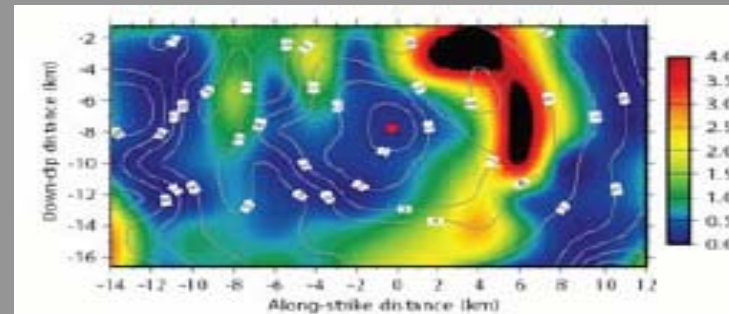




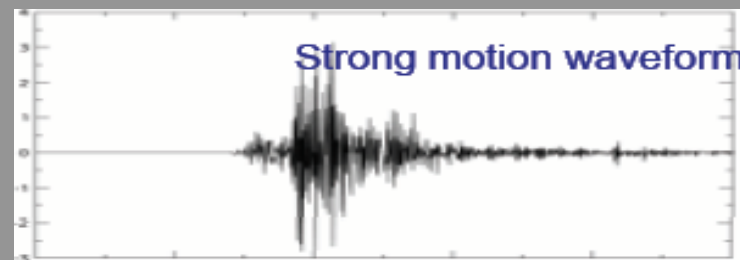
**Fault governing laws  
( constitutive equations )**

# Seismologists need traction

- ✓ To apply fracture mechanics on mathematical planes representing the fault surfaces;
- ✓ To numerically simulate the spontaneous rupture nucleation, propagation, healing and arrest in dynamic earthquake models;

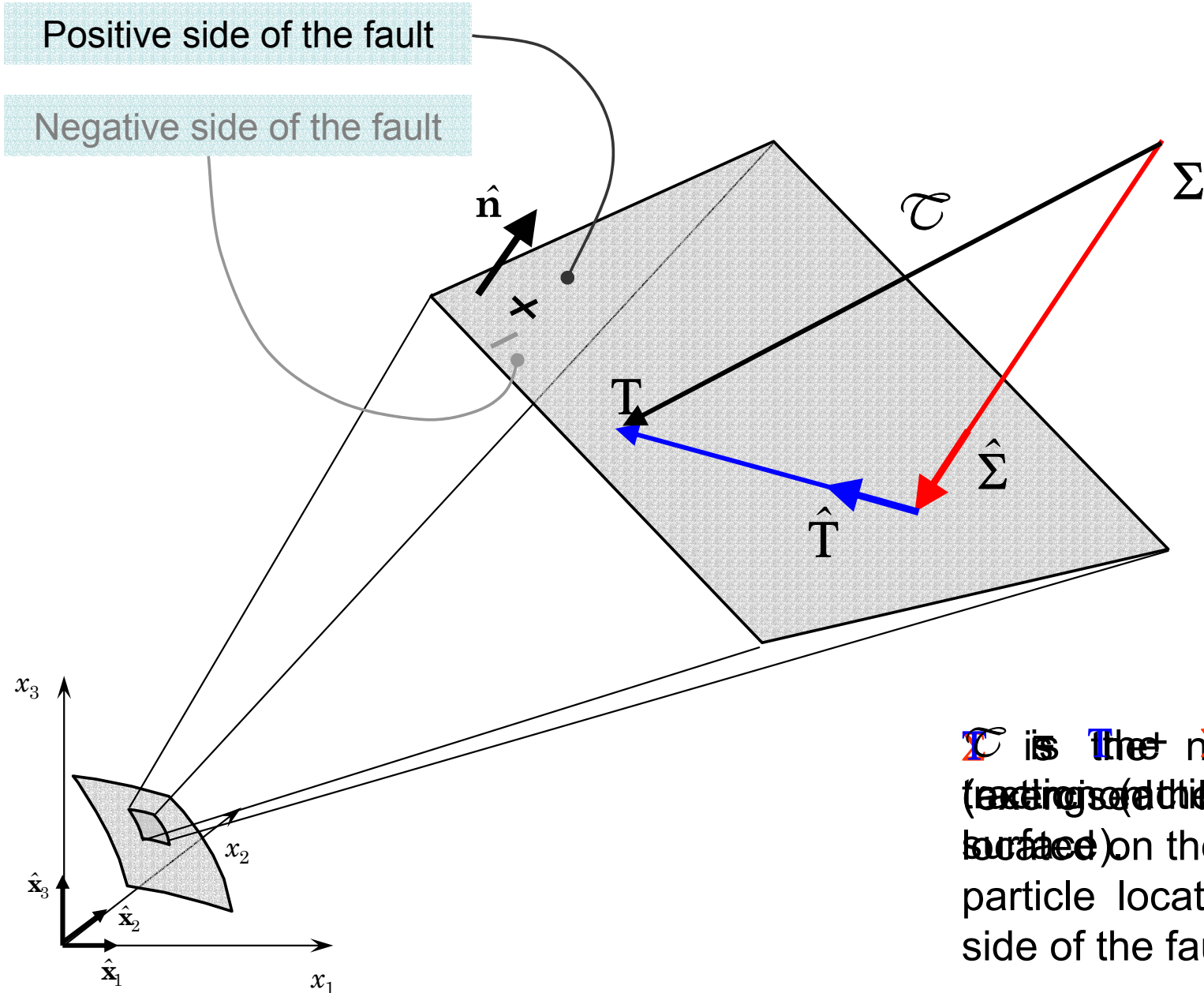


- ✓ To model seismic wave propagation in the surrounding medium;



- ✓ To predict ground shaking.

# Notations and symbols



$$\mathcal{T}^{(\hat{\mathbf{n}})} = \mathbf{T}^{(\hat{\mathbf{n}})} + \Sigma^{(\hat{\mathbf{n}})} \quad \text{total traction (acting on the fault surface).}$$

$$\mathcal{T}_j^{(\hat{\mathbf{n}})} = n_i \sigma_{ij}^{eff} \quad \text{Cauchy's formula, where } \mathcal{T}^{(\hat{\mathbf{n}})} = (\mathcal{T}_1^{(\hat{\mathbf{n}})}, \mathcal{T}_2^{(\hat{\mathbf{n}})}, \mathcal{T}_3^{(\hat{\mathbf{n}})}),$$

$$\mathbf{n} = (n_1, n_2, n_3) \text{ and}$$

$$\sigma_{ij}^{eff} = \sigma_{ij} + p_{fluid} \delta_{ij} = \begin{bmatrix} -\sigma_{n_1}^{eff} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & -\sigma_{n_2}^{eff} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & -\sigma_{n_3}^{eff} \end{bmatrix}$$

where:  $\sigma_{n_i}^{eff} = \sigma_{n_i} - p_{fluid} = -\sigma_{ii} - p_{fluid}$  and stresses are assumed to be negative for compression

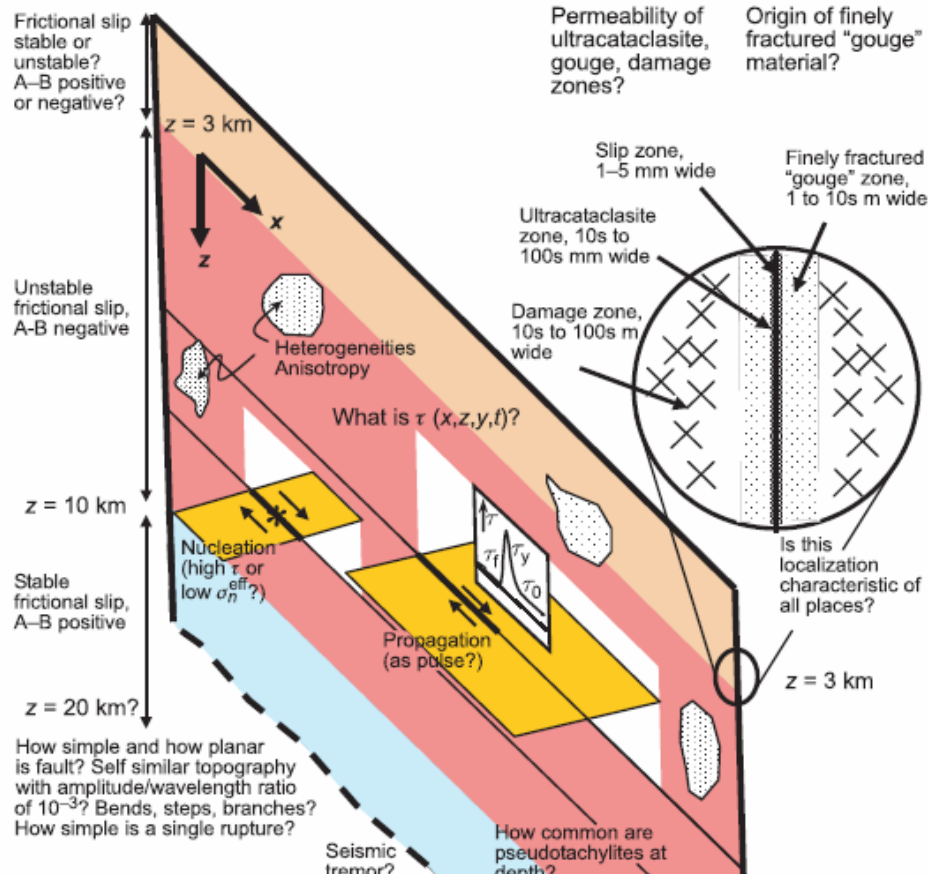
$$T_j^{(\hat{\mathbf{n}})} = n_i \sigma_{ij}^{eff} - n_j (n_i \sigma_{ik}^{eff} n_k) \quad \text{shear traction}$$

$$\Sigma_j^{(\hat{\mathbf{n}})} = n_j (n_i \sigma_{ik}^{eff} n_k) \quad \text{normal traction}$$

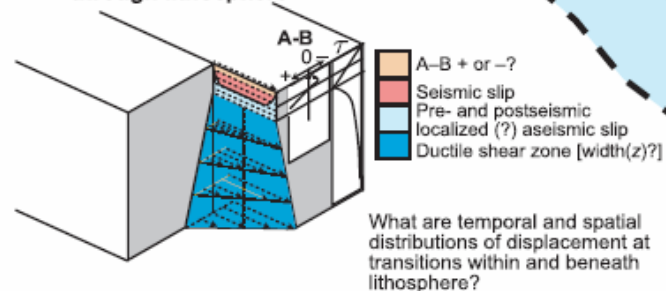


# Fault models

## (a) Seismogenic part of fault



## (b) Displacement styles through lithosphere



Tullis et al. ( 2007, MIT Press )

## Internal Structure of Principal Faults of the North Branch San Gabriel Fault

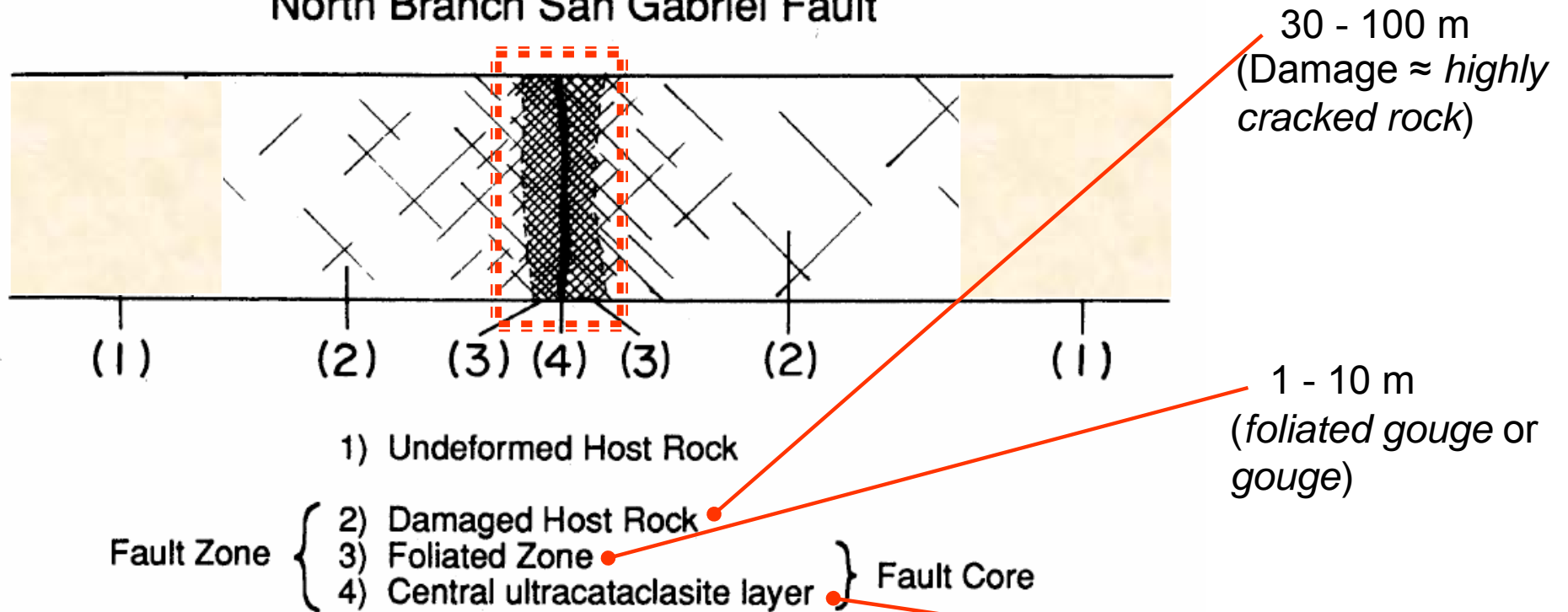
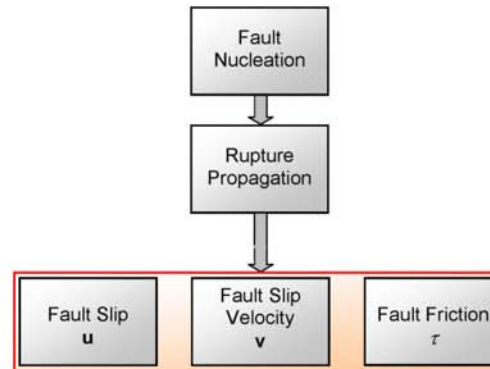
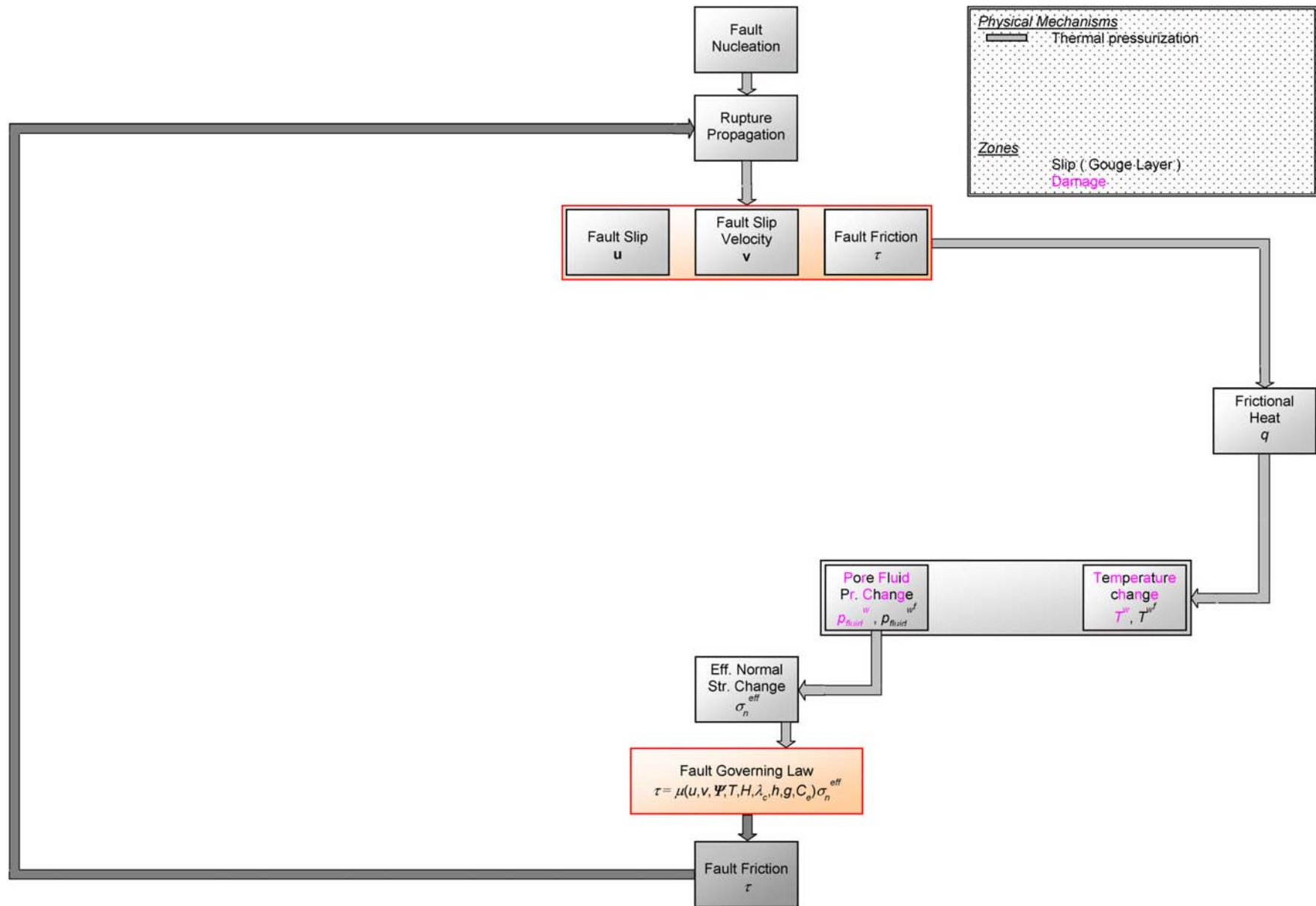


