friction Laws

$$\tau = \left\{ \mu^{ss}(v) + \left[F(u)\mu_i - \mu^{ss}(v) \right] e^{\frac{\ln(0.05)u}{d_0}} \right\} \sigma_n^{eff}$$

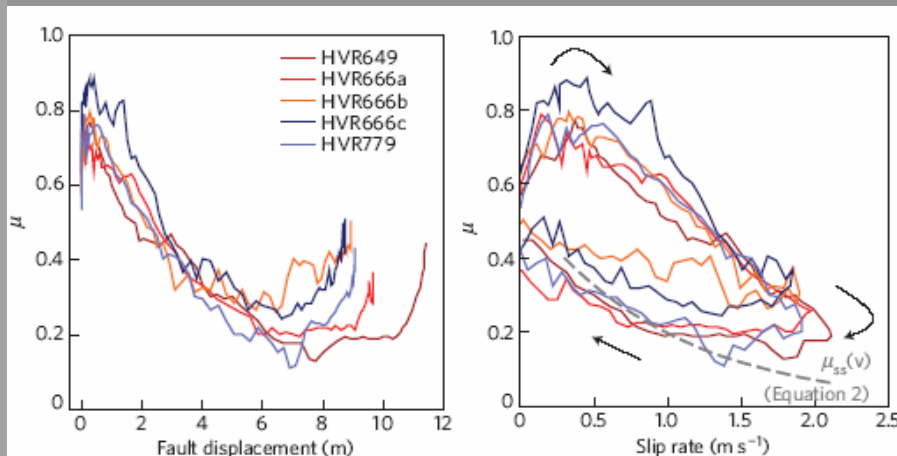
ilaw = 26

$$\mu^{ss}(v) = \mu^{ss}(0) e^{-\frac{v}{v_{SS}}}$$

$$F(u) = \alpha_{SS} + (1 - \alpha_{SS}) e^{\frac{\ln(0.05)u}{u_h}}$$

VW

Sone and Shimamoto (2009)



u_h controls the duration in slip of the slip – hardening phase, described by the function $F(u)$.

$$\mu^{ss}(0) = 0.55 \pm 0.09$$

$$\mu_i = 0.6$$

$$v_{SS} = 0.99 \pm 0.23 \text{ m/s}$$

$$\alpha_{SS} = 1.26 \div 1.54$$

$$u_h = 23 \div 160 \text{ mm}$$

Rate - Dependent Friction Law

$$\tau = \frac{U_*}{U + U_*} \mu_u \sigma_n^{eff}$$

*Burrige and Knopoff (1967),
Carlson and Langer (1989),
Madariaga and Cochard (1994),
Cochard and Madariaga (1994)*

Rate - and State - Dependent Friction Laws



1. DIETERICH IN REDUCED FORMULATION

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_*}{v} \right) + b \ln \left(\frac{\Psi v_*}{L} \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} \end{array} \right.$$

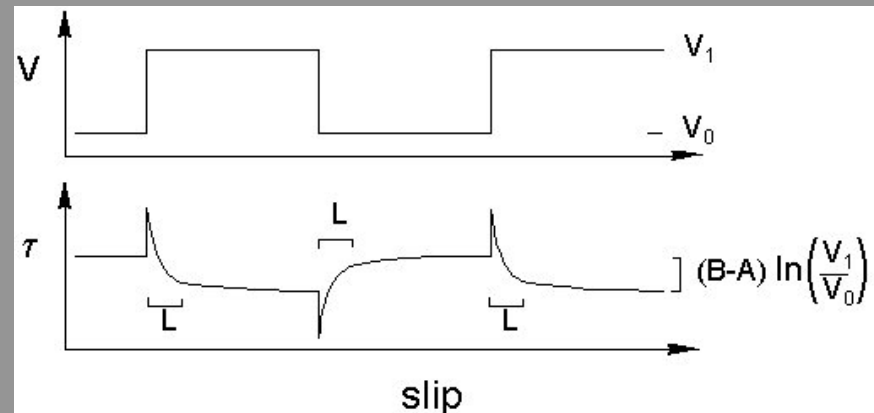
ilaw = 31

DR

However, while in velocity stepping experiments the traction response following the velocity variation is directly controlled by the parameter L , its effects are much less evident during the dynamic rupture propagation.