Fig. 2. Schematic section across the North Branch San Gabriel fault zone illustrating position of the structural zones of the fault. The diagram is not to scale.

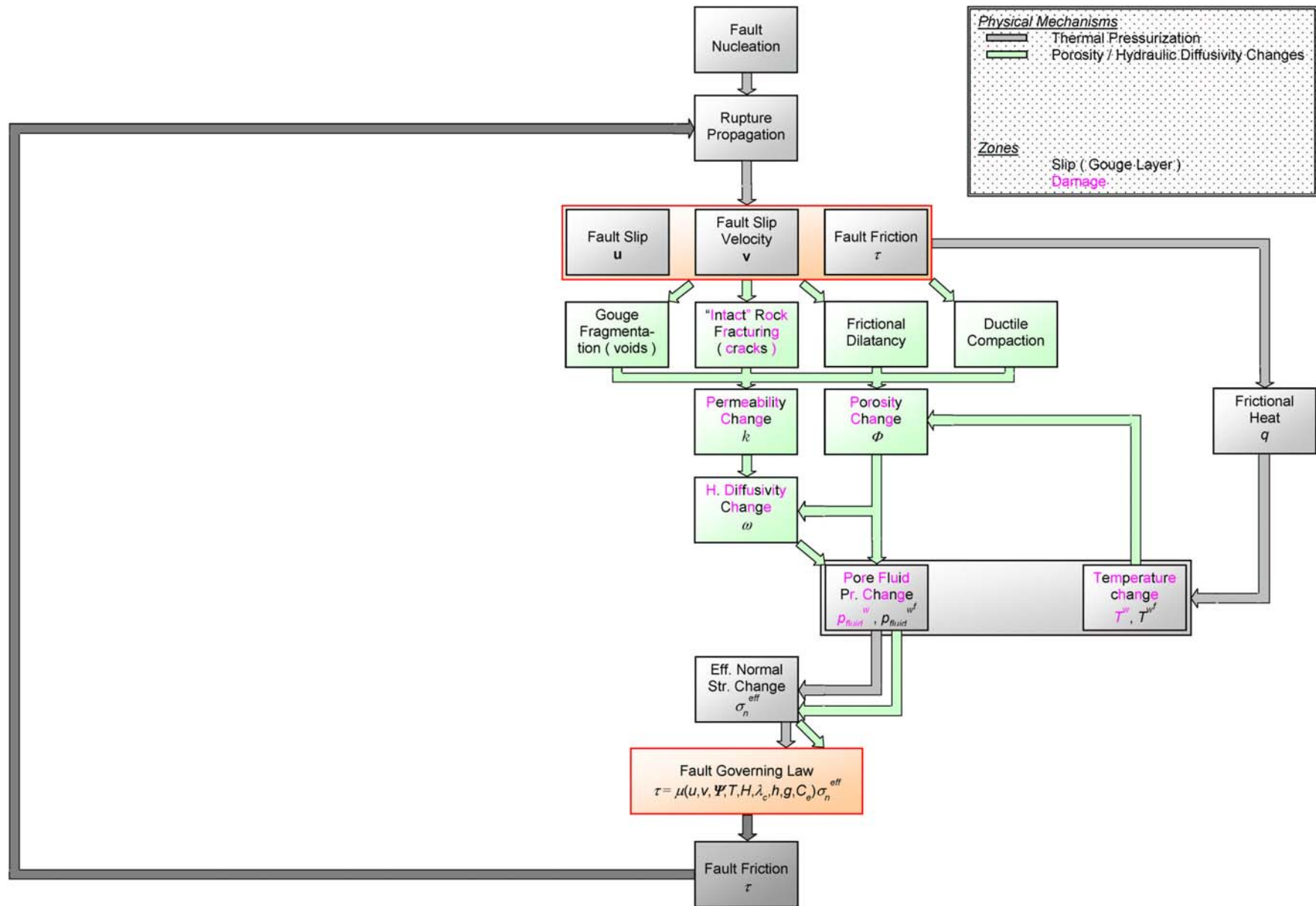
Chester, Evans and Biegel, *J. Geoph. Res.*, 1993  
 Sibson, BSSA, 2003  
 Chester and Chester, SSA, SCEC meetings 2004

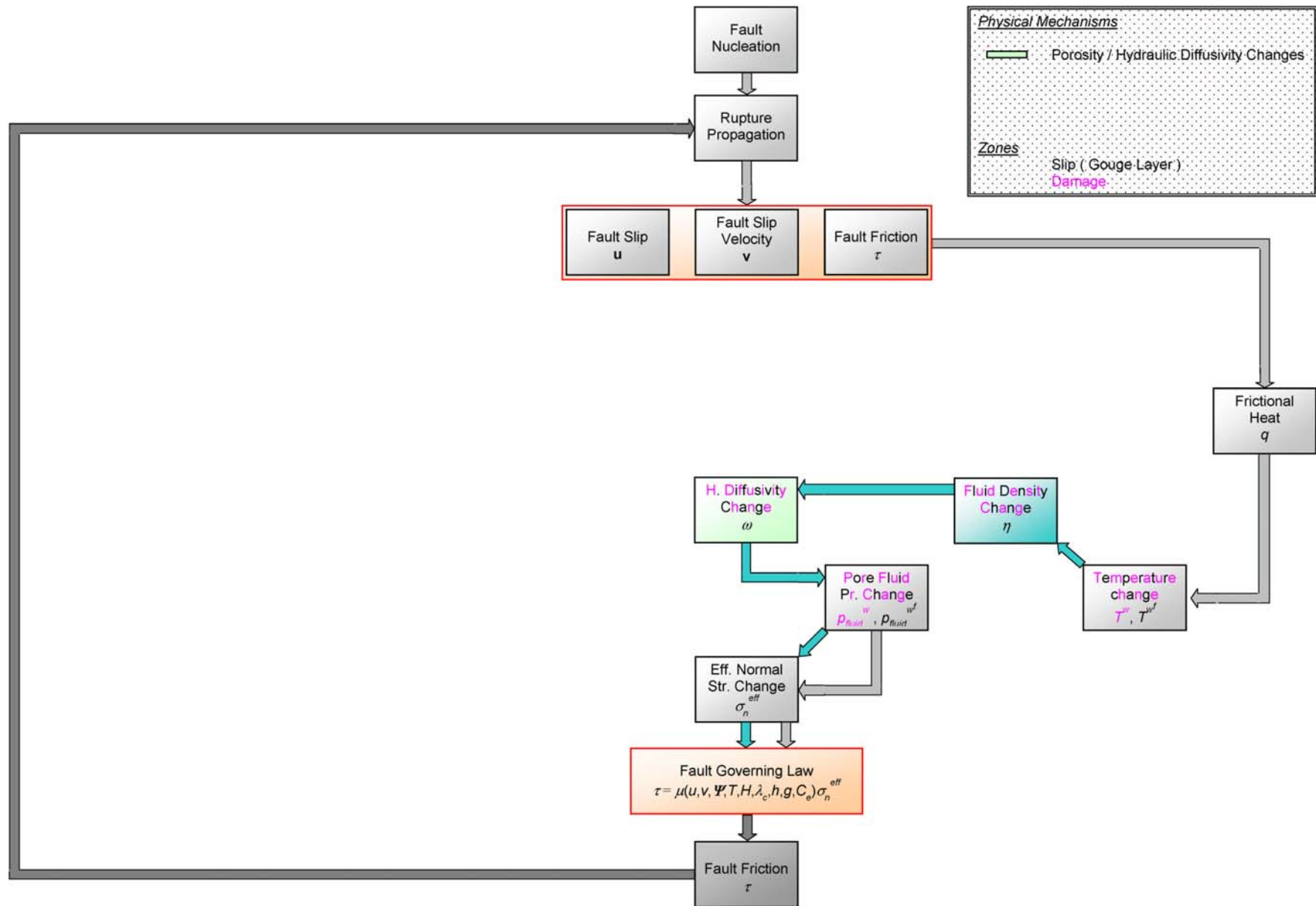
# Physical Phenomena in Faulting

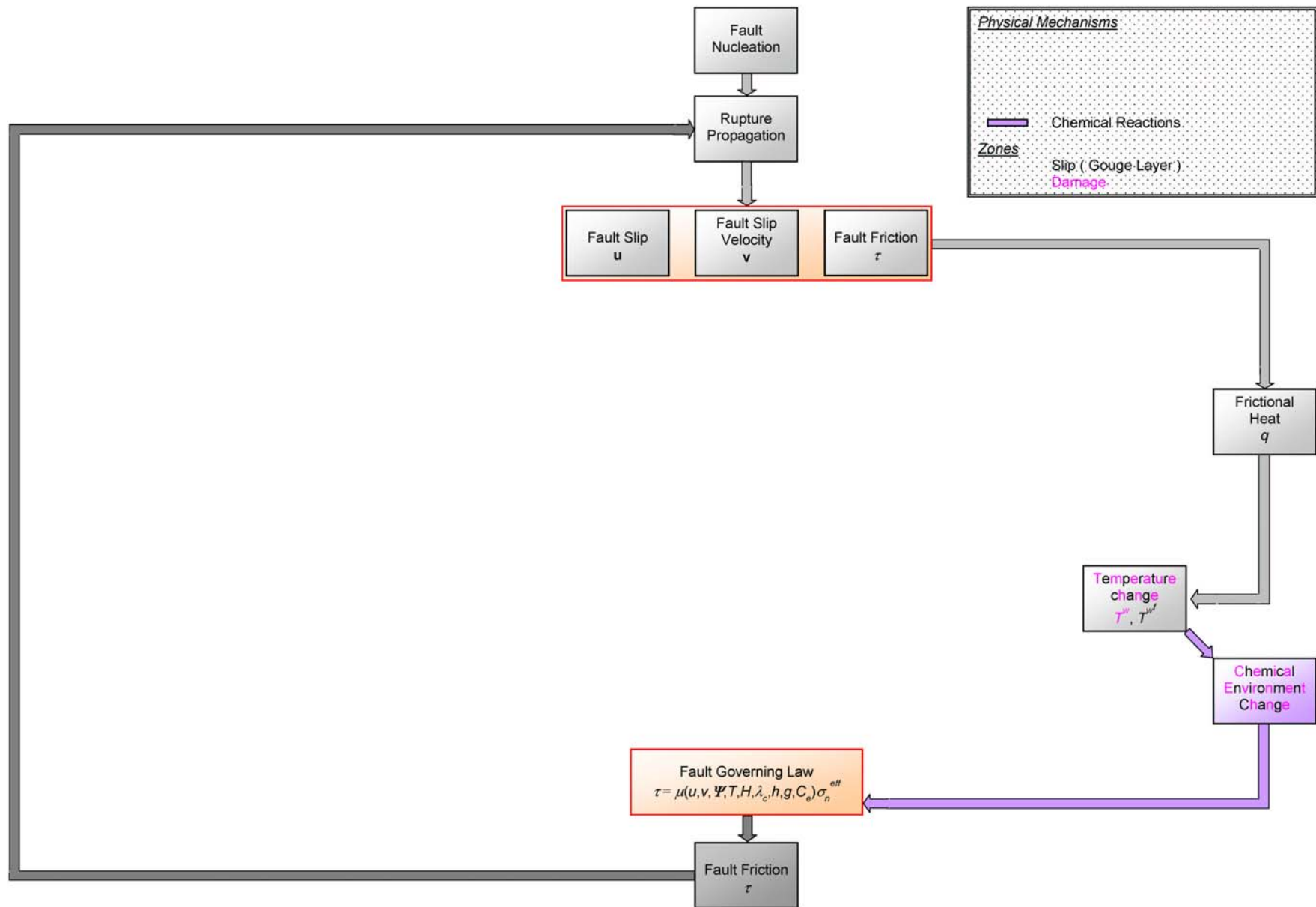


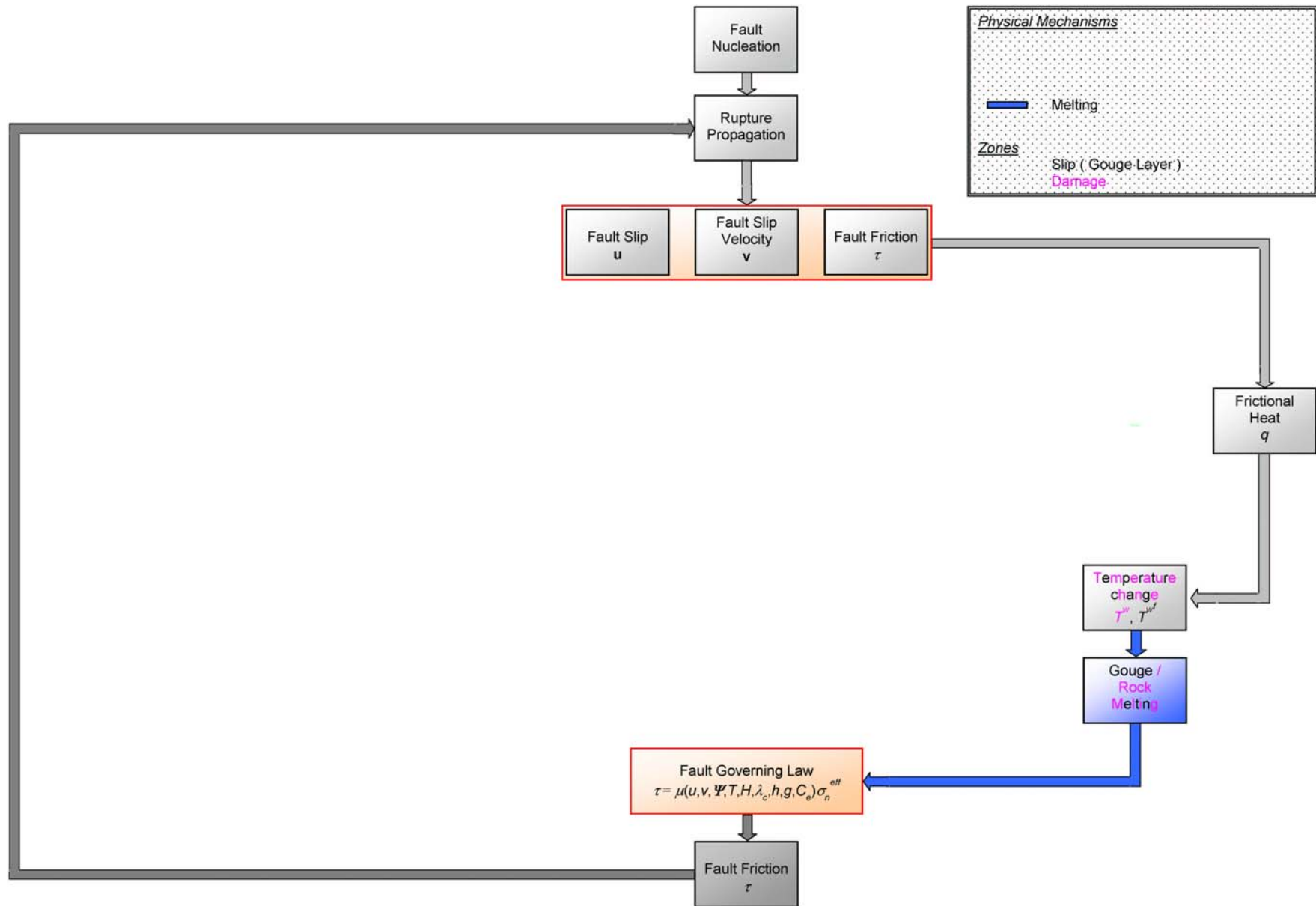




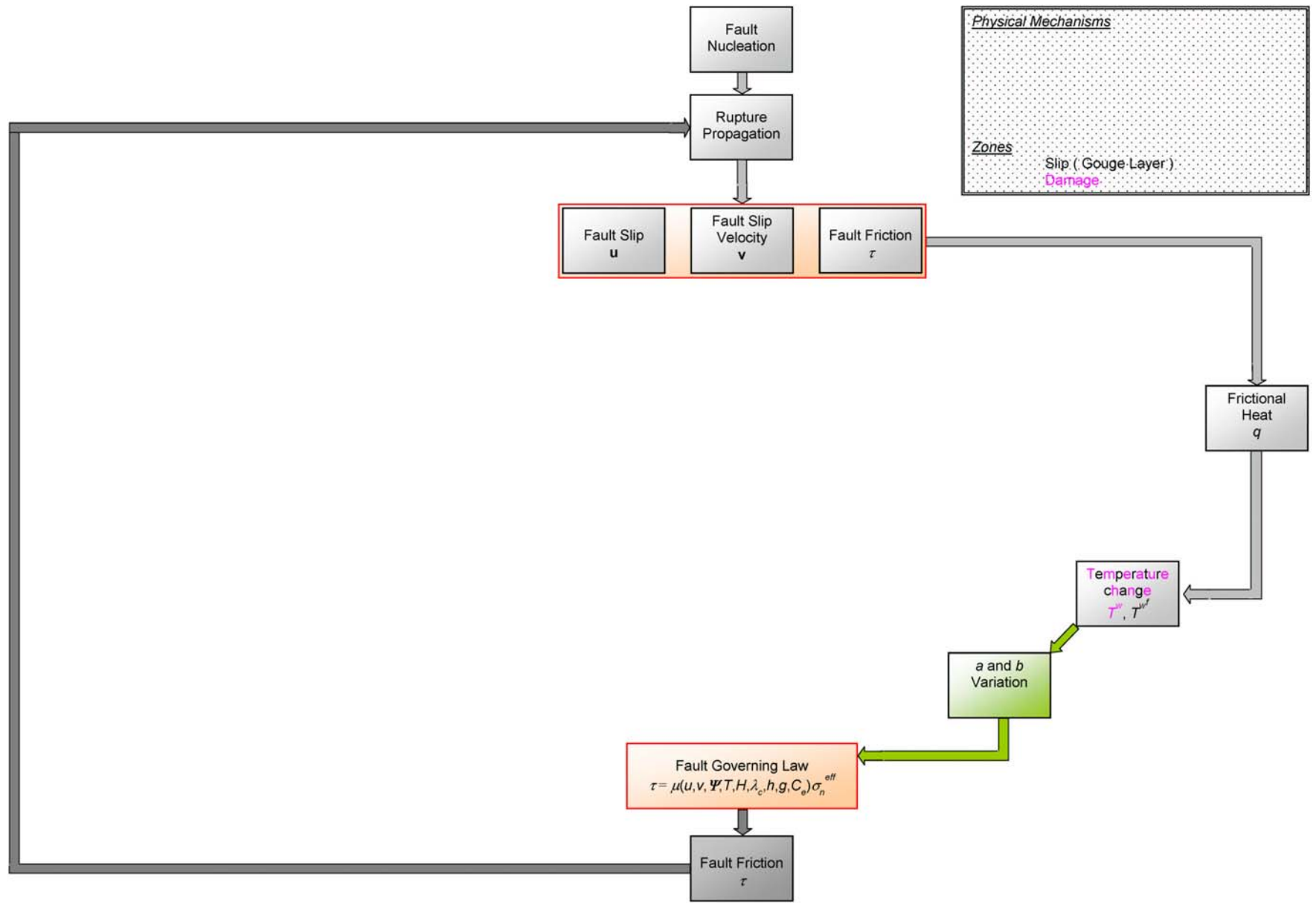


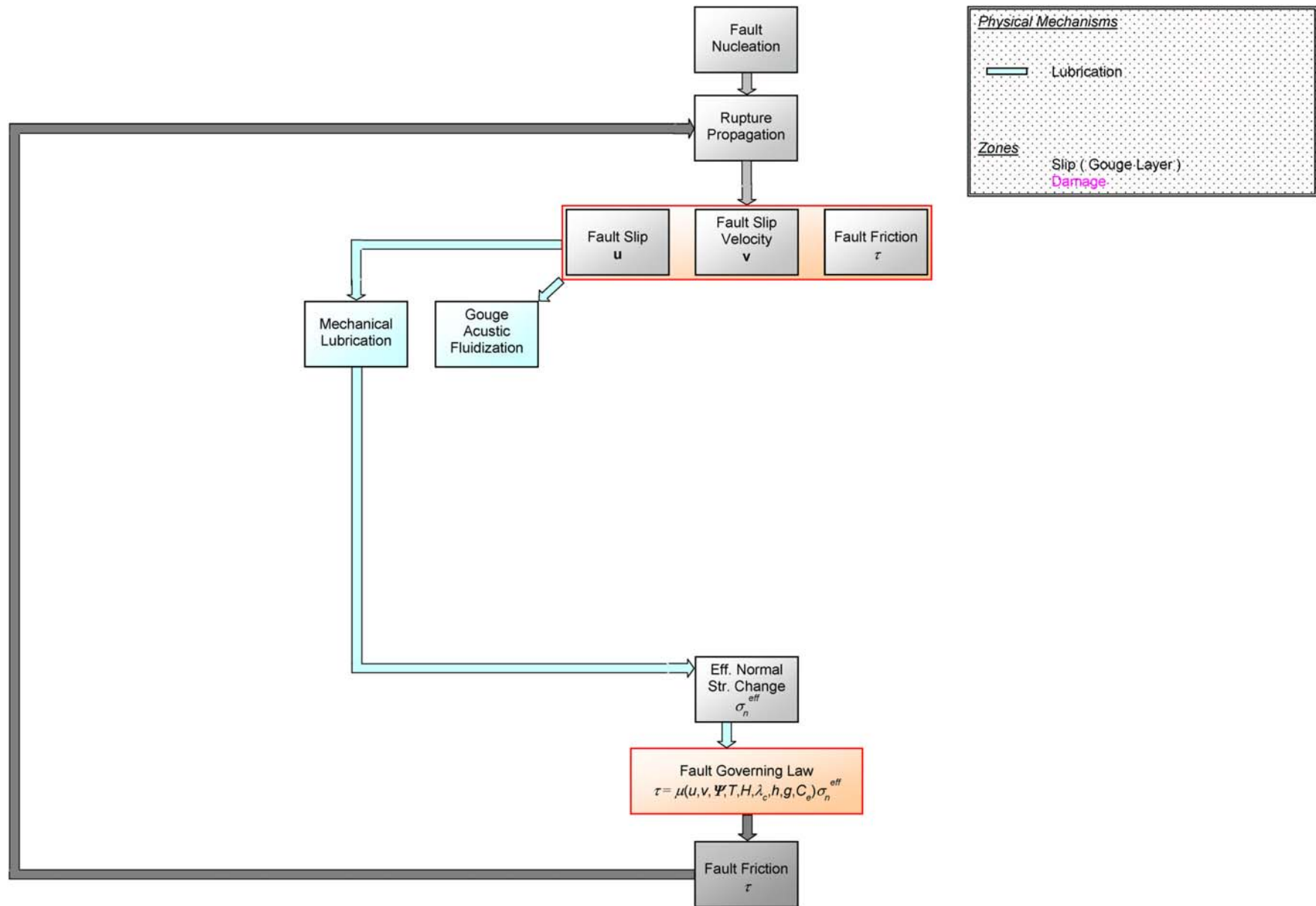


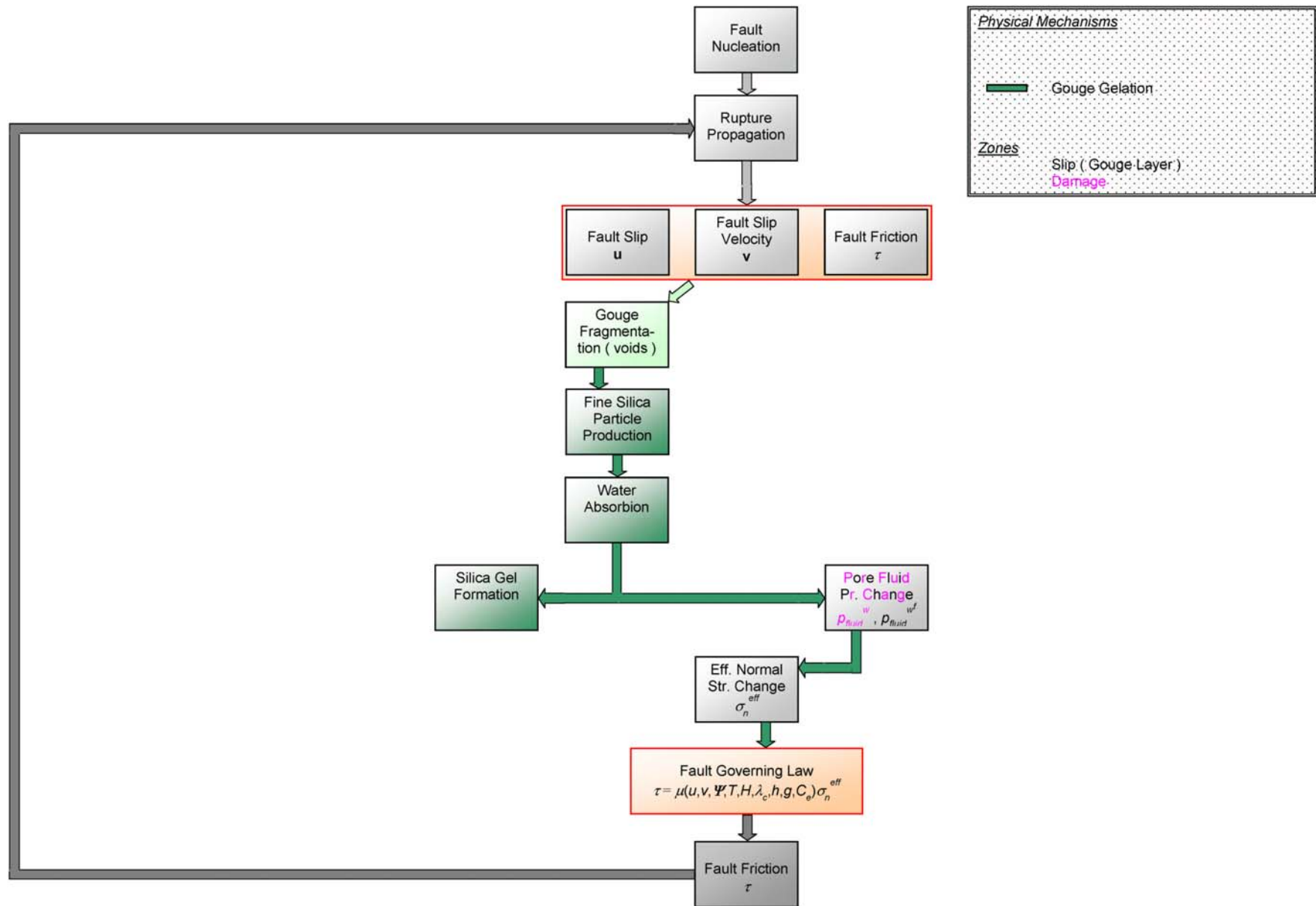


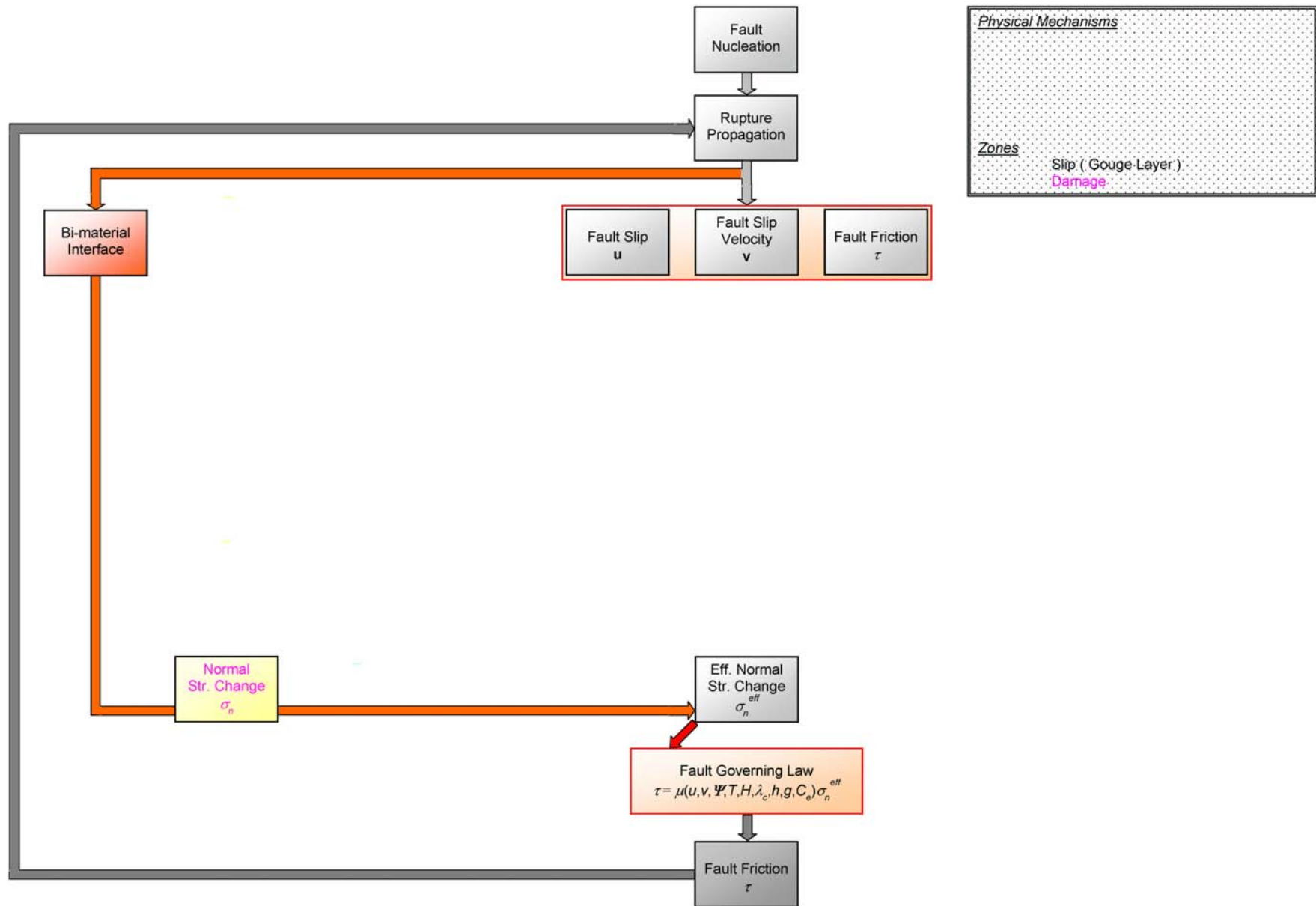




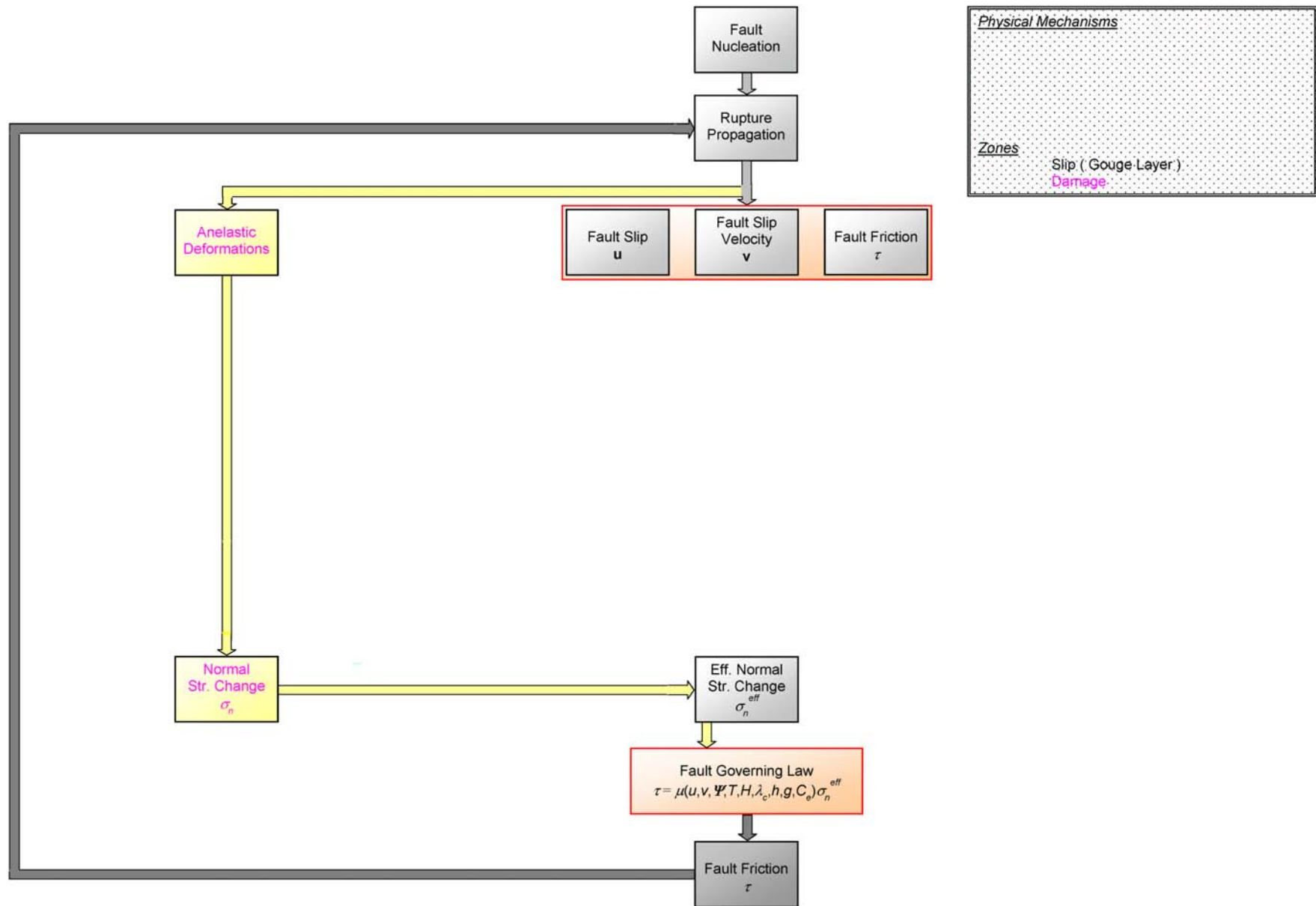


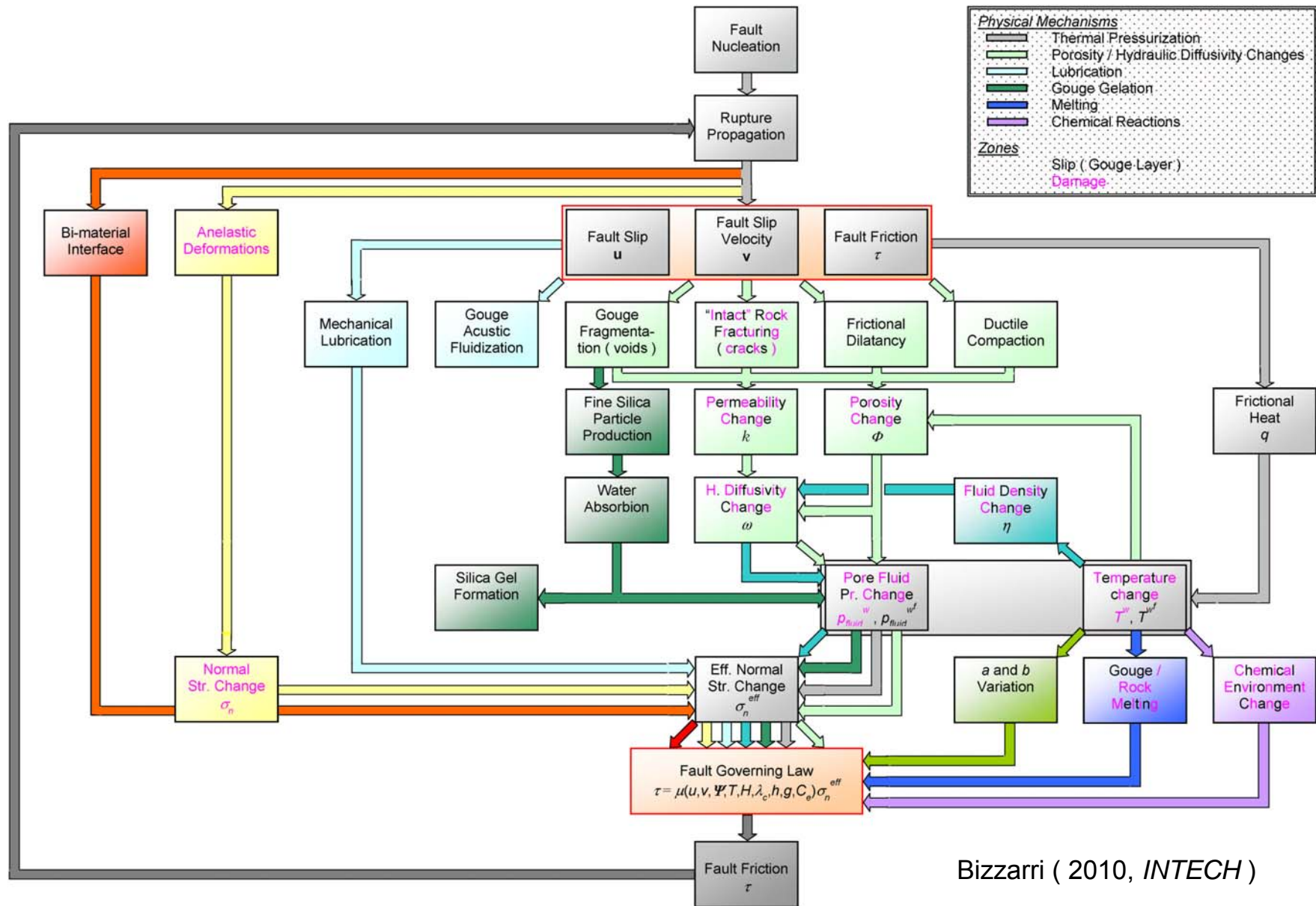