Bizzarri and Cocco (2005)

Response to an abrupt jump in load



2. RUINA – DIETERICH (RUINA ORIGINAL FORM)

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_*}{v} \right) + \theta \right] \sigma_n^{eff} \\ \frac{d}{dt} \theta = -\frac{v}{L} \left[\theta + b \ln \left(\frac{v}{v_*} \right) \right] \end{array} \right.$$

Ruina (1980, 1983)

2bis. RUINA – DIETERICH (RUINA MODERN FORM.)

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{\Psi v}{L} \ln\left(\frac{\Psi v}{L}\right) \end{array} \right.$$

ilaw = 32

RD

Beeler et al. (1994), Roy and Marone
(1996)

3. DIETERICH – RUINA WITH VARYING NORMAL STR.

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - \alpha \ln \left(\frac{v_*}{v} \right) + b \ln \left(\frac{\Psi v_*}{L} \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} - \left(\frac{\alpha_{LD} \Psi}{b \sigma_n^{eff}} \right) \frac{d}{dt} \sigma_n^{eff} \end{array} \right.$$

ilaw = 31

decis10=T

DR

Linker and Dieterich (1992), Dieterich and Linker (1992), Bizzarri and Cocco (2006a, 2006b)

4. RUINA – DIETERICH WITH VARYING NORMAL STR.

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - \alpha \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{\Psi v}{L} \ln\left(\frac{\Psi v}{L}\right) - \left(\frac{\alpha_{LD} \Psi}{b \sigma_n^{eff}}\right) \frac{d}{dt} \sigma_n^{eff} \end{array} \right.$$

ilaw = 32

decis10=T

RD

Linker and Dieterich (1992), Bizzarri and Cocco (2006a, 2006b)

5. DIETERICH IN REDUCED FORM REGULARIZED

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v + v_*}{v + v_r} \right) + b \ln \left(\frac{\Psi(v - v_r)}{L} + 1 \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi(v - v_r)}{L} \end{array} \right.$$

ilaw = 33

DE

v_r is a regularization fault slip velocity

Perrin et al. (1995), Cocco et al. (2004)

6. RUINA REGULARIZED

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_* - v_r}{v + v_r} \right) + \frac{\Psi}{\sigma_n^{eff}} \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{v + v_r}{L} \left(\Psi + b \ln \left(\frac{v + v_r}{v_* - v_r} \right) \right) \end{array} \right.$$

ilaw = 34

RE

v_r is a regularization fault slip velocity

Bizzarri (2002, unpublished work)

7. DIETERICH IN REDUCED FORM WITH HEALING

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v^*}{v} + 1 \right) + b \ln \left(\frac{\Psi v^*}{L} + 1 \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = \frac{\gamma_{fh} - \Psi}{t_{fh}} - \frac{\Psi v}{L} \end{array} \right.$$

ilaw = 35

DH

$$\gamma_{fh} = 1 \text{ s}$$

t_{fh} is the time for healing (slip duration)

Evolution law proposed by Nielsen et al. (2000) and by Nielsen and Carlson (2000). Used in this form by Cocco et al. (2004)

8. DIETERICH IN REDUCED FORM WITH 2 STATE VAR.

ilaw = 36

DW

Tullis and Weeks (1993). Used in this form by *Bizzarri (xxxx, unpublished work)*

9. PRAKASH – CLIFTON

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_*}{v} \right) + b \ln \left(\frac{\Psi v_*}{L} \right) \right] \left(\frac{d}{dt} \Psi_1 + \frac{d}{dt} \Psi_2 \right) \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} \\ \frac{d}{dt} \Psi_1 = - \frac{v}{L_1} \left(\Psi_1 - \alpha_{PC_1} \sigma_n^{eff} \right) \\ \frac{d}{dt} \Psi_2 = - \frac{v}{L_2} \left(\Psi_2 - \alpha_{PC_2} \sigma_n^{eff} \right) \end{array} \right.$$

ilaw = 37

PC

Ψ_1 and Ψ_2 are additional state variables accounting for the coupling with effective normal stress. The formulation of friction law is not based on the Amonton – Coulamb law.