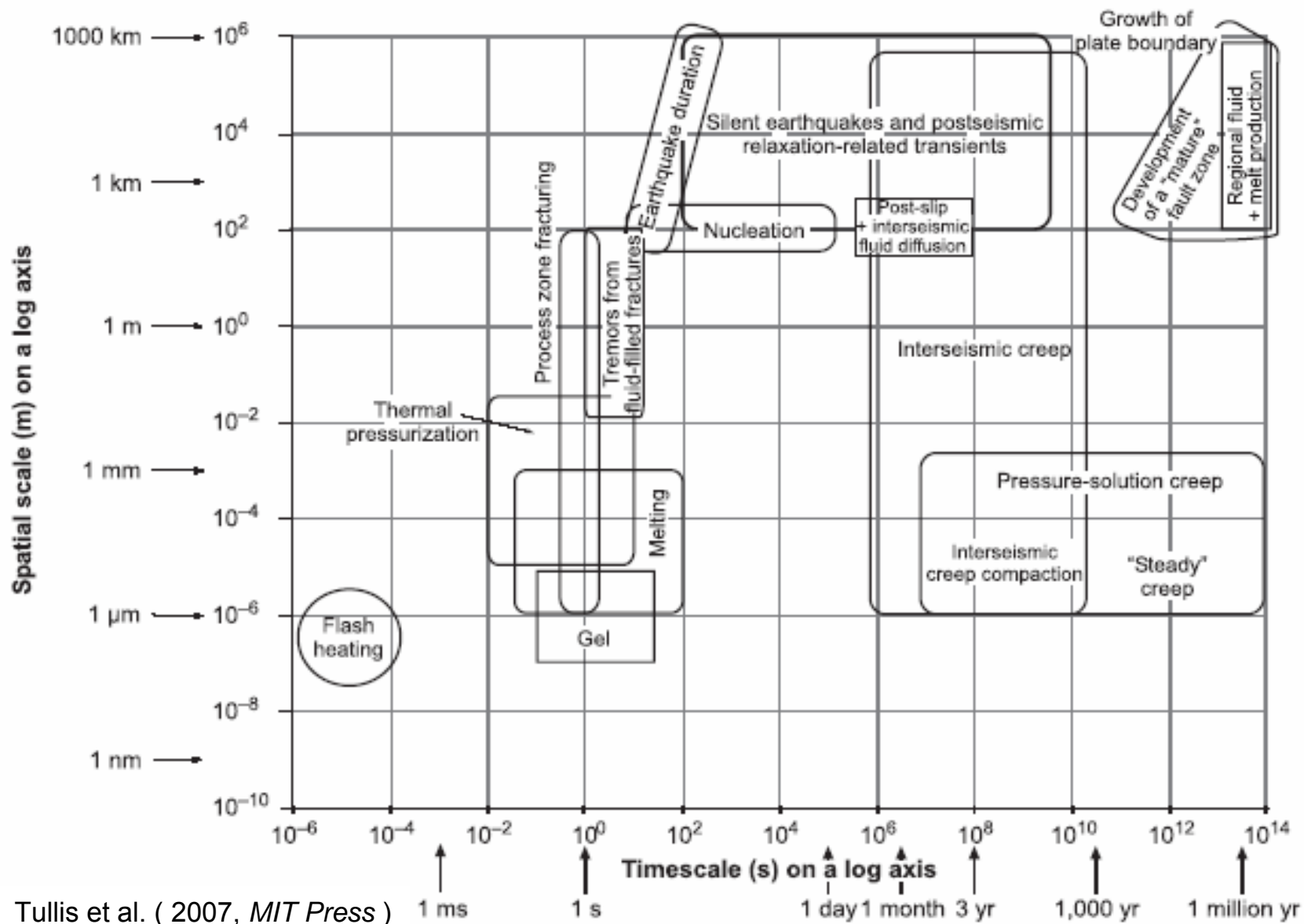
Bizzarri ( 2010, INTECH )

# Occam's razor

- ✓ We follow the logical principle of simplicity ( i.e., the Occam's razor ):

The simplest way to describe the fault complexity is to **start from the beginning** ( i.e., canonical formulations of the governing equations ) and then **add** to the model **one by one** all additional phenomena until the empirical ( instrumentally recorded ) evidence can be explained.

# Spatial and temporal scales



Tullis et al. (2007, MIT Press)





# Fracture Criteria & Constitutive Laws

## 1. FRACTURE CRITERION

- Condition that specifies, at a given fault point and at a given time, if there is a rupture or not.
- It can be expressed in terms of **energy**, in terms of **maximum frictional resistance**, and so on.
- It is based on (i) the *Benioff* ( 1951 ) hypothesis: The fracture occurs when the stress in a volume reaches the rock strength  
or, analogously,  
(ii) the *Reid* ( 1910 ) statement: The fracture takes place when the stress attains a value greater than the rock can endure.

## 2. CONSTITUTIVE LAW

- Analytical relation existing between the components of the stress tensor and physical observable(s), like the slip, the slip velocity, the state variable, etc..
- From a mathematical point of view it is a **Fault Boundary Condition ( FBC )** that controls earthquake dynamics and its complexity in space and in time.
- Its simplest form consider only **two frictional levels**,  $\tau_u$  and  $\tau_f$ ; it accounts for stress drop ( or stress release ), but the process is instantaneous: there is a singularity at crack tip. 
- **Cohesive zone models**: *Barenblatt ( 1959a, 1959b )*, *Ida ( 1972 )*, *Andrews ( 1976a, 1976b )*. In these models the singularity is removed and the stress release occurs over a breakdown zone distance  $X_b$  and in a breakdown zone time  $T_b$ . 
- Friction laws ( Rate and State dependent f. l. ): *Dieterich ( 1976 )*, *Ruina ( 1980, 1983 )*. They accounts for fault spontaneous nucleation, re – strengthening, healing, etc..

## ***CONSTITUTIVE LAW ( continues )***

- “ The central issue is *whether* faults obey simple friction laws, and if so, what is the friction coefficient associated with fault slip “ ( *Scholz and Hanks, 2004* ).

## CONSTITUTIVE LAW ( continues )

- In full of generality we can express the constitutive ( or governing ) as:

$$\tau = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{\text{eff}}(\sigma_n, p_f)$$



where:

1st – order dependencies

- $u$  is the Slip ( i. e. displ. disc. ) modulus, ←
- $v$  is the Slip Velocity modulus ( its time der. ), ←
- $\Psi = ( \Psi_1, \dots, \Psi_N )$  is the State Variable vector, ←
- $T$  is the Temperature ( accounting for Ductility, Plastic Flow, Melting and Vaporization ),
- $H$  is the Humidity,
- $\lambda_c$  is the Characteristic Length of surface ( accounting for Roughness and Topography of asperity contacts ),
- $h$  is the Hardness,
- $g$  is the Gouge ( accounting for Surface Consumption and Gouge formation ),
- $C_e$  is the Chemical Environment

# Strength & Constitutive Laws

## 1. THE STRENGTH PARAMETER

- Historically introduced by *Das and Aki* ( 1977a, 1977b ) to have a quantitative estimate of the ability to fracture for a fault
- Its expression can be generalized as:

$$S = (\mu_u \sigma_n^{eff} - \tau_0) / (\tau_0 - \mu_f \sigma_n^{eff})$$

where  $\mu$  are the friction coefficient.

- We can also define

## 2. THE FAULT STRENGTH

- Is the parameter that quantify the Strength in the more general case, in which a fault is described by a rhealistic friction laws

$$S^{fault} = \mu(u, v, \Psi, T, H, \lambda_c, h, g, C_e) \sigma_n^{eff}(\sigma_n, p_{fluid})$$

# Towards real – world conditions

$u_{tot} \sim$  several m

$v \sim$  several m/s

$\sigma_n^{eff} = 100 - 200$  MPa

Classical laboratory

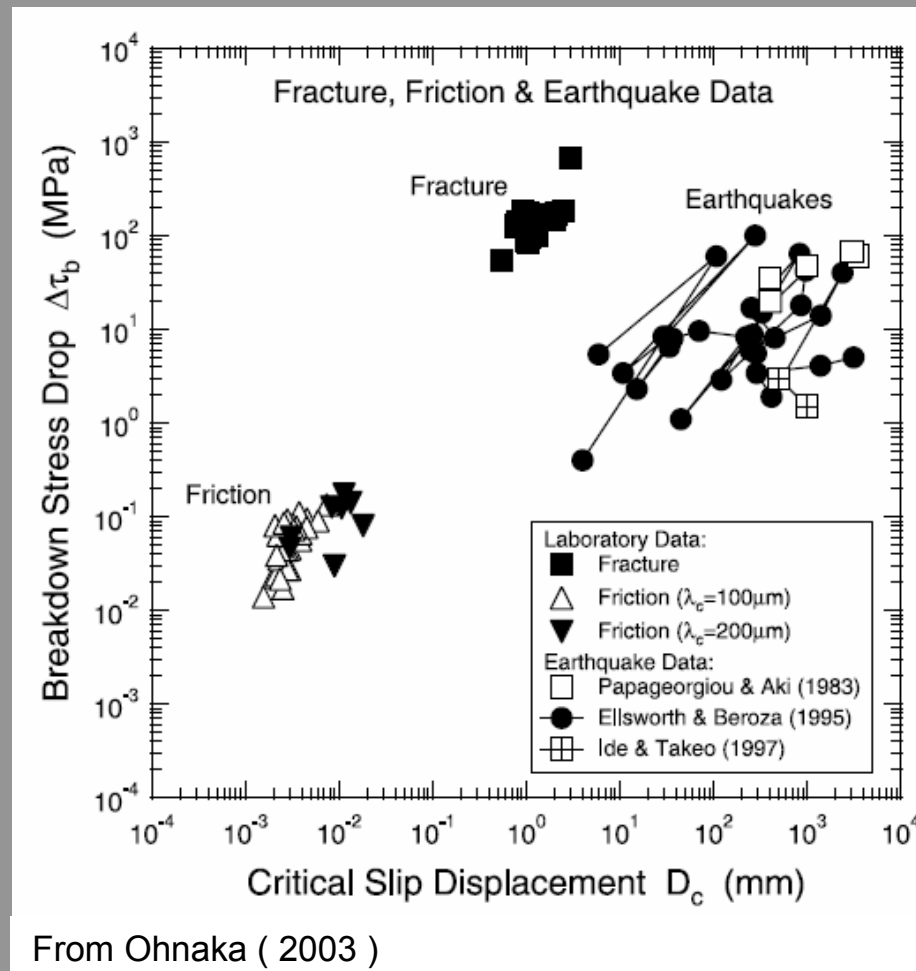
stick – slip experiments

( Dieterich, 1981 )

$u_{tot}$  up to 1.4 mm

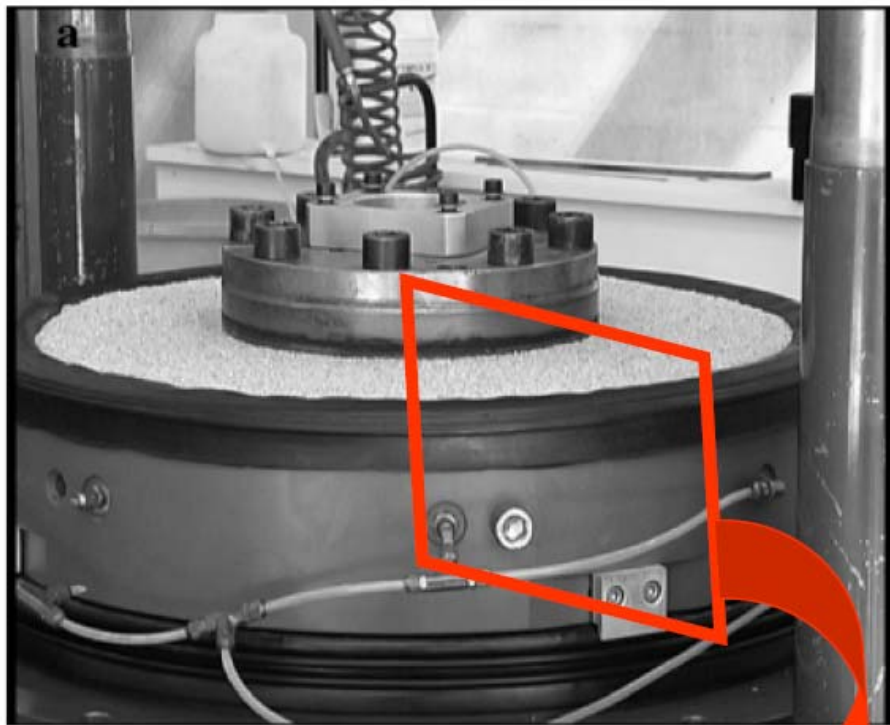
$v$  up to 25  $\mu\text{m/s}$

$\sigma_n^{eff} = 10$  MPa





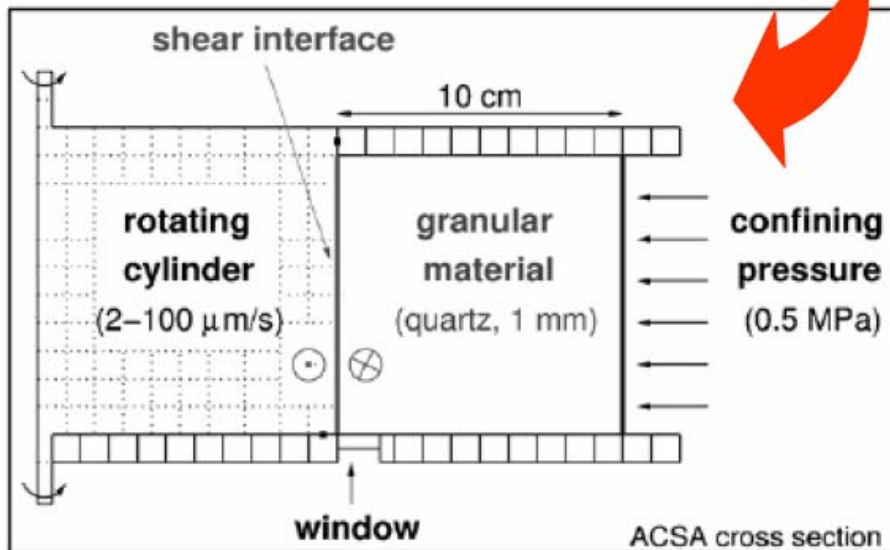
# Annular simple shear apparatus



$$u_{tot} < 50 \text{ m}$$

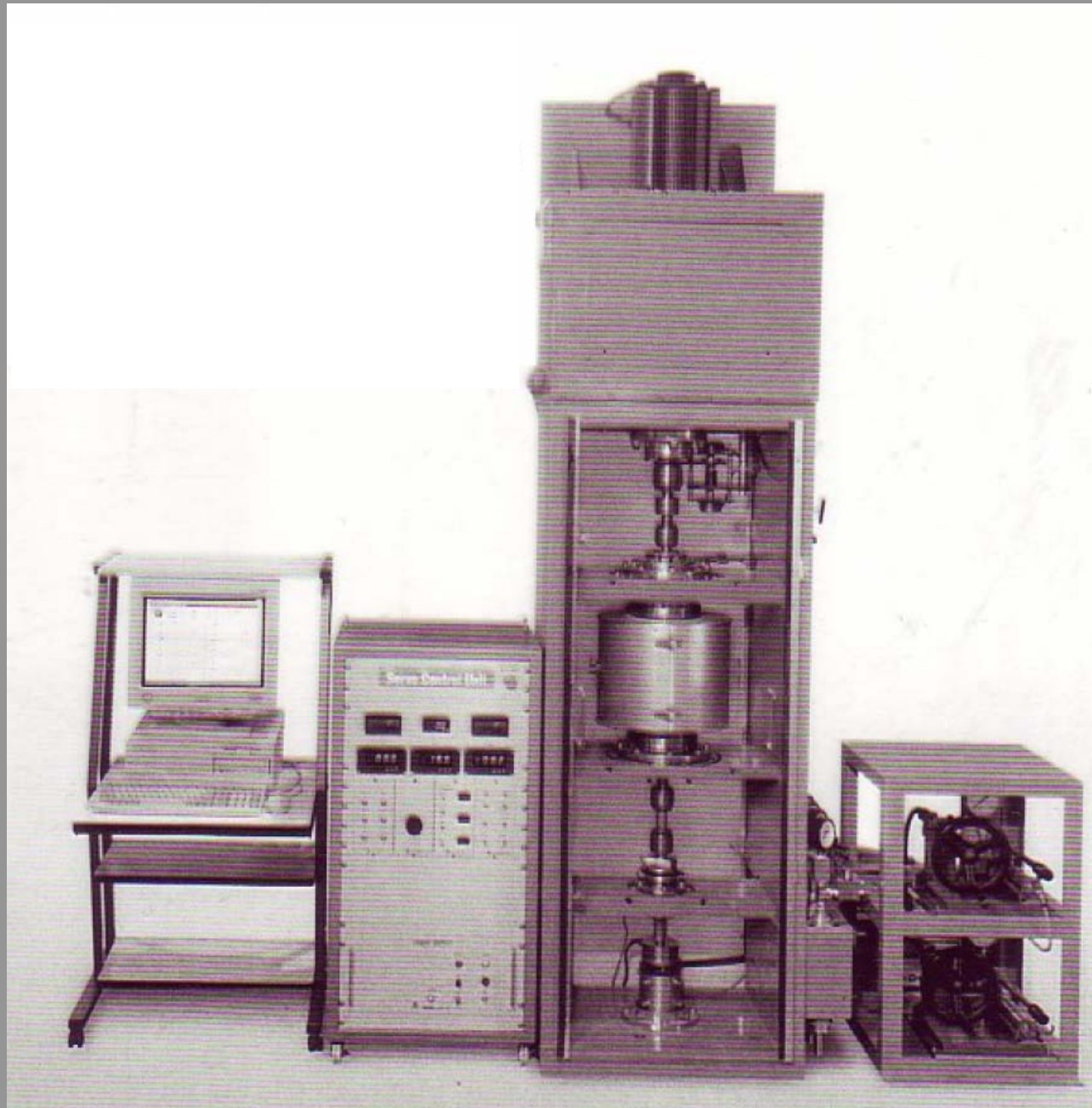
$$v = 1 \text{ } \mu\text{m/s} - 0.1 \text{ mm/s}$$

$$\sigma_n^{eff} < 1 \text{ MPa}$$



Chambon et al. ( 2006a, 2006b, *JGR*, 111, B09308, B09309 )