Coupling with effective normal stress proposed by Prakash and Clifton (1993) and Prakash (1998). Used in this form by Bizzarri (2005, unpublished work)

10. RUINA – DIETERICH WITH FLASH HEATING

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_*}{v} \right) + \theta \right] \sigma_n^{eff} \\ \frac{d}{dt} \theta = \begin{cases} -\frac{v}{L} \left[\theta + b \ln \left(\frac{v}{v_*} \right) \right] & , v \leq v_{fh} \\ -\frac{v}{L} \left[\theta + b \frac{v_{fh}}{v} \ln \left(\frac{v}{v_*} \right) + \left(1 - \frac{v_{fh}}{v} \right) \left(a \ln \left(\frac{v}{v_*} \right) + \mu_* - \mu_{fh} \right) \right] & , v > v_{fh} \end{cases} \end{array} \right.$$

ilaw = 38

FH

where $v_{fh} = \frac{\pi\chi}{D_{ac}} \left(c \frac{T_{weak} - T^{wf}}{\tau_{ac}} \right)^2$ is a weakening velocity above which flash heating is activated, T^{weak} is a weakening temperature, τ_{ac} is the (average) shear strength of asperity contacts and D_{ac} their (average) size.

Beeler and Tullis (2003); Tullis and Goldsby (2003a, 2003b). Rice (1999, 2006). Modified from Noda et al. (2009)

11. RUINA – DIETERICH WITH TEMPERATURE DEPEN.

$$\left\{ \begin{array}{l} \tau = \left[\mu_* - a \ln \left(\frac{v_*}{v} \right) + \theta \right] + \frac{a Q_a}{R} \left(\frac{1}{T} - \frac{1}{T_*} \right) \\ \frac{d}{dt} \theta = -\frac{v}{L} \left[\theta + b \ln \left(\frac{v}{v_*} \right) \right] + \frac{b Q_b}{R} \left(\frac{1}{T} - \frac{1}{T_*} \right) \end{array} \right. \sigma_n^{eff}$$

ilaw = 39

CH

where Q_a and Q_b are activation energies (Kato, 2001 assumes: $Q_a = Q_b = 0.1$ MJ/mol) and T_* is a reference absolute temperature.

Note that T is the absolute temperature.

Chester and Higgs (1992), Kato (2001)

Slip - and State - Dependent Friction Law

$$\tau = \begin{cases} \left[(\mu_u - \Delta\mu) \left(1 - \frac{u}{d_1} \right) \right] \sigma_n^{eff} & , u < d_1, \Psi \geq \Psi_1 \\ 0 & , u \geq d_1, \Psi \geq \Psi_0 \\ \mu_{sp} \left(1 - \frac{\Psi}{\Psi_0} \right) \sigma_n^{eff} & , \Psi < \Psi_0, \Psi < \Psi_1 \end{cases}$$

$$\frac{d}{dt} \Psi = - \frac{\beta_{CM}}{d_0} (\Psi - v)$$

ilaw = 41

CM

$\Delta\mu$ is an initial artificial stress drop

$$\Psi_1 \equiv \Psi_0 (u - u_1) / (d_1 - u_1)$$

$$U_1 \equiv - d_1 (\mu_{sp} - \mu_u + \Delta\mu) / (\mu_u - \Delta\mu)$$

d_0 and d_1 are characteristic lengths

$\mu_{sp} = 0 \Rightarrow$ linear SW with d_1 as characteristic length

Cochard and Madariaga (1994)