# High velocity rotary friction apparatus



$U_{tot} = \text{infinite}$

$v = 0.1 \mu\text{m/s} - 10 \text{ m/s}$

$\sigma_n^{eff} < 20 \text{ MPa}$

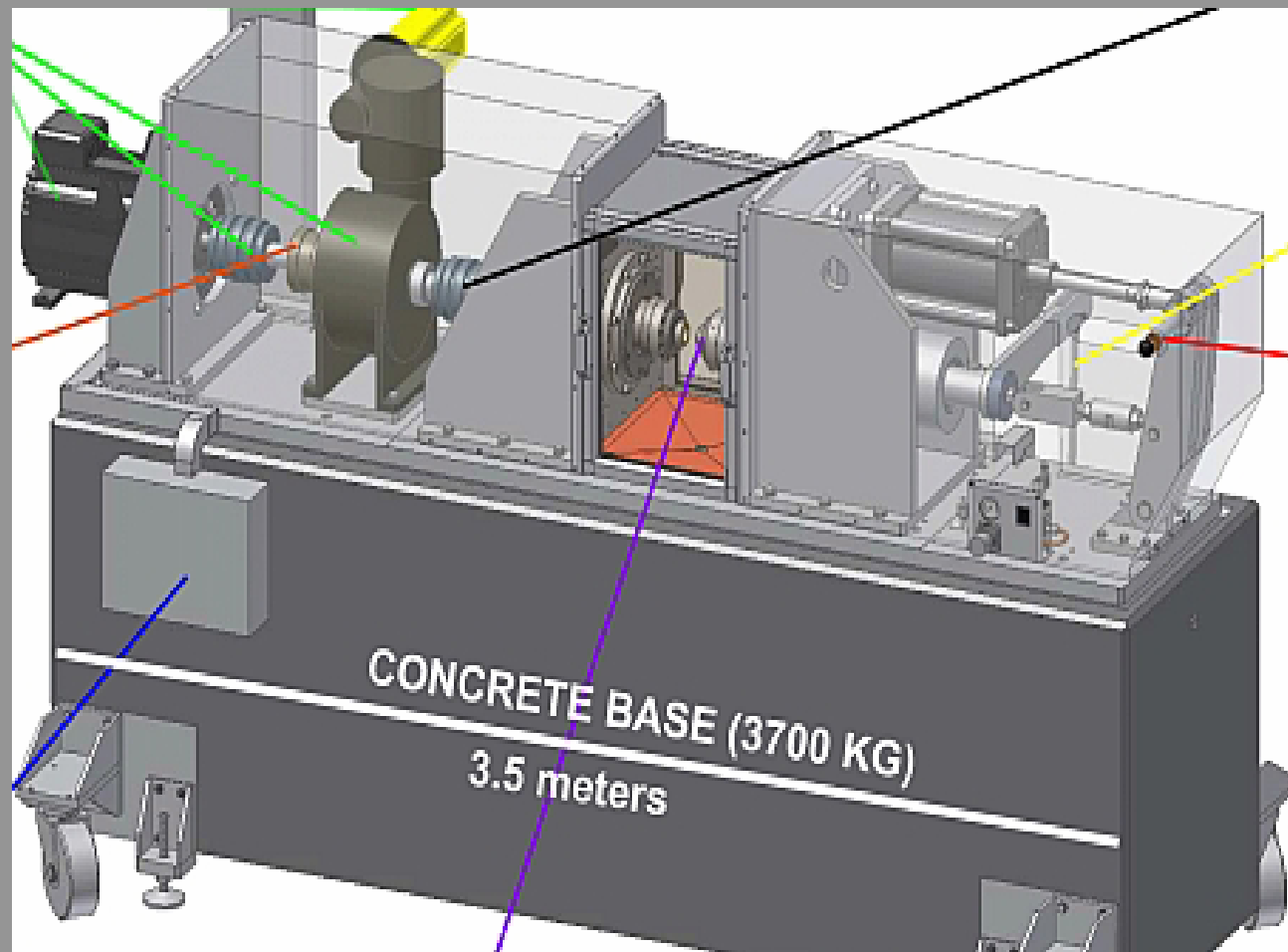
Shimamoto and Tsutumi ( 2004,  
*Str. Geol.*, **39** )

# High velocity rotary friction apparatus @ INGV

$U_{tot} = \text{infinite}$

$v = 1 \mu\text{m/s} - 9 \text{ m/s}$

$\sigma_n^{eff} < 70 \text{ MPa}$



Niemeijer et al. ( 2009, *AGU Fall Meeting* )

# Time - weakening Friction Law

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{(t - t_r)}{t_0} \right] \sigma_n^{eff} & , t - t_r < t_0 \\ \mu_f \sigma_n^{eff} & , t - t_r \geq t_0 \end{cases}$$

ilaw = 11

**TW**

$t_r = t_r(\xi)$  is the rupture onset time in every fault point  $\xi$  (when  $u > 0$ ).

Andrews ( 1985 ), Bizzarri et al. ( 2001 ) and other following Bizzarri' s papers

$t_0$  is the characteristic time – weakening duration.

# Position - weakening Friction Law

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{x}{R_0} \right] \sigma_n^{eff} & , -R_0 < x < 0 \\ \mu_f \sigma_n^{eff} & , -L < x < -R_0 \end{cases}$$

PW

$x$  is the position on the fault (extending up to  $-L$ ).

Palmer and Rice (1973)

$R_0$  is the characteristic position – weakening distance.

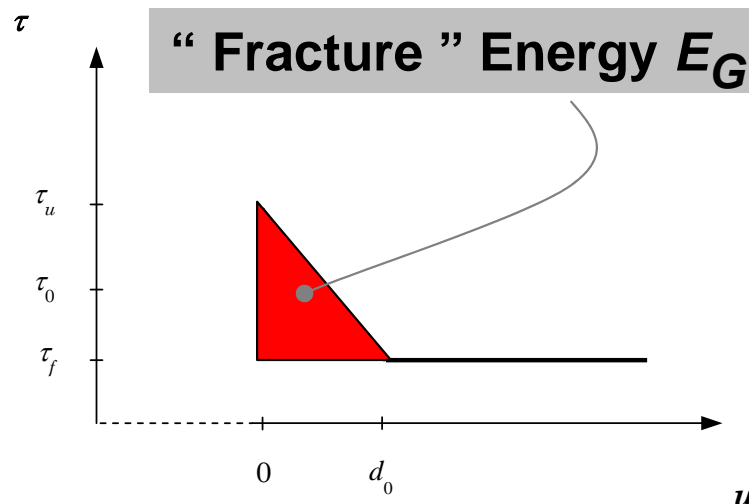
# Slip - Dependent Friction Laws

## 1. LINEAR SLIP – WEAKENING LAW

$$\tau = \begin{cases} \left[ \mu_u - (\mu_u - \mu_f) \frac{u}{d_0} \right] \sigma_n^{eff} & , u < d_0 \\ \mu_f \sigma_n^{eff} & , u \geq d_0 \end{cases}$$

ilaw = 21

SW



*Barenblatt ( 1959a, 1959b ), Ida ( 1972 ), Andrews ( 1976a, 1976b ), and many authors thereafter*

$d_0$  is the characteristic slip – weakening distance



## 2. NON – LINEAR SLIP – WEAKEING LAW

$$\tau = \begin{cases} \left[ \mu_u - \frac{\mu_u - \mu_f}{d_0} \left( u - \frac{(1 - p_{IW})d_0}{2\pi} \sin\left(\frac{2\pi u}{d_0}\right) \right) \right] \sigma_n^{eff} & , u < d_0 \\ \mu_f \sigma_n^{eff} & , u \geq d_0 \end{cases}$$

Ionescu and Campillo ( 1999 )

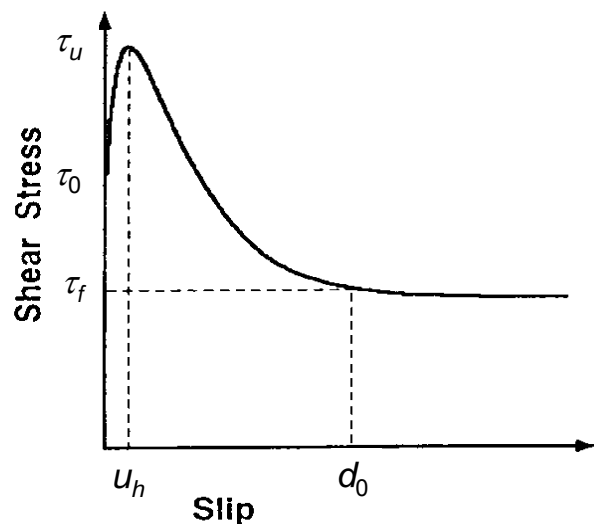
### 3. NON LINEAR SLIP – WEAKENING LAW WITH SLIP – HARDENING

$$\tau = \left\{ \left[ \left( \frac{\tau_0}{\sigma_n^{eff}} - \mu_f \right) \left( 1 + \alpha_{OY} \ln \left( 1 + \frac{u}{\beta_{OY}} \right) \right) \right] e^{-\frac{u}{d_0}} + \mu_f \right\} \sigma_n^{eff}$$

ilaw = 23

OW

$$u_h : \left. \frac{d\tau}{du} \right|_{u_h} = 0; \quad \begin{cases} u_h = r d_0 \quad (\text{e.g. } r = 0.1) \\ \tau(u_h) = \tau_u \end{cases}$$



Ohnaka and Yamashita (1989) and the following papers by Ohnaka and coworkers

$u_h$  is associated with the preparatory phase of the imminent macroscopic failure in the cohesive zone. It accounts for micro-cracking

## 4. NON LINEAR SLIP – WEAKENING LAW WITH EXPONENTIAL DECAY

$$\tau = \left[ (\mu_u - \mu_f) e^{-\frac{u}{d_0}} + \mu_f \right] \sigma_n^{eff}$$

ilaw = 24

**EW**

## 5a. POWER LAW SLIP – WEAKENING

$$\tau = \left\{ \mu_u - (\mu_u - \mu_f) \left[ \left( \frac{p_{PW}}{p_{PW} + 1} \right) \frac{u}{d_0} \right]^{p_{PW}} \right\} \sigma_n^{eff}$$

$$i_{law} = 25$$

PW

## 5b. POWER LAW SLIP – WEAKENING II

$$\tau = \left\{ \mu_f + \frac{\alpha_{CEA}}{\sigma_n^{eff}} (u - d_0)^{p_{CEA}} \right\}$$

$$\alpha_{CEA} = 5.6 \times 10^{-2} \text{ MPa m}^{p_{CEA}}$$

$$p_{CEA} = 0.4$$

Chambon et al. (2006b)

# Slip - and Rate - Dependent Friction Laws

$$\tau = \left\{ \mu^{ss}(v) + \left[ F(u)\mu_i - \mu^{ss}(v) \right] e^{\frac{\ln(0.05)u}{d_0}} \right\} \sigma_n^{eff}$$

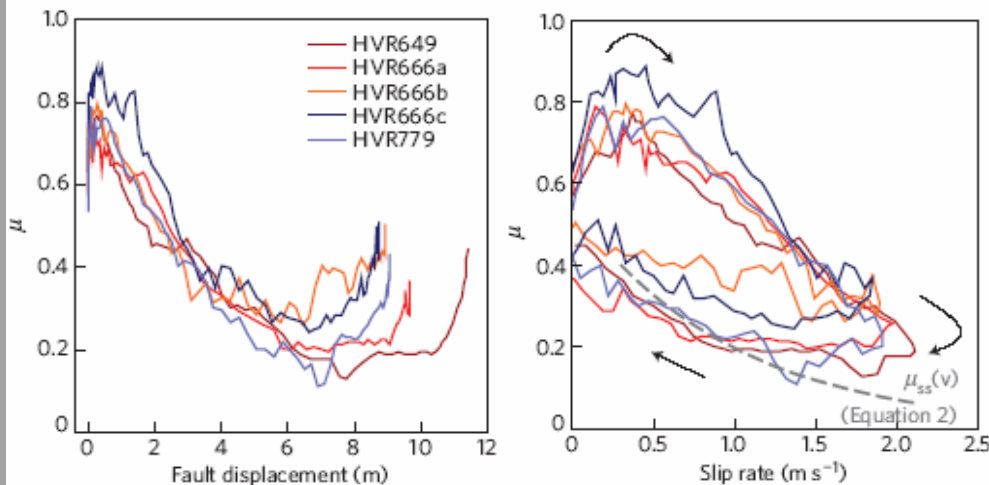
ilaw = 26

$$\mu^{ss}(v) = \mu^{ss}(0) e^{-\frac{v}{v_{SS}}}$$

$$F(u) = \alpha_{SS} + (1 - \alpha_{SS}) e^{\frac{\ln(0.05)u}{u_h}}$$

VW

## Sone and Shimamoto (2009)



$u_h$  controls the duration in slip of the slip – hardening phase, described by the function  $F(u)$ .

$$\mu^{ss}(0) = 0.55 \pm 0.09$$

$$\mu_i = 0.6$$

$$v_{SS} = 0.99 \pm 0.23 \text{ m/s}$$

$$\alpha_{SS} = 1.26 \div 1.54$$

$$u_h = 23 \div 160 \text{ mm}$$

# Rate - Dependent Friction Law

$$\tau = \frac{v_*}{v + v_*} \mu_u \sigma_n^{eff}$$

*Burrige and Knopoff ( 1967 ),  
Carlson and Langer ( 1989 ),  
Madariaga and Cochard ( 1994 ),  
Cochard and Madariaga ( 1994 )*

# Rate - and State - Dependent Friction Laws



## 1. DIETERICH IN REDUCED FORMULATION

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v_*}{v} \right) + b \ln \left( \frac{\Psi v_*}{L} \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} \end{array} \right.$$

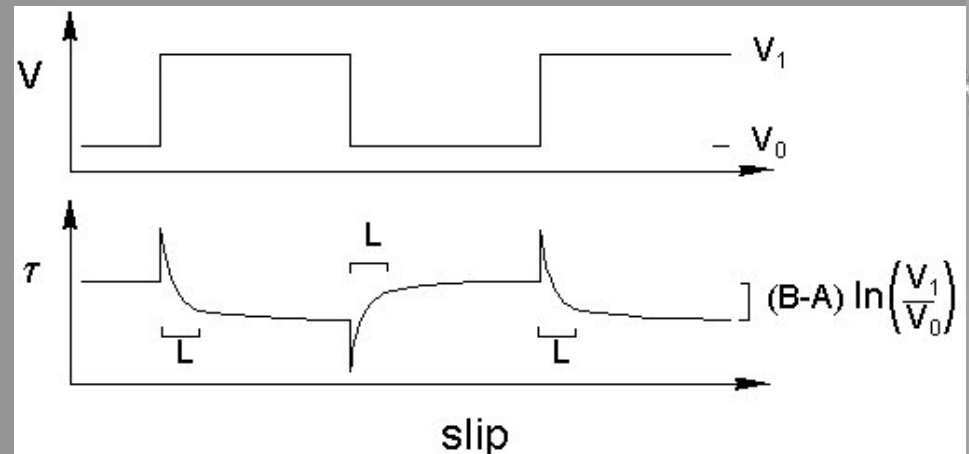
ilaw = 31

DR

However, while in velocity stepping experiments the traction response following the velocity variation is directly controlled by the parameter  $L$ , its effects are much less evident during the dynamic rupture propagation.

Bizzarri and Cocco (2005)

**Response to an abrupt jump in load**





## 2. RUINA – DIETERICH ( RUINA ORIGINAL FORM )

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v_*}{v} \right) + \theta \right] \sigma_n^{eff} \\ \frac{d}{dt} \theta = - \frac{v}{L} \left[ \theta + b \ln \left( \frac{v}{v_*} \right) \right] \end{array} \right.$$

Ruina ( 1980, 1983 )

## **2bis. RUINA – DIETERICH ( RUINA MODERN FORM. )**

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{\Psi v}{L} \ln\left(\frac{\Psi v}{L}\right) \end{array} \right.$$

ilaw = 32

RD

Beeler et al. ( 1994 ), Roy and Marone ( 1996 )

### 3. DIETERICH – RUINA WITH VARYING NORMAL STR.

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} - \left( \frac{\alpha_{LD} \Psi}{b \sigma_n^{eff}} \right) \frac{d}{dt} \sigma_n^{eff} \end{array} \right.$$

ilaw = 31

decis10=T

**DR**

Linker and Dieterich ( 1992 ), Dieterich and Linker ( 1992), Bizzarri and Cocco ( 2006a, 2006b )

## 4. RUINA – DIETERICH WITH VARYING NORMAL STR.

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{\Psi v}{L} \ln\left(\frac{\Psi v}{L}\right) - \left(\frac{\alpha_{LD} \Psi}{b \sigma_n^{eff}}\right) \frac{d}{dt} \sigma_n^{eff} \end{array} \right.$$

ilaw = 32

decis10=T

**RD**

Linker and Dieterich ( 1992 ), Bizzarri and Cocco ( 2006a, 2006b )

## 5. DIETERICH IN REDUCED FORM REGULARIZED

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v + v_*}{v + v_r} \right) + b \ln \left( \frac{\Psi(v - v_r)}{L} + 1 \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi(v - v_r)}{L} \end{array} \right.$$

ilaw = 33

**DE**

$v_r$  is a regularization fault slip velocity

Perrin et al. (1995), Cocco et al. (2004)

## 6. RUINA REGULARIZED

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - \alpha \ln \left( \frac{v_* - v_r}{v + v_r} \right) + \frac{\Psi}{\sigma_n^{eff}} \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = -\frac{v + v_r}{L} \left( \Psi + b \ln \left( \frac{v + v_r}{v_* - v_r} \right) \right) \end{array} \right.$$

ilaw = 34

**RE**

$v_r$  is a regularization fault slip velocity

Bizzarri (2002, unpublished work)

## 7. DIETERICH IN REDUCED FORM WITH HEALING

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v^*}{v} + 1 \right) + b \ln \left( \frac{\Psi v^*}{L} + 1 \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \Psi = \frac{\gamma_{fh} - \Psi}{t_{fh}} - \frac{\Psi v}{L} \end{array} \right.$$

ilaw = 35

DH

$$\gamma_{fh} = 1 \text{ s}$$

$t_{fh}$  is the time for healing (slip duration)

Evolution law proposed by Nielsen et al. (2000) and by Nielsen and Carlson (2000). Used in this form by Cocco et al. (2004)



## 8. DIETERICH IN REDUCED FORM WITH 2 STATE VAR.

ilaw = 36

DW

Tullis and Weeks ( 1993 ). Used in this form by *Bizzarri ( xxxx, unpublished work )*

## 9. PRAKASH – CLIFTON

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln\left(\frac{v_*}{v}\right) + b \ln\left(\frac{\Psi v_*}{L}\right) \right] \left( \frac{d}{dt} \Psi_1 + \frac{d}{dt} \Psi_2 \right) \\ \frac{d}{dt} \Psi = 1 - \frac{\Psi v}{L} \\ \frac{d}{dt} \Psi_1 = -\frac{v}{L_1} \left( \Psi_1 - \alpha_{PC_1} \sigma_n^{eff} \right) \\ \frac{d}{dt} \Psi_2 = -\frac{v}{L_2} \left( \Psi_2 - \alpha_{PC_2} \sigma_n^{eff} \right) \end{array} \right.$$

ilaw = 37

PC

$\Psi_1$  and  $\Psi_2$  are additional state variables accounting for the coupling with effective normal stress. The formulation of friction law is not based on the Amonton – Coulamb law.

Coupling with effective normal stress proposed by Prakash and Clifton (1993) and Prakash (1998). Used in this form by Bizzarri (2005, unpublished work)

## 10. RUINA – DIETERICH WITH FLASH HEATING

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v_*}{v} \right) + \theta \right] \sigma_n^{eff} \\ \frac{d}{dt} \theta = \begin{cases} -\frac{v}{L} \left[ \theta + b \ln \left( \frac{v}{v_*} \right) \right] & , v \leq v_{fh} \\ -\frac{v}{L} \left[ \theta + b \frac{v_{fh}}{v} \ln \left( \frac{v}{v_*} \right) + \left( 1 - \frac{v_{fh}}{v} \right) \left( a \ln \left( \frac{v}{v_*} \right) + \mu_* - \mu_{fh} \right) \right] & , v > v_{fh} \end{cases} \end{array} \right.$$

ilaw = 38

**FH**

where  $v_{fh} = \frac{\pi\chi}{D_{ac}} \left( c \frac{T_{weak} - T^{wf}}{\tau_{ac}} \right)^2$  is a weakening velocity above which flash heating is activated,  $T^{weak}$  is a weakening temperature,  $\tau_{ac}$  is the ( average ) shear strength of asperity contacts and  $D_{ac}$  their ( average ) size.

*Beeler and Tullis ( 2003 ); Tullis and Goldsby ( 2003a, 2003b ). Rice ( 1999, 2006). Modified from Noda et al. ( 2009 )*

# 11. RUINA – DIETERICH WITH TEMPERATURE DEPENDEN.

$$\left\{ \begin{array}{l} \tau = \left[ \mu_* - a \ln \left( \frac{v_*}{v} \right) + \theta \frac{a Q_a}{R} \left( \frac{1}{T} - \frac{1}{T_*} \right) \right] \sigma_n^{eff} \\ \frac{d}{dt} \theta = -\frac{v}{L} \left[ \theta + b \ln \left( \frac{v}{v_*} \right) + \frac{b Q_b}{R} \left( \frac{1}{T} - \frac{1}{T_*} \right) \right] \end{array} \right.$$

ilaw = 39

CH

where  $Q_a$  and  $Q_b$  are activation energies ( Kato, 2001 assumes:  $Q_a = Q_b = 0.1$  MJ/mol ) and  $T_*$  is a reference absolute temperature.

Note that  $T$  is the absolute temperature.

Chester and Higgs ( 1992 ), Kato ( 2001 )

# Slip - and State - Dependent Friction Law

$$\tau = \begin{cases} \left[ (\mu_u - \Delta\mu) \left( 1 - \frac{u}{d_1} \right) \right] \sigma_n^{eff} & , u < d_1, \Psi \geq \Psi_1 \\ 0 & , u \geq d_1, \Psi \geq \Psi_0 \\ \mu_{sp} \left( 1 - \frac{\Psi}{\Psi_0} \right) \sigma_n^{eff} & , \Psi < \Psi_0, \Psi < \Psi_1 \end{cases}$$

$$\frac{d}{dt} \Psi = - \frac{\beta_{CM}}{d_0} (\Psi - v)$$

ilaw = 41

CM

$\Delta\mu$  is an initial artificial stress drop

$$\Psi_1 \equiv \Psi_0 (u - u_1) / (d_1 - u_1)$$

$$U_1 \equiv - d_1 (\mu_{sp} - \mu_u + \Delta\mu) / (\mu_u - \Delta\mu)$$

$d_0$  and  $d_1$  are characteristic lengths

$\mu_{sp} = 0 \Rightarrow$  linear SW with  $d_1$  as characteristic length

Cochard and Madariaga ( 1994 )

# Free Volume Friction law

$$\left\{ \begin{array}{l} \tau = \sigma_d \operatorname{Arcsinh} \left( \frac{v}{v_*} \frac{e^{\frac{f_* + \chi_s + \chi_h}{\chi}}}{1 - m_0} \right) \\ \frac{d}{dt} \chi = -R_c e^{-\frac{\chi_c}{\chi}} + \alpha_{FV} \tau v \\ m_0 = \begin{cases} 1 & , \tau \leq \tau_0 e^{\frac{\chi_h}{\chi}} \\ \frac{\tau_0}{\tau} e^{\frac{\chi_h}{\chi}} & , \tau > \tau_0 e^{\frac{\chi_h}{\chi}} \end{cases} \end{array} \right.$$

ilaw = 51

**FV**

$\chi \equiv \Phi - \Phi_0$  free volume variable

$\chi_s$  reference value of  $\chi$  for shearing

$\chi_h$  FV value required to create a Shear Transformation Zone ( STZ )

$\chi_c$  FV value for compaction

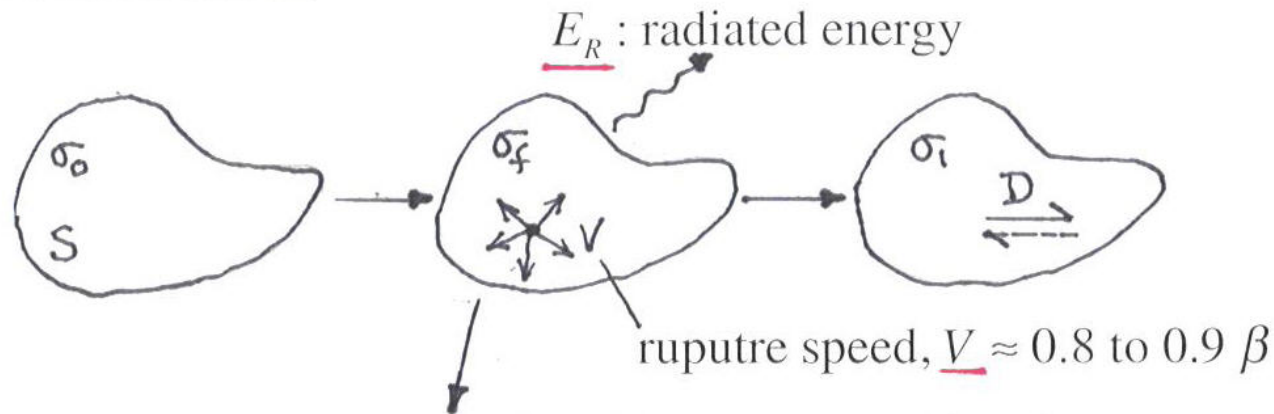
$R_c$  rate of compaction

$\alpha_{FV}$  scaled dilatancy coefficient

*Falk and Langer ( 1998, 2000 );  
Lemaitre ( 2002 ); Daub and Carlson  
( 2008 )*

# How to relate relevant quantities to constitutive parameters

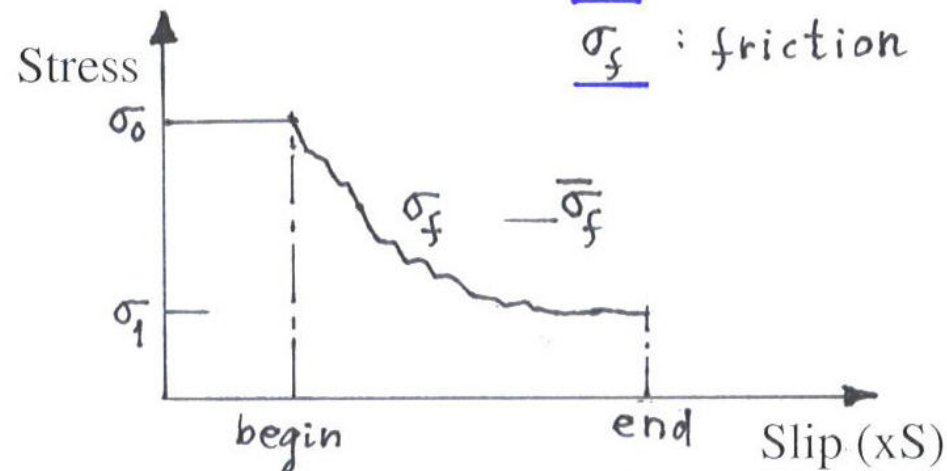
## Dynamic Parameters



$E_{NR} = E_F + E_G + \dots$  : non-radiated energy

$E_F$  : friction (heat),  $E_G$  : fracture energy

$\sigma_f$  : friction

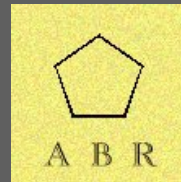


$\Delta\sigma_s = \sigma_0 - \sigma_1$  : static stress drop

$\Delta\sigma_d = \sigma_0 - \bar{\sigma}_f$  : dynamic stress drop



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# **Support Slides: Parameters, Notes, etc.**