Free Volume Friction law

$$\left\{ \begin{array}{l} \tau = \sigma_d \operatorname{Arcsinh} \left(\frac{v}{v_*} \frac{e^{\frac{f_* + \chi_s + \chi_h}{\chi}}}{1 - m_0} \right) \\ \frac{d}{dt} \chi = -R_c e^{-\frac{\chi_c}{\chi}} + \alpha_{FV} \tau v \\ m_0 = \begin{cases} 1 & , \tau \leq \tau_0 e^{\frac{\chi_h}{\chi}} \\ \frac{\tau_0}{\tau} e^{\frac{\chi_h}{\chi}} & , \tau > \tau_0 e^{\frac{\chi_h}{\chi}} \end{cases} \end{array} \right.$$

ilaw = 51

FV

$\chi \equiv \Phi - \Phi_0$ free volume variable

χ_s reference value of χ for shearing

χ_h FV value required to create a Shear Transformation Zone (STZ)

χ_c FV value for compaction

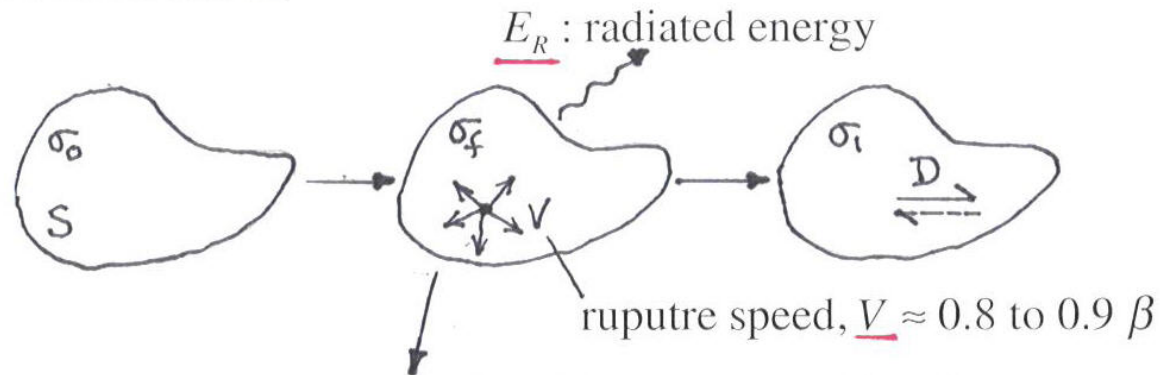
R_c rate of compaction

α_{FV} scaled dilatancy coefficient

*Falk and Langer (1998, 2000);
Lemaitre (2002); Daub and Carlson (2008)*

How to relate relevant quantities to constitutive parameters

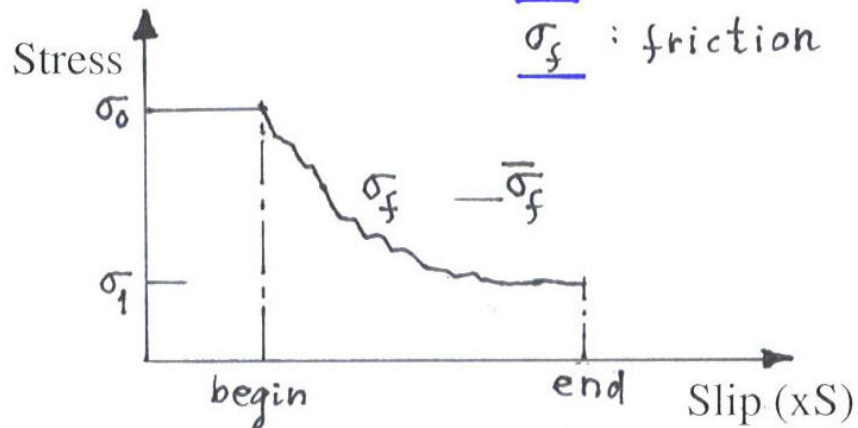
Dynamic Parameters



$E_{NR} = E_F + E_G + \dots$: non-radiated energy

E_F : friction (heat), E_G : fracture energy

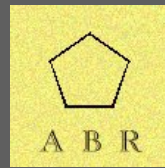
σ_f : friction



$\Delta\sigma_s = \sigma_0 - \sigma_1$: static stress drop

$\Delta\sigma_d = \sigma_0 - \bar{\sigma}_f$: dynamic stress drop

This slide is empty intentionally.



Support Slides: Parameters, Notes, etc.

To not be displayed directly. Referenced above.



Thermal pressurization:

*Sibson (1973); Lachenbruch (1980); Mase and Smith (1985, 1987);
Andrews (2002); Bizzarri and Cocco (2006b, 2006c) .*

Morrow et al. (1984) show that gouge contains water



Gouge behaviour:

*Marone et al. (1990); Marone and Kilgore (1993); Mair and Marone (1999);
Mair et al. (2002); Chambon et al. (2002); Mizogichi et al. (2007)*

Frictional melting:

Jeffreys (1942); McKenzie and Brune (1972); Richards (1977); Sibson (1977); Cardwell et al. (1978); Allen (1979); Nielsen et al. (2007)



Pseudo -
tachylyte: Fault
vein (*Sibson,*
1975)



Mechanical lubrication:

Spray (1993); Brodsky and Kanamori (2001); Kanamori and Brodsky (2001)

Acoustic fluidization:

Melosh (1979, 1996)



Gouge gelation:

Goldbsy and Tullis (2002); Di Toro et al. (2004)

Bi – material interface:

Andrews and Ben – Zion (1997); Harris and Day (1997); Andrews and Harris (2005); Ben – Zion (2006a, 2006b); Dunham and Rice (2008)



MTL: Fractured mylonite,
cataclasite and gouge



Humidity effects:

Dieterich and Conrad (1984); Hirose and Bystricky (2007)

Characteristic length of surface effects:

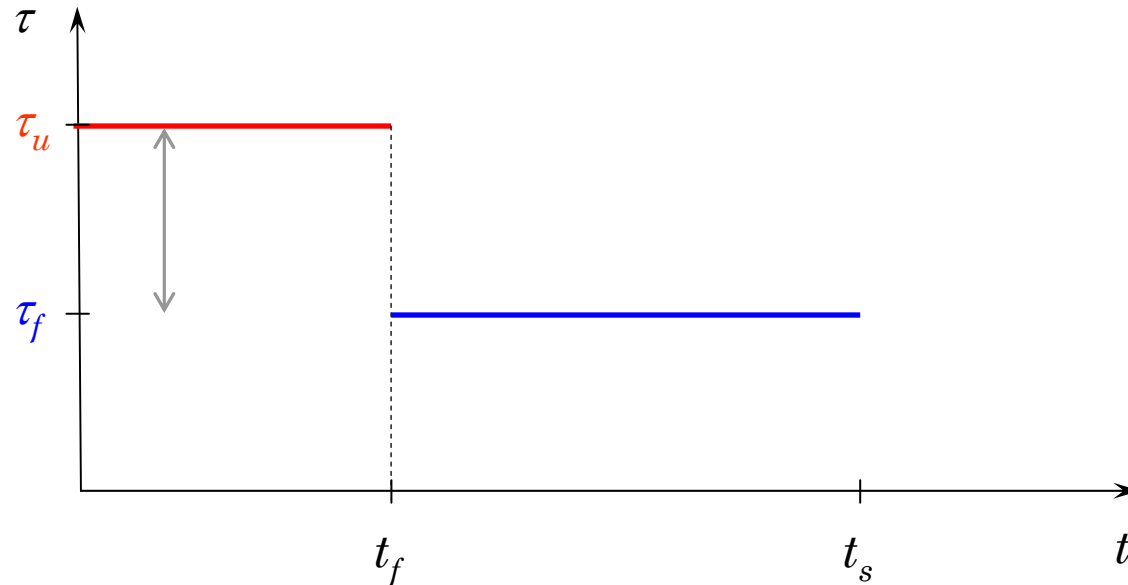
Ohnaka and Shen (1999); Ohnaka (2003)

Simplest friction models

At a particular fault point ξ (following *Savage and Wood, 1971; Scholz, 1990*)

Maximum (or upper, or yield) stress

Kinetic (or frictional) stress



Strength excess: $\tau_u - \tau_0 = 0$

Dynamic stress drop: $\Delta\tau_d = \tau_0 - \tau_f$

In the Dugdale' s model (*Dugdale, 1960; Barenblatt, 1962*) the drop occurs when $u = d_0$.