*To not be displayed directly. Referenced above.*

## **Thermal pressurization:**

*Sibson ( 1973 ); Lachenbruch ( 1980 ); Mase and Smith ( 1985, 1987 );  
Andrews ( 2002 ); Bizzarri and Cocco ( 2006b, 2006c ) .*

*Morrow et al. ( 1984 ) show that gouge contains water*

## **Gouge behaviour:**

*Marone et al. ( 1990 ); Marone and Kilgore ( 1993 ); Mair and Marone ( 1999 );  
Mair et al. ( 2002 ); Chambon et al. ( 2002 ); Mizogichi et al. ( 2007 )*

## Frictional melting:

*Jeffreys ( 1942 ); McKenzie and Brune ( 1972 ); Richards ( 1977 ); Sibson ( 1977 ); Cardwell et al. ( 1978 ); Allen ( 1979 ); Nielsen et al. ( 2007 )*



Pseudo -  
tachylyte: Fault  
vein ( *Sibson,  
1975* )

**Mechanical lubrication:**

*Spray ( 1993 ); Brodsky and Kanamori ( 2001 ); Kanamori and Brodsky ( 2001 )*

**Acoustic fluidization:**

*Melosh ( 1979, 1996 )*

**Gouge gelation:**

*Goldbsy and Tullis ( 2002 ); Di Toro et al. ( 2004 )*



## **Bi – material interface:**

*Andrews and Ben – Zion ( 1997 ); Harris and Day ( 1997 ); Andrews and Harris ( 2005 ); Ben – Zion ( 2006a, 2006b ); Dunham and Rice ( 2008 )*



MTL: Fractured mylonite,  
cataclasite and gouge



**Humidity effects:**

*Dieterich and Conrad ( 1984 ); Hirose and Bystricky ( 2007 )*

**Characteristic length of surface effects:**

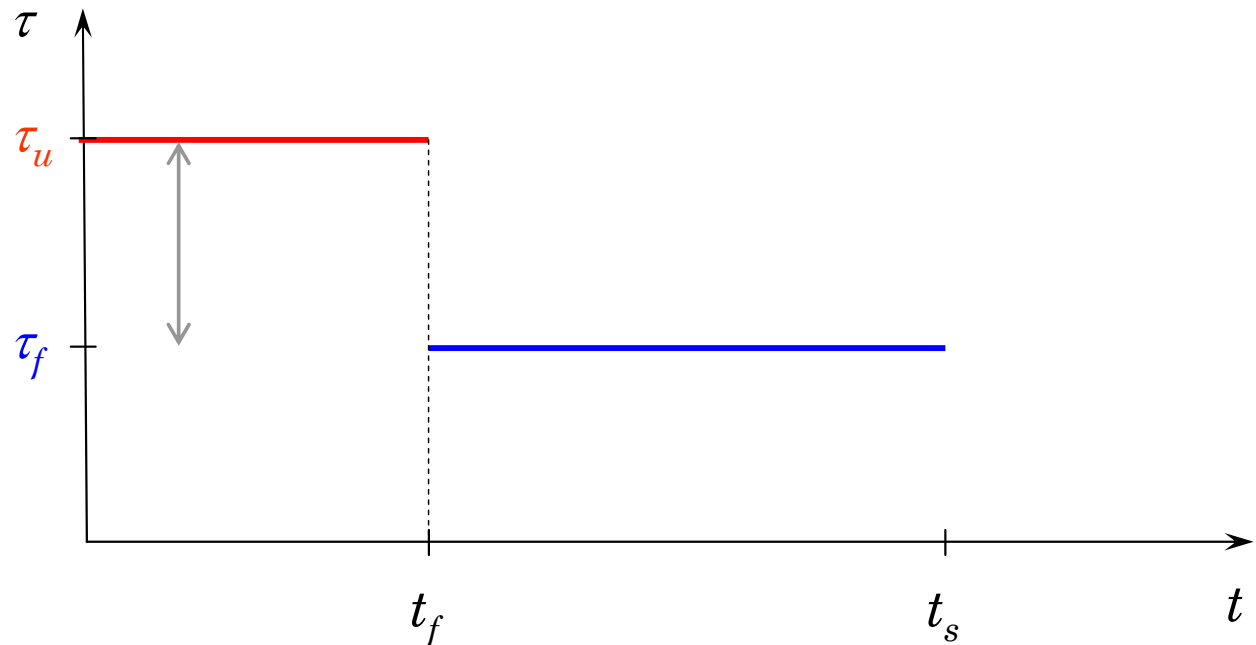
*Ohnaka and Shen ( 1999 ); Ohnaka ( 2003 )*

# Simplest friction models

At a particular fault point  $\xi$  ( following *Savage and Wood, 1971; Scholz, 1990* )

Maximum ( or upper, or yield ) stress

Kinetic ( or frictional ) stress



Strength excess:  $\tau_u - \tau_0 = 0$

Dynamic stress drop:  $\Delta\tau_d = \tau_0 - \tau_f$

In the Dugdale' s model ( *Dugdale, 1960; Barenblatt, 1962* ) the drop occurs when  $u = d_0$ .

Failure time ( or rupture onset )

Rupture arrest

# Simplest friction models

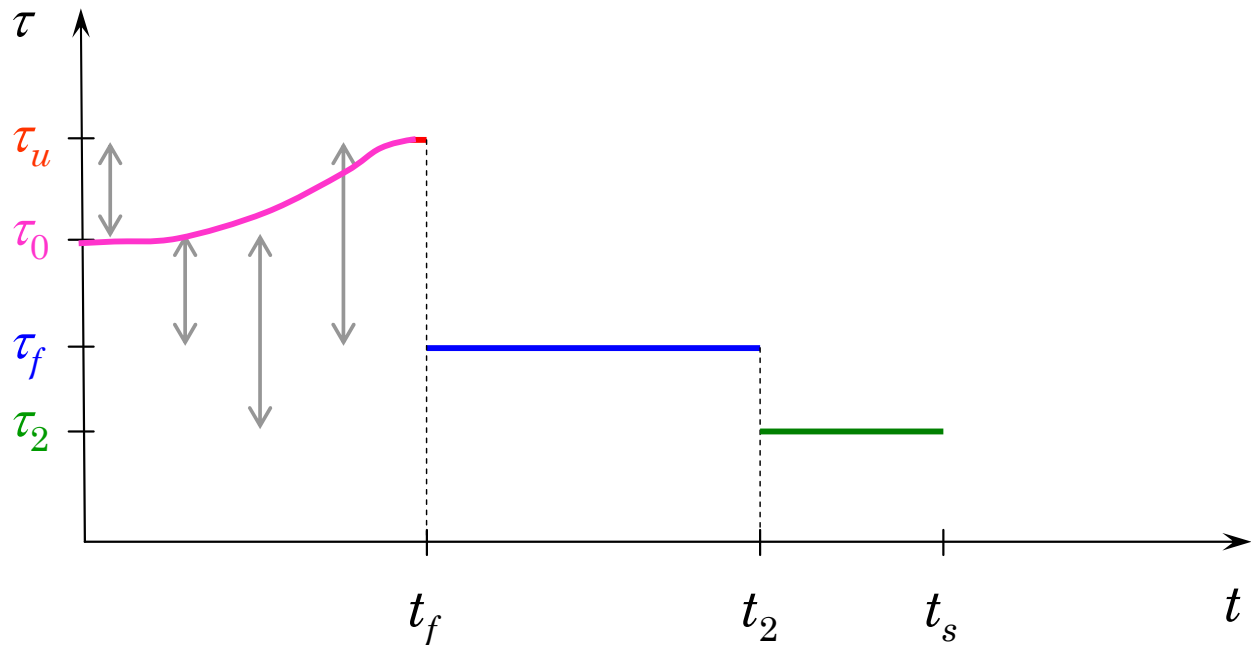
At a particular fault point  $\xi$  ( following *Savage and Wood, 1971; Scholz, 1990* )

Maximum ( or upper, or yield ) stress

Initial stress

Kinetic ( or frictional ) stress

Residual stress



Strength excess:  $\tau_u - \tau_0$

Dynamic stress drop:  $\Delta\tau_d = \tau_0 - \tau_f$

Static stress drop:  $\Delta\tau_s = \tau_0 - \tau_2$

Breakdown str. drop:  $\Delta\tau_b = \tau_u - \tau_f$

Failure time ( or rupture onset )

Dynamic overshoot

Rupture arrest

- *Savage and Wood ( 1971 )* also define:

Mean stress:  $\langle \tau \rangle = \frac{1}{2} (\tau_u + \tau_2)$

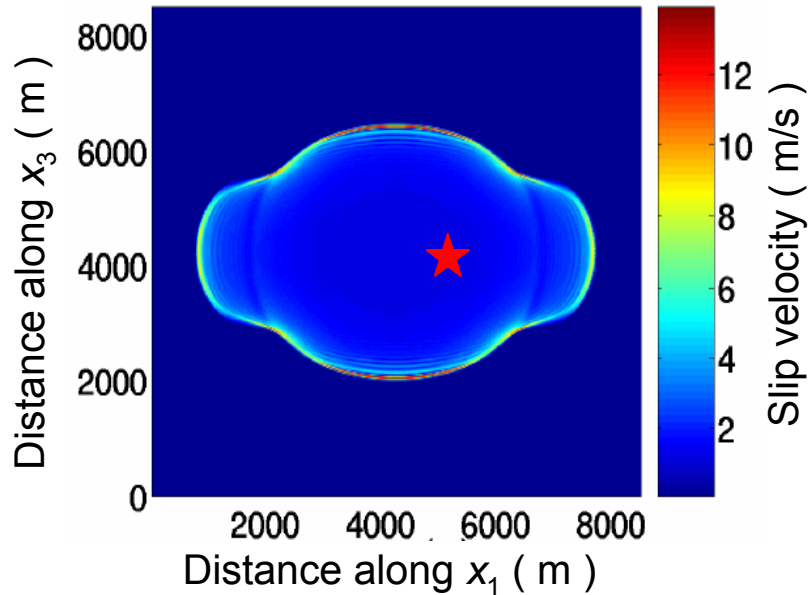
Seismic efficiency:  $\eta = E_s/E$ , where:  $E_s$  is the seismic energy  
 $E$  is the total available energy

Apparent stress:  $\tau_a = \eta \langle \tau \rangle$

- Direct observation of the absolute stress near an earthquake is not feasible, but it is possible ( *Wyss and Brune, 1968* ) calculate  $\tau_a$  and stress drop from physical observables.

# The cohesive zone

Time snapshot  
( t = 0.8 s ) - SW law



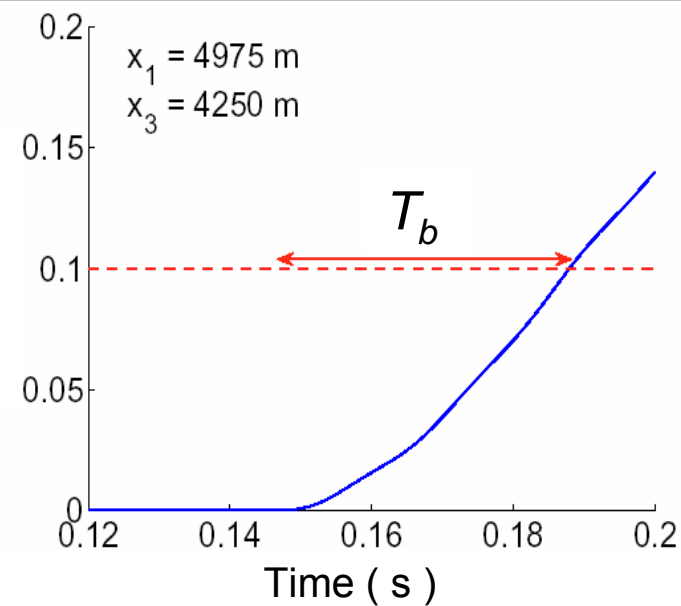
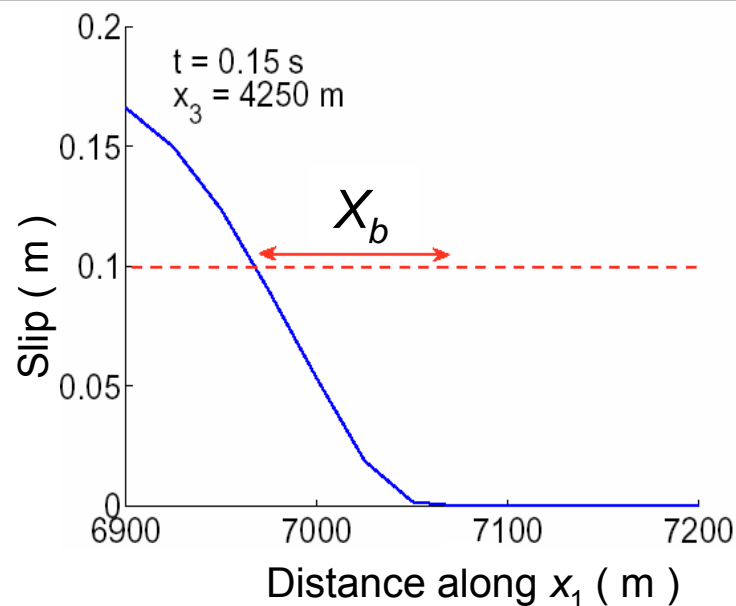
In the target location we can estimate:

$$X_b = 105 \text{ m} \quad T_b = 0.04 \text{ s}$$

From these quantities:

$$V_{rupt} = X_b / T_b = 2625 \text{ m/s}$$

Local estimate



# Slip - hardening effect



- \* The slip – hardening ( **SH** ) phenomenon has been also found in seismological inversion studies ( e. g. *Quin, 1990; Miyatake, 1992; Mikumo and Miyatake, 1993; Beroza and Mikumo, 1996; Ide, 1997; Bouchon, 1997* ).

# Interpretation of the state variable