Failure time (or rupture onset)

Rupture arrest

Simplest friction models

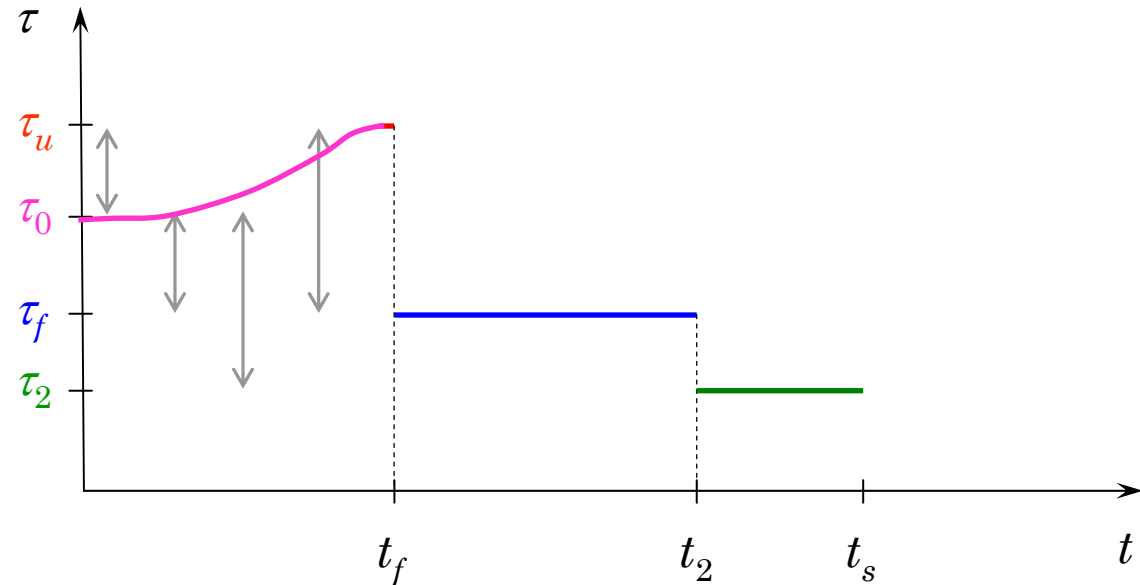
At a particular fault point ξ (following *Savage and Wood, 1971; Scholz, 1990*)

Maximum (or upper, or yield) stress

Initial stress

Kinetic (or frictional) stress

Residual stress



Strength excess: $\tau_u - \tau_0$

Dynamic stress drop: $\Delta\tau_d = \tau_0 - \tau_f$

Static stress drop: $\Delta\tau_s = \tau_0 - \tau_2$

Breakdown str. drop: $\Delta\tau_b = \tau_u - \tau_f$

Failure time (or rupture onset)

Dynamic overshoot

Rupture arrest



- *Savage and Wood (1971)* also define:

Mean stress: $\langle \tau \rangle = \frac{1}{2} (\tau_u + \tau_2)$

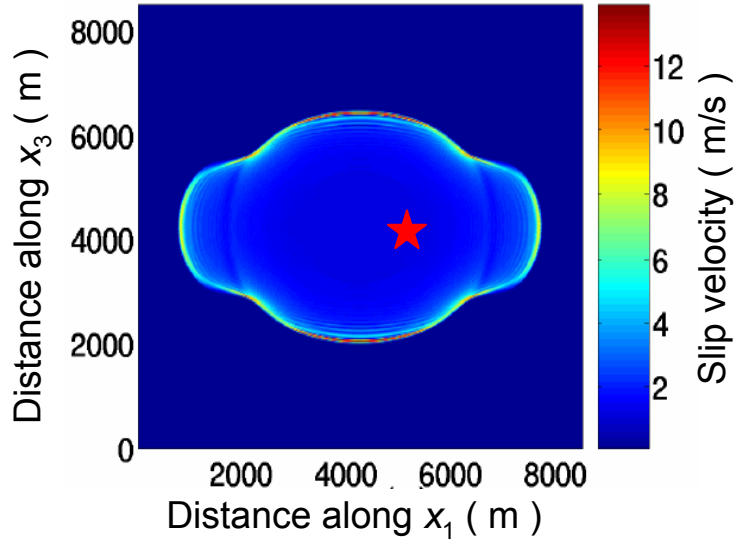
Seismic efficiency: $\eta = E_s / E$, where: E_s is the seismic energy
 E is the total available energy

Apparent stress: $\tau_a = \eta \langle \tau \rangle$

- Direct observation of the absolute stress near an earthquake is not feasible, but it is possible (*Wyss and Brune, 1968*) calculate τ_a and stress drop from physical observables.

The cohesive zone

Time snapshot
($t = 0.8 \text{ s}$) - SW law



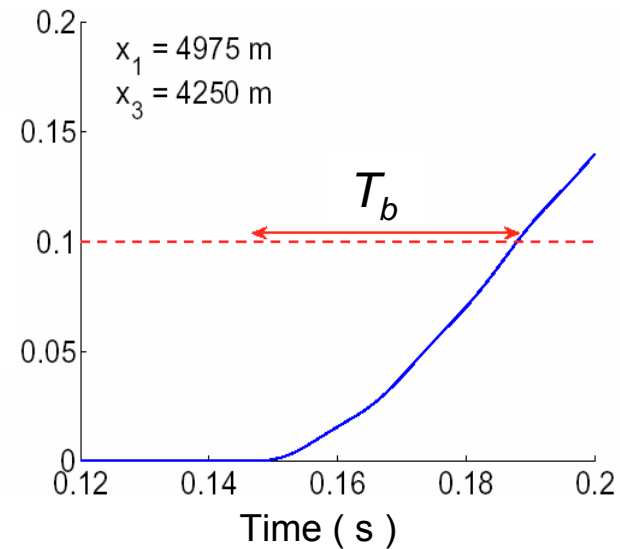
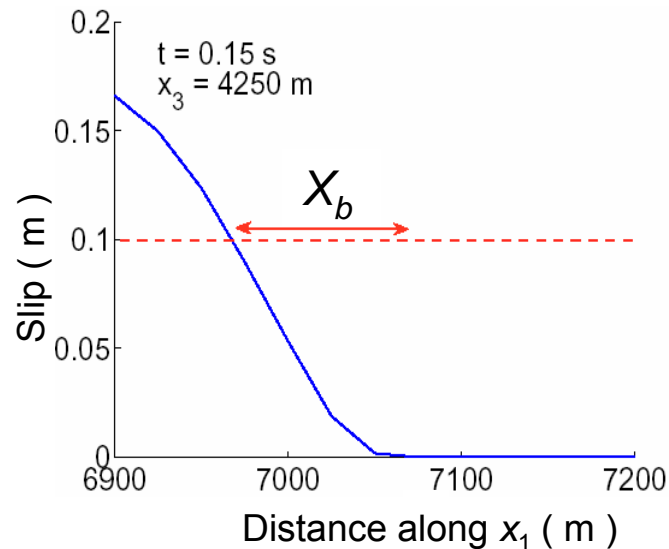
In the target location we can estimate:

$$X_b = 105 \text{ m} \quad T_b = 0.04 \text{ s}$$

From these quantities:

$$v_{rupt} = X_b / T_b = 2625 \text{ m/s}$$

Local estimate



Slip - hardening effect



- * The slip – hardening (**SH**) phenomenon has been also found in seismological inversion studies (e. g. *Quin, 1990; Miyatake, 1992; Mikumo and Miyatake, 1993; Beroza and Mikumo, 1996; Ide, 1997; Bouchon, 1997*).

Interpretation of the state variable

